

Yukawa Alignment in Two-Higgs-Doublet Models

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1 Introduction

2 Aligned Two-Higgs-Doublet Model

3 Phenomenology

4 Conclusions

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3 Phenomenology

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Introduction

The doublets and the benefits

Two Higgs doublets ϕ_a ($a=1,2$) with $Y = \frac{1}{2}$ whose neutral components acquire VEV's:

$$\langle 0 | \phi_a^T(x) | 0 \rangle = \frac{1}{\sqrt{2}} (0, v_a e^{i\theta_a}) \quad v = \sqrt{v_1^2 + v_2^2} \quad \text{Choice: } \theta_1 = 0, \theta \equiv \theta_2 - \theta_1$$

$$\begin{pmatrix} \Phi_1 \\ -\Phi_2 \end{pmatrix} \equiv \Omega \begin{pmatrix} \phi_1 \\ e^{-i\theta} \phi_2 \end{pmatrix} \quad ; \quad \Omega \equiv \frac{1}{v} \begin{bmatrix} v_1 & v_2 \\ v_2 & -v_1 \end{bmatrix}$$

Higgs basis

$$\Phi_1 = \begin{bmatrix} G^+ \\ \frac{1}{\sqrt{2}}(v + S_1 + iG^0) \end{bmatrix} \quad ; \quad \Phi_2 = \begin{bmatrix} H^+ \\ \frac{1}{\sqrt{2}}(S_2 + iS_3) \end{bmatrix}$$

$$S_1, S_2, S_3 \xrightarrow{\mathcal{R}} H, h, A$$

Benefits of having more than one Higgs doublet

- Present or required in many new-physics scenarios (SUSY)
- Potential new sources of CP symmetry breaking (also Spontaneous CP violation, Axion phenomenology, dark matter candidates...)

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General Two-Higgs-Doublet Model

$$\mathcal{L} = \mathcal{L}^{SM} + \overbrace{T_H + V_H}^{\mathcal{L}_H} + \mathcal{L}_Y$$

$$\mathcal{L}_Y = -\overline{Q}'_L(\Gamma_1\phi_1 + \Gamma_2\phi_2)d'_R - \overline{Q}'_L(\Delta_1\tilde{\phi}_1 + \Delta_2\tilde{\phi}_2)u'_R - \overline{L}'_L(\Pi_1\tilde{\phi}_1 + \Pi_2\tilde{\phi}_2)l'_R + h.c.$$

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Fermion-mass-eigenstate basis $\mathcal{L}_Y[f' \rightarrow f]$ ($f = u, d, l$):

- $M'_f \rightarrow M_f$ diagonal
- $Y'_f \rightarrow Y_f$ NON diagonal and unrelated to M_f

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Avoiding FCNC

How to avoid FCNC:

- Yukawa couplings: $g_{ij} \propto \sqrt{m_i m_j}$ \leftarrow particular Yukawa textures (type III)
- Heavy enough $M_{H_{bosons}}$ \rightarrow suppressed FCNC ('phenomenologically-non-relevant' 2HDM)
- Imposing discrete \mathcal{Z}_2 symmetry

$$\phi_1 \rightarrow \phi_1, \phi_2 \rightarrow -\phi_2, Q_L \rightarrow Q_L, L_L \rightarrow L_L$$

Only one scalar doublet is coupling to a given right-handed fermion field

Different implementations of \mathcal{Z}_2 symmetry

- ϕ_2 to all-fermions (type I)
- ϕ_1 to d and l and ϕ_2 to u (type II)
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Since \mathcal{Z}_2 is scalar-basis dependent:

- Φ_1 to all-fermions required! (inert or dark model) \rightarrow natural frame for Dark Matter

NO-FCNC but also NO new potential CP violating sources

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Alignment in flavor space of the Yukawa couplings of the doublets

$$\Gamma_2 = \xi_d e^{-i\theta} \Gamma_1 \quad \Delta_2 = \xi_u^* e^{i\theta} \Delta_1 \quad \Pi_2 = \xi_l e^{-i\theta} \Pi_1$$

$$Y_{d,l} = \zeta_{d,l} M_{d,l} \quad Y_u = \zeta_u M_u \quad ; \quad \zeta_f \equiv \frac{\xi_f - \tan \beta}{1 + \xi_f \tan \beta} \quad (\tan \beta = \frac{v_2}{v_1})$$

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Recovering usual \mathcal{Z}_2 models, warning and noting

Model	(ξ_d, ξ_u, ξ_l)	ζ_d	ζ_u	ζ_l
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Type II	$(0, \infty, 0)$	$-\tan \beta$	$\cot \beta$	$-\tan \beta$
Type X	$(\infty, \infty, 0)$	$\cot \beta$	$\cot \beta$	$-\tan \beta$
Type Y	$(0, \infty, \infty)$	$-\tan \beta$	$\cot \beta$	$\cot \beta$
Inert	$\tan \beta$	0	0	0

Table 1

Note on ζ_l

Lepton sector: FCNC are identically zero to all orders in PT

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A comment on Radiative Corrections

Alignment Yukawa couplings is not directly protected by any symmetry: radiative FCNC

Nevertheless ...

- $f_i \rightarrow f_i e^{i\theta_{f_i}}$, $V_{ij} \rightarrow e^{i\theta_{u_i}} V_{ij} e^{-i\theta_{d_j}}$
 - Loops cannot generate LPV
 - FCNCs have a particular structure:

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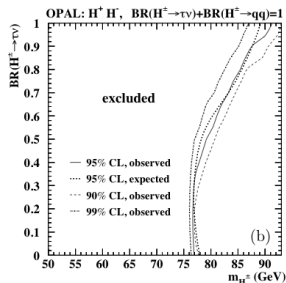
3 Phenomenology

4 Conclusions

Phenomenology in A2HDM

$$H^\pm \rightarrow \tau^\pm \nu_\tau$$

OPAL Collaboration, arXiv:0812.0267 [hep-ex]



M_H (GeV)	> 79.6	> 79.2	> 91.2
B_τ	0	0.5	1

- $H^\pm \rightarrow \tau^\pm \nu_\tau, q \bar{q}'$
- $B_\tau = 0$ (1) $\longleftrightarrow \zeta_I = 0$ ($\zeta_{u,d} = 0$)

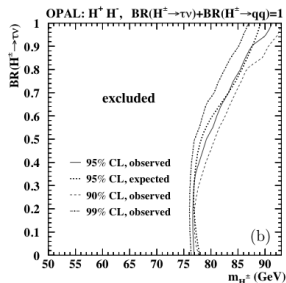
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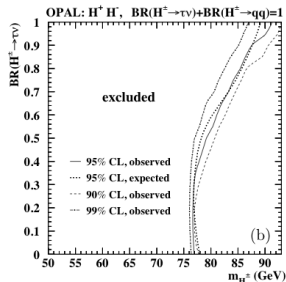
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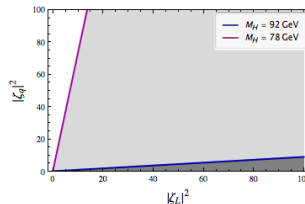
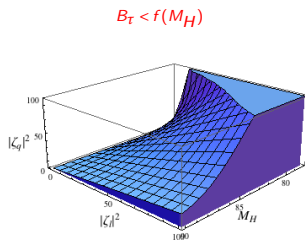


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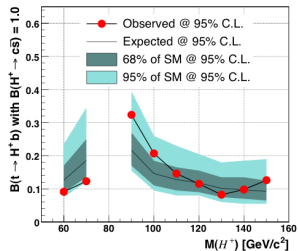
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Phenomenology in A2HDM

$t \rightarrow H^+ b$, $H^+ \rightarrow c\bar{s}$

CDF Collaboration, arXiv:0907.1269 [hep-ex]



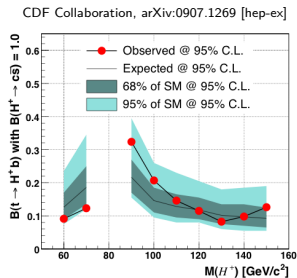
$$|\zeta_u| < 0.5|\zeta_l| + 1$$

$$\bullet B(t \rightarrow H^+ b) \cdot \overbrace{B(H^+ \rightarrow c\bar{s})}^{1-B_\tau}$$

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Phenomenology in A2HDM

$t \rightarrow H^+ b$, $H^+ \rightarrow c\bar{s}$



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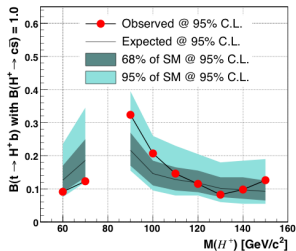
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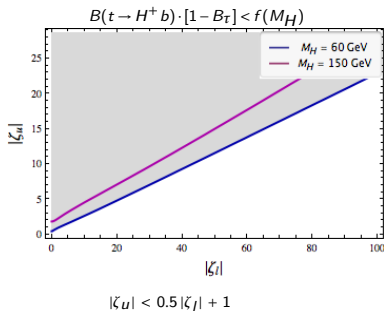
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CDF Collaboration, arXiv:0907.1269 [hep-ex]



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Phenomenology in A2HDM

$$B^\pm \rightarrow \tau^\pm \nu_\tau \text{ and } D_S^\pm \rightarrow \tau^\pm \nu_\tau$$

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$$[\text{M. Antonelli et al '09}] \quad f_B = 216 \pm 22$$

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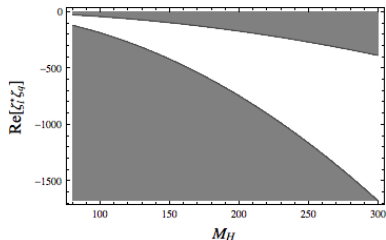
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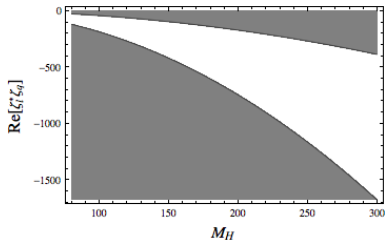
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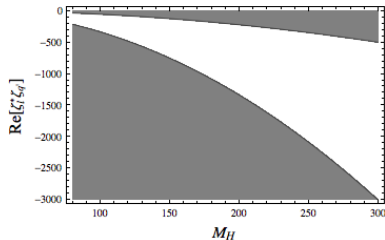


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$M_H = 150 \text{ GeV}, \zeta_{u,d} \text{ reals}$

Work in progress...

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- CP asymmetry

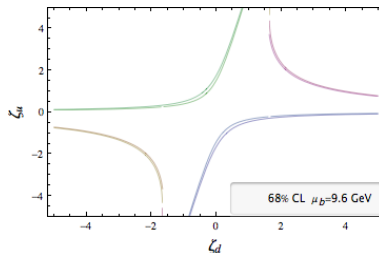
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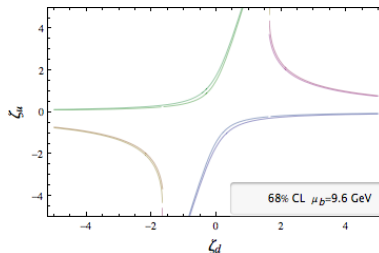
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- 1 Introduction
- 2 Aligned Two-Higgs-Doublet Model
- 3 Phenomenology
- 4 Conclusions**

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- parametrizes the phenomenology with **only three parameters**: ζ_u , ζ_d and ζ_l
 - Z_2 -2HDMs are included in this approach
- These parameters are complex \rightarrow potential new **sources of CP violation**
- The *aligned* Lagrangian has a **symmetry** restricting the (possible) radiative FCNC structures

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Thanks !

General \mathcal{R}

$$y_{d,l}^\phi = \mathcal{R}_{i1} + (\mathcal{R}_{i2} + i\mathcal{R}_{i3})\zeta_{d,l} \quad y_u^\phi = \mathcal{R}_{i1} + (\mathcal{R}_{i2} + -\mathcal{R}_{i3})\zeta_u^*$$

CP-symmetric potential:

$$\begin{aligned} y_{d,l}^H &= \cos(\alpha - \beta) + \sin(\alpha - \beta)\zeta_{d,l} & y_u^H &= \cos(\alpha - \beta) + \sin(\alpha - \beta)\zeta_u^* \\ y_{d,l}^h &= -\sin(\alpha - \beta) + \cos(\alpha - \beta)\zeta_{d,l} & y_u^h &= s \sin(\alpha - \beta) + \sin(\alpha - \beta)\zeta_u^* \\ y_{d,l}^A &= i\zeta_{d,l} & y_u^A &= -i\zeta_u^* \end{aligned}$$