

A Systematical Renormalization of the Type I/II 2HDM and the Automated Generation of NLO Amplitudes

Lukas Altenkamp

In Collaboration with
S. Dittmaier and H. Rzehak

02.09.2014, Workshop on Multi-Higgs Models, Lisbon



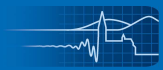
Tests for the scalar sector

Needed:

Accurate experimental measurements & precise theoretical predictions

↑
our contribution

- 1 **Type I/II Two-Higgs Doublet Model**
 - Define parameters at higher order
- 2 **Renormalization**
 - Consistent set of renormalization conditions
- 3 **Implementation**
 - Automatization to be able to calculate many processes in an efficient way
- 4 **Outlook & Application**
 - Higgs decay into 4 fermions



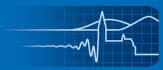
Tests for the scalar sector

Needed:

Accurate experimental measurements & precise theoretical predictions

↑
our contribution

- 1 **Type I/II Two-Higgs Doublet Model**
 - Define parameters at higher order
- 2 **Renormalization**
 - Consistent set of renormalization conditions
- 3 **Implementation**
 - Automatization to be able to calculate many processes in an efficient way
- 4 **Outlook & Application**
 - Higgs decay into 4 fermions



Bare Higgs Lagrangian:

$$\mathcal{L}_{\text{Higgs,bare}} = (D_\mu \Phi_1)^\dagger (D^\mu \Phi_1) + (D_\mu \Phi_2)^\dagger (D^\mu \Phi_2) - V,$$

with

$$\begin{aligned} V = & m_{11} \Phi_1^\dagger \Phi_1 + m_{22} \Phi_2^\dagger \Phi_2 - [m_{12} \Phi_1^\dagger \Phi_2 + h.c.] \\ & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\ & + \frac{1}{2} \lambda_5 [(\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2]. \end{aligned}$$

Constraints:

- ▶ CP conservation
- ▶ No FCNCs at tree level

Expand the fields around the vevs: $\Phi_i = \left(\begin{array}{c} \phi_i^+ \\ \frac{1}{\sqrt{2}}(\eta_i + ic_i^0) + v_i \end{array} \right)$



LO Mass eigenstates:

$$\begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} \cos \beta_{\text{LO}} & -\sin \beta_{\text{LO}} \\ \sin \beta_{\text{LO}} & \cos \beta_{\text{LO}} \end{pmatrix} \begin{pmatrix} G_0 \\ A_0 \end{pmatrix}$$
$$\begin{pmatrix} \phi_1^\pm \\ \phi_2^\pm \end{pmatrix} = \begin{pmatrix} \cos \beta_{\text{LO}} & -\sin \beta_{\text{LO}} \\ \sin \beta_{\text{LO}} & \cos \beta_{\text{LO}} \end{pmatrix} \begin{pmatrix} G^\pm \\ H^\pm \end{pmatrix}$$
$$\begin{pmatrix} \eta_1 \\ \eta_2 \end{pmatrix} = \begin{pmatrix} \cos \alpha_{\text{LO}} & -\sin \alpha_{\text{LO}} \\ \sin \alpha_{\text{LO}} & \cos \alpha_{\text{LO}} \end{pmatrix} \begin{pmatrix} H \\ h \end{pmatrix},$$

with

$$\tan \beta = v_2/v_1$$

Mass mixing terms occur at $\mathcal{O}(\alpha)$:

$$\mathcal{L}_{\text{mass}} = \delta M_{Hh} H h + \delta M_{AG} A_0 G_0 + \delta M_{H^\pm G^\pm} (H^+ G^- + G^+ H^-) \\ + \text{ordinary mass terms}$$



Tadpole terms do not vanish in $\mathcal{O}(\alpha_s)$:

$$\left. \frac{\partial V}{\partial H} \right|_{\text{Fields}=0} = \delta t_H,$$

$$\left. \frac{\partial V}{\partial h} \right|_{\text{Fields}=0} = \delta t_h$$

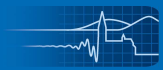
Yukawa Couplings to Fermions:

- So far: Neglect CKM-mixing

Downtype fermions: $\mathcal{L}_{\text{Yukawa}} = \sqrt{2} m_f / v_{1,2} \bar{F}^L f^R \Phi_{1,2}$

Uptype fermions: $\mathcal{L}_{\text{Yukawa}} = \sqrt{2} m_f / v_{1,2} \bar{F}^L f^R \tilde{\Phi}_{1,2}$

	u_i	d_i	e_i
Type I	Φ_2	Φ_2	Φ_2
Type II	Φ_2	Φ_1	Φ_1



Tadpole terms do not vanish in $\mathcal{O}(\alpha_s)$:

$$\left. \frac{\partial V}{\partial H} \right|_{\text{Fields}=0} = \delta t_H,$$

$$\left. \frac{\partial V}{\partial h} \right|_{\text{Fields}=0} = \delta t_h$$

Yukawa Couplings to Fermions:

- So far: Neglect CKM-mixing

Downtype fermions: $\mathcal{L}_{\text{Yukawa}} = \sqrt{2} m_f / v_{1,2} \bar{F}^L f^R \Phi_{1,2}$

Uptype fermions: $\mathcal{L}_{\text{Yukawa}} = \sqrt{2} m_f / v_{1,2} \bar{F}^L f^R \tilde{\Phi}_{1,2}$

	u_i	d_i	e_i
Type I	Φ_2	Φ_2	Φ_2
Type II	Φ_2	Φ_1	Φ_1



THDM

Fields:

$H, h, A_0, H^\pm, G, G^\pm$

parameters:

"input" parameters:

$m_{11}, m_{22}, m_{12}, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5$

In terms of "physical parameters"

$M_H, M_h, M_{A_0}, M_{H^\pm}, M_W, \tan \beta, \sin \alpha_{LO}, \lambda_5$



- ▶ Method 1: Renormalization of the "input" parameters and fields

$$p_i \rightarrow p_i + \delta p_i, \quad \Phi_j \rightarrow \Phi_j + \delta Z_{H_j} \Phi_j, \quad v_j \rightarrow \delta Z_{H_j} (v_j - \delta v_j)$$

- ▶ Method 2: Renormalization of the "physical parameters"

$$p_p \rightarrow p_p + \delta p_p, \quad \delta t_{H/h}, \quad \delta M_{Hh}(\delta \lambda_3), \delta M_{AG}, \delta M_{H^\pm G^\pm}$$

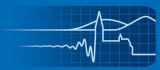
and fields

$$\begin{pmatrix} H \\ h \end{pmatrix} \rightarrow \begin{pmatrix} 1 + \frac{1}{2} \delta Z_H & \frac{1}{2} \delta Z_{Hh} \\ \frac{1}{2} \delta Z_{hH} & 1 + \frac{1}{2} \delta Z_h \end{pmatrix} \begin{pmatrix} H \\ h \end{pmatrix},$$
$$\begin{pmatrix} A_0 \\ G_0 \end{pmatrix} \rightarrow \begin{pmatrix} 1 + \frac{1}{2} \delta Z_{A_0} & \frac{1}{2} \delta Z_{AG} \\ \frac{1}{2} \delta Z_{GA} & 1 + \frac{1}{2} \delta Z_G \end{pmatrix} \begin{pmatrix} A_0 \\ G_0 \end{pmatrix},$$
$$\begin{pmatrix} H^\pm \\ G^\pm \end{pmatrix} \rightarrow \begin{pmatrix} 1 + \frac{1}{2} \delta Z_{H^\pm} & \frac{1}{2} \delta Z_{HG^\pm} \\ \frac{1}{2} \delta Z_{G^\pm H} & 1 + \frac{1}{2} \delta Z_{G^\pm} \end{pmatrix} \begin{pmatrix} H^\pm \\ G^\pm \end{pmatrix},$$

- ▶ Both approaches are equivalent for the parameter renormalization
- ▶ Second approach allows for on-shell scheme in the field renormalization



Renormalization conditions



As in SM:

- ▶ Masses δM_i : Poles of the propagators have to be at the mass

$$\hat{\Gamma}^a(k^2 = M_a^2) = \text{---}^a \text{---} \bigcirc \text{---}^a \Big|_{k^2=M_a^2} = 0$$

- ▶ Fields δZ_i : The residuum of the propagators have to be i

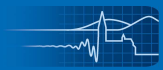
$$\lim_{k^2 \rightarrow M_a^2} \frac{\hat{\Gamma}^a(k)}{k^2 - M_a^2} = \lim_{k^2 \rightarrow M_a^2} \frac{1}{k^2 - M_a^2} \text{---}^a \text{---} \bigcirc \text{---}^a \Big|_{k^2=M_a^2} = i$$

- ▶ Field mixing δZ_{ij} : The mixings of the fields have to vanish on-shell

$$\hat{\Gamma}^{ab}(k^2 = M_a^2 \text{ or } M_b^2) = \Big|_{k^2=M_a^2 \text{ or } M_b^2} = 0$$

- ▶ Tadpoles $\delta t_{H,h}$: All tadpole terms have to vanish

$$\hat{T}_H = \quad = 0$$



As in SM:

- ▶ Masses δM_i : Poles of the propagators have to be at the mass

$$\hat{\Gamma}^a(k^2 = M_a^2) = \text{---}^a \text{---} \text{---} \text{---} \text{---} \left|_{k^2=M_a^2} = 0$$

- ▶ Fields δZ_i : The residuum of the propagators have to be i

$$\lim_{k^2 \rightarrow M_a^2} \frac{\hat{\Gamma}^a(k)}{k^2 - M_a^2} = \lim_{k^2 \rightarrow M_a^2} \frac{1}{k^2 - M_a^2} \text{---}^a \text{---} \text{---} \text{---} \left|_{k^2=M_a^2} = i$$

- ▶ Field mixing δZ_{ij} : The mixings of the fields have to vanish on-shell

$$\hat{\Gamma}^{ab}(k^2 = M_a^2 \text{ or } M_b^2) = \text{---}^{a \neq b} \text{---} \text{---} \left|_{k^2=M_a^2 \text{ or } M_b^2} = 0$$

- ▶ Tadpoles $\delta t_{H,h}$: All tadpole terms have to vanish

$$\hat{T}_H = \text{---} = 0$$



Renormalization conditions



As in SM:

- ▶ Masses δM_i : Poles of the propagators have to be at the mass

$$\hat{\Gamma}^a(k^2 = M_a^2) = \text{---}^a \text{---} \text{---} \text{---} \left(\text{---} \text{---} \right) \text{---}^a \bigg|_{k^2=M_a^2} = 0$$

- ▶ Fields δZ_i : The residuum of the propagators have to be i

$$\lim_{k^2 \rightarrow M_a^2} \frac{\hat{\Gamma}^a(k)}{k^2 - M_a^2} = \lim_{k^2 \rightarrow M_a^2} \frac{1}{k^2 - M_a^2} \text{---}^a \text{---} \left(\text{---} \text{---} \right) \text{---}^a \bigg|_{k^2=M_a^2} = i$$

- ▶ Field mixing δZ_{ij} : The mixings of the fields have to vanish on-shell

$$\hat{\Gamma}^{ab}(k^2 = M_a^2 \text{ or } M_b^2) = \text{---}^{a \neq b} \text{---} \left(\text{---} \text{---} \right) \text{---}^b \bigg|_{k^2=M_a^2 \text{ or } M_b^2} = 0$$

- ▶ Tadpoles $\delta t_{H,h}$: All tadpole terms have to vanish

$$\hat{T}_H = \text{---}^{\text{H,h}} \text{---} \left(\text{---} \text{---} \right) = 0$$



- ▶ $\delta \tan \beta$: \overline{MS} scheme

Condition for the vevs: [Sperling, Stöckinger, Voigt 13]

$$\delta v_1/v_1|_{UV-div} = \delta v_2/v_2|_{UV-div}$$

$\Rightarrow \delta \tan \beta$ includes only contributions from the field renormalization
 $\Rightarrow UV$ -Divergences can be calculated from field renormalization constants:

$$\delta \tan \beta = \frac{\delta Z_h - \delta Z_H}{c_\alpha^2 - s_\alpha^2} \frac{\tan \beta}{2} \Big|_{UV-div}$$

- ▶ $\delta \lambda_3$: \overline{MS} -Scheme

$\delta \lambda_3 = \delta \lambda_3(\delta M_{Hh}, \{\delta \dots\})$ with

$$\delta M_{Hh} \Big|_{UV-div} = -\frac{s_\alpha c_\alpha}{\tan \beta} \delta \tan \beta (M_H^2 - M_h^2) - \Sigma^{Hh}(M_H^2) \Big|_{UV-div},$$



- ▶ $\delta\lambda_5$: \overline{MS} scheme

$$\hat{\Gamma}^{HA_0A_0} \Big|_{\text{UV-div}} = \begin{array}{c} A_0 \\ \text{---} \\ \text{---} \\ A_0 \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} H \\ \text{---} \end{array} \Big|_{\text{UV-div}} = 0$$

The diagram shows a central shaded circle with four external lines. Two lines on the left are dashed and labeled A_0 . One line on the right is dashed and labeled H . The fourth line is a solid vertical line on the far right. The entire diagram is enclosed in a large vertical line on the right side, with the label UV-div at the bottom of this line.

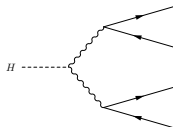


- ▶ MATHEMATICA code to calculate the counterterm potential in dependence of "physical" parameters and renormalization constants
- ▶ Implemented the whole Lagrangian in FEYNRULES
[Alloul, Christensen, Degrande, Duhr, Fuks 09], for generating Feynman Rules and a FEYNARTS [Hahn 99] model file at NLO level
- ▶ A MATHEMATICA routine, which includes the renormalization conditions consistently into the FEYNARTS model file
- ▶ Generation and Calculation of amplitudes and squared matrix elements with FEYNARTS & FORMCALC [Hahn 99]

⇒ Applicable to arbitrary processes! NLO EW & QCD corrections in general possible



Higgs decay into 4 fermions:



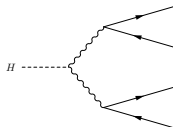
Significant effects in the THDM with a SM like h ?

For SM calculation: Programm PROPHECY4F [Bredestein, Denner, Dittmaier 06]

- ▶ Computation of EW & QCD corrections
- ▶ Arbitrary light final states
- ▶ Subtraction and Slicing method for the cancellation of IR divergences
- ▶ Full off-shell effects
- ▶ Complex mass scheme
- ▶ Generation of unweighted events



Higgs decay into 4 fermions:



Significant effects in the THDM with a SM like h ?

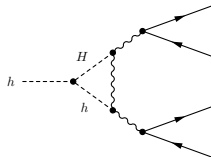
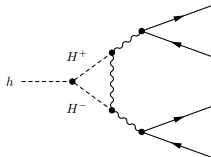
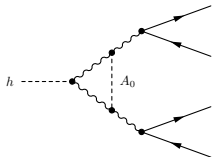
For SM calculation: Programm PROPHECY4F [Bredestein, Denner, Dittmaier 06]

- ▶ Computation of EW & QCD corrections
- ▶ Arbitrary light final states
- ▶ Subtraction and Slicing method for the cancellation of IR divergences
- ▶ Full off-shell effects
- ▶ Complex mass scheme
- ▶ Generation of unweighted events



Implementation in the THDM

- ▶ Use as much of the PROPHECY4F code as possible!
- ▶ LO & photon/gluon emission: Change of the HVV coupling
- ▶ virtual corrections: New diagrams



⇒ Calculation with FEYNARTS & FORMCALC

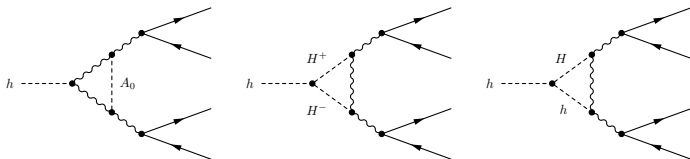
- ▶ Interesting observable: $\Delta\phi_{ij}$, azimuthal angle of the momenta of the fermion pairs

We expect results soon!



Implementation in the THDM

- ▶ Use as much of the PROPHECY4F code as possible!
- ▶ LO & photon/gluon emission: Change of the HVV coupling
- ▶ virtual corrections: New diagrams



⇒ Calculation with FEYNARTS & FORMCALC

- ▶ Interesting observable: $\Delta\phi_{ij}$, azimuthal angle of the momenta of the fermion pairs

We expect results soon!



Thank you for your Attention!



Backup





renormalization conditions: **On-shell scheme** whenever it's possible,
otherwise **\overline{MS} scheme**

parameters	
Electroweak sector:	$\delta M_Z^2, \delta M_W^2, \delta e$
Fermion masses:	$\delta m_{f,i}, \quad f = l, u, d$
Higgs masses:	$\delta M_H^2, \delta M_h^2, \delta M_{A_0}^2, \delta M_{H^\pm}^2$
Higgs potential:	$\delta \lambda_3 (\delta M_{Hh}), \delta \lambda_5, \delta \tan \beta$
Tadpoles:	$\delta t_H, \delta t_h$
fields	
Electroweak sector:	$\delta Z_W, \delta Z_{ZZ}, \delta Z_{ZA}, \delta Z_{AZ}, \delta Z_{ZZ}$
Left fermions:	$\delta Z_i^{f,L}, \quad f = \nu, l, u, d$
Right fermions:	$\delta Z_i^{f,R}, \quad f = l, u, d$
Higgs:	$\delta Z_H, \delta Z_{Hh}, \delta Z_{hH}, \delta Z_h$ $\delta Z_{A_0}, \delta Z_{AG}, \delta Z_{GA}, \delta Z_G$ $\delta Z_{H^\pm}, \delta Z_{HG^\pm}, \delta Z_{GH^\pm}, \delta Z_{G^\pm}$



Renormalization conditions

- ▶ Masses δM_i : Poles of the propagators have to be at the mass

$$\hat{\Gamma}^a(k^2 = M_a^2) = \text{---}^a \text{---} \text{---} \text{---} \left. \vphantom{\hat{\Gamma}^a} \right|_{k^2=M_a^2} = 0$$

- ▶ Fields δZ_i : The residuum of the propagators have to be i

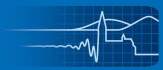
$$\lim_{k^2 \rightarrow M_a^2} \frac{\hat{\Gamma}^a(k)}{k^2 - M_a^2} = \lim_{k^2 \rightarrow M_a^2} \frac{1}{k^2 - M_a^2} \text{---}^a \text{---} \text{---} \left. \vphantom{\hat{\Gamma}^a} \right|_{k^2=M_a^2} = i$$

- ▶ Field mixing δZ_{ij} : The mixings of the fields have to vanish on-shell

$$\hat{\Gamma}^{ab}(k^2 = M_a^2 \text{ or } M_b^2) = \left. \vphantom{\hat{\Gamma}^{ab}} \right|_{k^2=M_a^2 \text{ or } M_b^2} = 0$$

- ▶ electric charge δe : Corrections to $\gamma \bar{f} f$ -vertex in the Thomson limit have to vanish

$$\left. \vphantom{\Gamma} \right|_{p=p', p^2=p'^2=m_e^2} = ie \bar{u}(p) \gamma_\mu u(p)$$



Renormalization conditions

- ▶ Masses δM_i : Poles of the propagators have to be at the mass

$$\hat{\Gamma}^a(k^2 = M_a^2) = \text{---}^a \text{---} \text{---} \text{---} \left. \vphantom{\hat{\Gamma}^a} \right|_{k^2=M_a^2} = 0$$

- ▶ Fields δZ_i : The residuum of the propagators have to be i

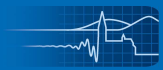
$$\lim_{k^2 \rightarrow M_a^2} \frac{\hat{\Gamma}^a(k)}{k^2 - M_a^2} = \lim_{k^2 \rightarrow M_a^2} \frac{1}{k^2 - M_a^2} \text{---}^a \text{---} \text{---} \left. \vphantom{\hat{\Gamma}^a} \right|_{k^2=M_a^2} = i$$

- ▶ Field mixing δZ_{ij} : The mixings of the fields have to vanish on-shell

$$\hat{\Gamma}^{ab}(k^2 = M_a^2 \text{ or } M_b^2) = \text{---}^{a \neq b} \text{---} \text{---} \left. \vphantom{\hat{\Gamma}^{ab}} \right|_{k^2=M_a^2 \text{ or } M_b^2} = 0$$

- ▶ electric charge δe : Corrections to $\gamma \bar{f} f$ -vertex in the Thomson limit have to vanish

$$\left. \vphantom{\gamma \bar{f} f} \right|_{p=p', p^2=p'^2=m_e^2} = ie \bar{u}(p) \gamma_\mu u(p)$$



Renormalization conditions

- ▶ Masses δM_i : Poles of the propagators have to be at the mass

$$\hat{\Gamma}^a(k^2 = M_a^2) = \text{---}^a \text{---} \left(\text{shaded circle} \right) \text{---}^a \text{---} \Big|_{k^2=M_a^2} = 0$$

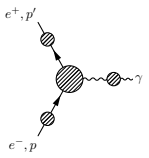
- ▶ Fields δZ_i : The residuum of the propagators have to be i

$$\lim_{k^2 \rightarrow M_a^2} \frac{\hat{\Gamma}^a(k)}{k^2 - M_a^2} = \lim_{k^2 \rightarrow M_a^2} \frac{1}{k^2 - M_a^2} \text{---}^a \text{---} \left(\text{shaded circle} \right) \text{---}^a \text{---} \Big|_{k^2=M_a^2} = i$$

- ▶ Field mixing δZ_{ij} : The mixings of the fields have to vanish on-shell

$$\hat{\Gamma}^{ab}(k^2 = M_a^2 \text{ or } M_b^2) = \text{---}^{a \neq b} \text{---} \left(\text{shaded circle} \right) \text{---}^b \text{---} \Big|_{k^2=M_a^2 \text{ or } M_b^2} = 0$$

- ▶ electric charge δe : Corrections to $\gamma \bar{f} f$ -vertex in the Thomson limit have to vanish



$$\Big|_{p=p', p^2=p'^2=m_e^2} = ie\bar{u}(p)\gamma_\mu u(p)$$