Renormalization Group Analysis of the xSM model and implementation in the ScannerS Tool

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Outline

- Introduction
 - Motivation
 - The xSM Model
- Renormalization of the xSM
 - The effective potential
 - Renormalization Group Equations
- Scans over the Parameter Space
 - Theoretical and Experimental Constraints
 - Dark Matter Phase
 - Broken Phase
- Conclusions



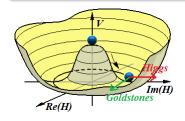
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The Standard Model $\rightarrow SU(3)_c \times SU(2)_L \times U(1)_Y$

$$\mathcal{L}_{SM} = -\frac{1}{4}F_{\mu\nu}^{a}F^{a\mu\nu} + i\overline{\psi}\not{D}\psi + h.c. - \overline{\psi}_{i}(y_{i})_{ij}\psi_{j}\phi + h.c. + (D^{\mu}H)^{\dagger}(D_{\mu}H) - V(H^{\dagger}H)$$



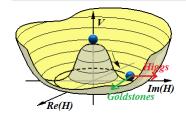
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$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} G^+ \\ v + \frac{h}{h} + iG^0 \end{pmatrix}$$
, $v = 246 \text{ GeV}$



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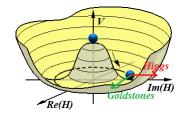
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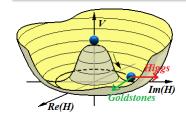
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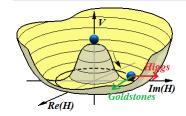
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- Charge quantization
- Fermion masses and mixings
- Hard to reconcile with the theory of General Relativity



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- ② $v_A \neq 0 \rightarrow broken\ phase$: A no longer dark $\rightarrow \frac{\delta_2 v_A}{4} A h^2 + \frac{d_2 v_A}{4} A S^2 \rightarrow$ S, A and h mix



Symmetric or Dark Matter phase:



$$\begin{pmatrix} H_1 \\ H_2 \\ A' \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} h \\ S \\ A \end{pmatrix}$$



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Broken phase:



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② Broken phase: → 3 mixed scalars

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$$\bullet \ \kappa_j \equiv \frac{\lambda_{H_j}^{(p)}}{\lambda_{h_{SM}}^{(p)}}$$

- $\lambda_{h_{SM}}^{(p)}$ Coupling of particle p to the SM Higgs
- $\lambda_{H_i}^{(p)}$ Coupling of particle p to the new scalar



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• Effective potential $V_{\it eff}$ defined as the effective energy density at a constant field ϕ_0

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- Scale invariance of $G^{(n)}(\phi_1,\ldots,\phi_n)$ (Callan-Symanzyk eq.)



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- V_{eff} is scale invariant



What is our aim?

Investigate the stability of the xSM model with energy/RG scale μ

- Use the scale invariance of the effective potential to derive the RGEs
- Apply stability constraints to the RG evolved couplings using ScannerS (bottom-up approach)
- ullet For each generated point, determine the scale up to where V_{xSM} is stable



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Generic loop expansion of V_{eff}

$$V_{eff} = \sum_{n=0}^{+\infty} \varepsilon^n V^{(n)} (L, \mathbf{v}_i, \mu) , \quad \varepsilon = \frac{\hbar}{16\pi^2}$$



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• 1-loop \longrightarrow truncate to first order n = 1

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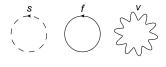


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$$V^{(1)} = rac{1}{64\pi^4} \sum_i n_i \operatorname{Tr} \left[m_i^4 \left(\log rac{m_i^2}{\mu^2} - k_i
ight)
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$$\hbar = 1$$
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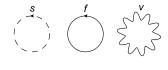




$$V_{eff} = V^{(0)} + rac{\hbar}{16\pi^2}V^{(1)} + \mathcal{O}(\hbar^2)$$

- $V^{(0)}$ is V_{xSM} evaluated at the minimum
- We use for $V^{(1)}$ the Coleman-Weinberg potential

$$V^{(1)} = rac{1}{64\pi^4} \sum_i n_i \, Tr \left[m_i^4 \left(\log rac{m_i^2}{\mu^2} - k_i
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 - Ongoing work: Full two-loop RGEs for xSM



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Theoretical Constraints

- Check boundedness from below of the scalar potential (all scales)
- Check perturbative unitarity in 2 → 2 processes (all scales)

$$|\lambda| \leq 16\pi, |d2| \leq 16\pi, |\delta_2| \leq 16\pi, \left| \frac{3}{2}\lambda + d_2 \pm \sqrt{\left(\frac{3}{2}\lambda + d_2\right)^2 + d_2^2} \right| \leq 16\pi$$



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- Check that the minimum we chose is global (low scale only)
- Compare with EW precision observables S, T, U (low scale only)



Collider searches for SM Higgs boson



Collider searches for SM Higgs boson

- ullet Experiments provide exclusion limits on various signal strengths μ_i normalized to the SM
- Cross sections σ and decay widths Γ for the new scalars are rescaled by κ_i^2

$$\mu_i = \frac{\sigma_{\textit{new}}(\textit{H}_i)\textit{Br}_{\textit{new}}\left(\textit{H}_i \rightarrow \textit{X}_{\textit{SM}}\right)}{\sigma_{\textit{SM}}(\textit{h}_{\textit{SM}})\textit{Br}_{\textit{SM}}\left(\textit{h}_{\textit{SM}} \rightarrow \textit{X}_{\textit{SM}}\right)} = \kappa_i^2 \frac{\textit{Br}_{\textit{new}}\left(\textit{H}_i \rightarrow \textit{X}_{\textit{SM}}\right)}{\textit{Br}_{\textit{SM}}\left(\textit{h}_{\textit{SM}} \rightarrow \textit{X}_{\textit{SM}}\right)}$$

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- Provide Γ for each scalar to HiggsBounds and HiggsSignals
 - exclude points with μ_i above limits
 - exclude points beyond 3σ

ade

Dark Matter searches

- $\Omega_{cdm}h^2 = 0.1196 \pm 0.0031$ (WMAP)
- LUX bounds
- micrOMEGAS (v2.4.5) to calculate relic density $\Omega_A h^2$ and exclude a point if above WMAP bounds
- Calculate σ_{scaled} for SI WIMP-nucleon scattering
- Reject point if $\sigma_{scaled} > \sigma_{LUX}$ with $\sigma_{scaled} = \sigma_A \frac{\Omega_A h^2}{0.1196}$



Input Parameters

Broken phase inputs

- Fixed input parameters: v = 246 GeV, $m_h = 125.7 \text{ GeV}$
- Free parameters: $v_A, v_S, m_{H_{1,2}} \in [0, 500] \text{ GeV }, \kappa_{1,2,3} \in [0, 1]$
- ullet Our convention $o m_{H_{1,2}} \equiv m_{H_{light\,,heavy}}$

Dark Matter phase inputs

- Fixed input parameters: $v_A = 0 \text{ GeV}$
- \bullet Free parameters: $\textit{a}_{1} \in \left[-10^{8}\,,0\right]~\mathrm{GeV}^{3}$, $\phi \in \left[0\,,1\right]$
- Our convention $\rightarrow m_{H_{1,2}} \equiv m_{H_{DM,new}}$

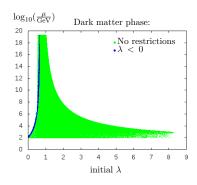


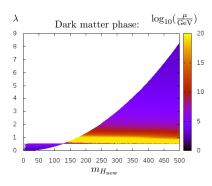
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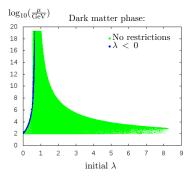


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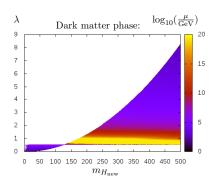




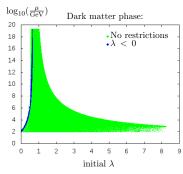




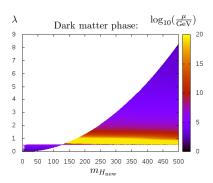
• Run from $\mu = M_z$ to $M_{Pl} \sim 10^{19} \ {\rm GeV}$



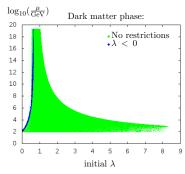


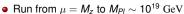


- Run from $\mu = M_z$ to $M_{Pl} \sim 10^{19} \text{ GeV}$
- Perturbative unitarity conditions dominant
- ullet Blue layer o V becomes UFB

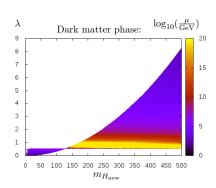




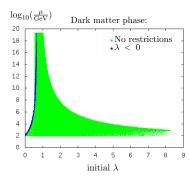




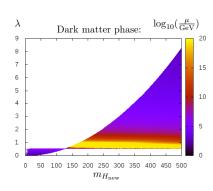
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- \bullet Favours a new visible particle with mass $m_{H_{new}} \gtrsim 140~\text{GeV}$ if we insist in RG stability







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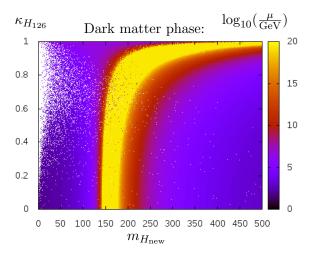


Minimum conditions + Higgs mass

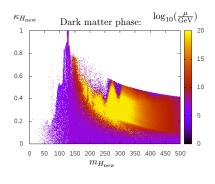
$$\lambda = \frac{1}{v^2} \left[m_{H_{\rm new}}^2 + m_{H_{126}}^2 \pm \sqrt{(m_{H_{\rm new}}^2 - m_{H_{126}}^2) - (v_s \delta_2)^2} \right]$$

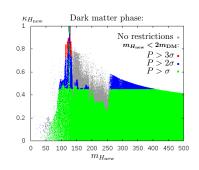
• No mixing limit ($v_s \rightarrow 0$) we obtain $\lambda = 2m_H^2 / v^2$ or $\lambda = 2m_{H_{co}}^2 / v^2$



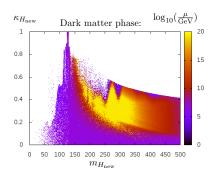


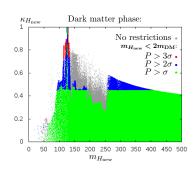






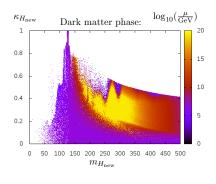


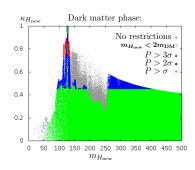




• Contour cut from HiggsBounds and P value from Higgs signals

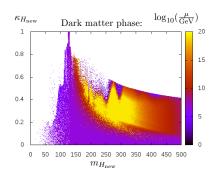


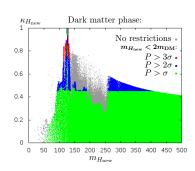




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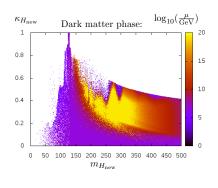


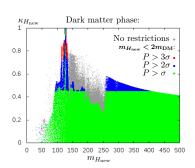




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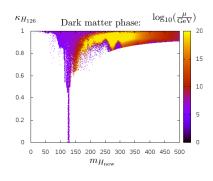


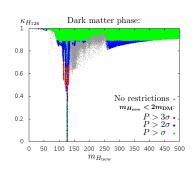




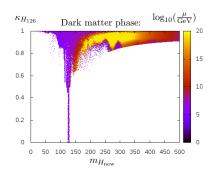
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- Grey points \rightarrow $P > 3\sigma$ and new particle decays to DM

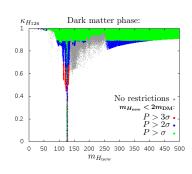






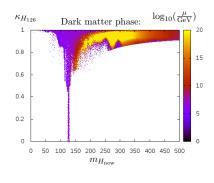


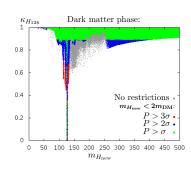




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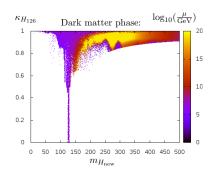


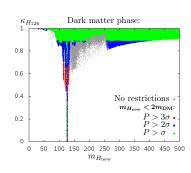


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Dark Matter Phase: Phenomenological tests

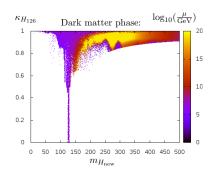


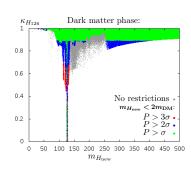


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Dark Matter Phase: Phenomenological tests





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- If not at the decoupling limit ($\kappa_{\textit{H}_{new}} \to 0$ and $\kappa_{\textit{H}_{126}} \to 1$) region may be probed at the 14 TeV LHC runs



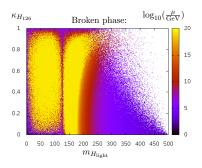
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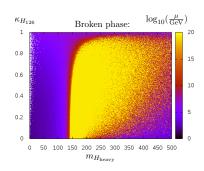
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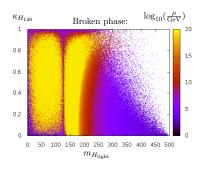
Broken Phase: Theoretical tests

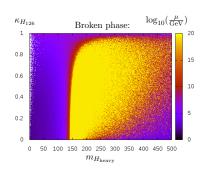






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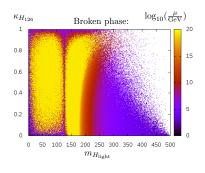


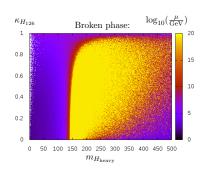


Three states mixing



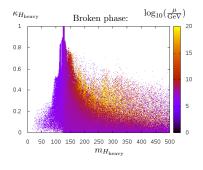
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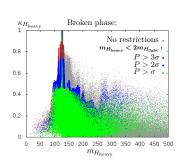




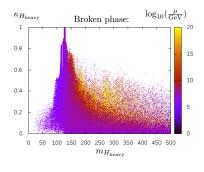
- Three states mixing
- ullet New heavy particle predicted to be $m_{H_{heavy}} \gtrsim 140~{
 m GeV}$ if we insist in RG stability

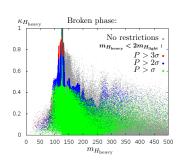






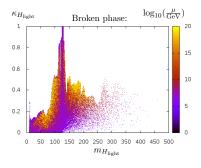


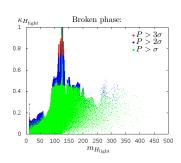




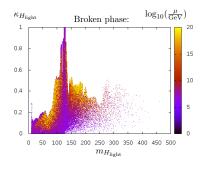
- Harder to find RG stable points with $P > \sigma$ in comparison with the broken phase
- For $\kappa_{H_{heavy}} \sim$ 0.4 and $P > \sigma \rightarrow$ stable up to $\sim 10^{15}~{\rm GeV}$

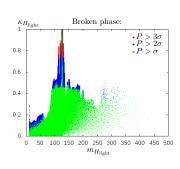












- Coupling
- $\bullet~$ For $\kappa_{\textit{H}_{\textit{heavy}}} \sim \text{0.4}$ and $\textit{P} > \sigma$ stable up to $\sim 10^{15}~\text{GeV}$



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Decay widths for $H_i \rightarrow H_j H_k$ and $H_i \rightarrow H_j H_j$

$$\Gamma(H_{i} \to H_{j}H_{k}) = \frac{g_{ijk}^{2}}{16\pi m_{i}} \sqrt{1 - \frac{(m_{j} + m_{k})^{2}}{m_{i}^{2}}} \sqrt{1 - \frac{(m_{j} - m_{k})^{2}}{m_{i}^{2}}}$$

$$\Gamma(H_{i} \to H_{j}H_{j}) = \frac{g_{ijj}^{2}}{32\pi m_{i}} \sqrt{1 - \frac{4m_{j}^{4}}{m_{i}^{2}}}$$

