

Multi-Higgs aspect of 3-3-1 models

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- Introduction
- 3-3-1 models
- m331 and HL331
 - The scalar sector
 - fermions
 - leptons
 - quarks
 - Gauge fields
- Predictions of the models
- Reduction to the SM
- Phenomenological aspects
- Conclusions

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Almost all the interesting extensions of the standard model (SM) include extra scalar multiplets: a real [Hill 1987, Davoudiasl 2004], or complex [Chikasige 1980] scalar singlet, one or more extra doublets [Branco 2011, this workshop], Hermitian [Brdar 2013] or non-Hermitian triplet [Cheng1980], single charged h^+ [Zee1980,1985] or doubly charged singlet k^{++} [Babu 1988], exotic doublet $(\phi^{++} \phi^+)^T$ [Law 2013],

.
Some of them appear in models which extend the electroweak gauge symmetry to $SU(3)_C \otimes SU(3)_L \otimes U(1)_X$, called by short 3-3-1 models

- Models with $SU(3)_C \otimes SU(3)_L \otimes U(1)_X$ gauge symmetry broken in two steps.



$$\begin{array}{ccc}
 SU(3)_C \otimes SU(3)_L \otimes U(1)_X & \xrightarrow{v_\chi} & SU(3)_C \otimes SU(2)_L \otimes U(1)_Y \\
 & \xrightarrow{v_\eta, v_\rho, v_{\sigma_2}} & SU(3)_C \otimes U(1)_{em} \quad (1)
 \end{array}$$

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- There are several representation contents

- Two triplets are the minimal number to break the 331 symmetry down to 321 symmetry: χ and one of η, ρ and σ_2 . Fermion masses require the other two.
- Three neutral gauge bosons A, Z and Z'
- In addition to $W_\mu^\pm$ there are dilepton gauge bosons $V_\mu^\pm$ and $U_\mu^{\pm\pm}$ that carry lepton number $L(V_\mu^+) = L(U_\mu^{++}) = -2$

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Two 3-3-1 models

- electric charge operator is given by

$$\frac{Q}{|e|} = T_3 - \sqrt{3}T_8 + X$$

where e is the electron charge, $T_{3,8} = \lambda_{3,8}/2$, and X is the hypercharge operator associated to the $U(1)_X$ group.

- the minimal 3-3-1 (m331 for short): the lepton sector is the same as in the standard model.
- the model with heavy leptons (HL331 for short) in which heavy leptons are introduced.
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The scalar sector

$$\eta = \begin{pmatrix} \Phi_1 \\ \eta_2^+ \end{pmatrix} \sim (\mathbf{3}, 0), \rho = \begin{pmatrix} \Phi_2 \\ \rho^{++} \end{pmatrix} \sim (\mathbf{3}, 1), \chi = \begin{pmatrix} \Phi_C \\ \chi^0 \end{pmatrix} \sim (\mathbf{3}, -1).$$

The triplet η and the sextet S

$$S = \begin{pmatrix} T & \frac{\Phi_3}{\sqrt{2}} \\ \frac{\Phi_3}{\sqrt{2}} & H_2^{--} \end{pmatrix}, \quad T = \begin{pmatrix} \sigma_1^0 & \frac{h^+}{\sqrt{2}} \\ \frac{h^+}{\sqrt{2}} & H_1^{++} \end{pmatrix},$$

couple to leptons. Under $SU(2) \otimes U(1)_Y$:

$$\Phi_1 = \begin{pmatrix} \eta^0 \\ -\eta_1^- \end{pmatrix}, \Phi_2 = \begin{pmatrix} \rho^+ \\ \rho^0 \end{pmatrix}, \Phi_3 = \begin{pmatrix} h^+ \\ h^0 \end{pmatrix}, \Phi_C = \begin{pmatrix} \chi^- \\ \chi^{--} \end{pmatrix},$$

$Y = -1, +1, +1, -3$; $\eta_2^+, \rho^{++}, \chi^0, H_2^{--}$ are singlet of $SU(2)$ and have $Y = +2, +4, 0, +4$, respectively and carry lepton number:

$$L(T, \eta_2^-) = -2, L(\Phi_C, \rho^{--}, H_2^{--}) = +2, L(\Phi_{1,2,3}) = 0$$

The lepton number

LV explicitly via soft and hard terms. Or, $v_{\sigma_1^0} \neq 0$ inducing spontaneously breaking of L . Three possibilities (Liu & Ng 1993):

- 1 $\langle \sigma_1^0 \rangle = 0$ and no explicit LV . But V_μ^- and U_μ^{--} violate $L_{e,\mu,\tau}$. $\langle \sigma_1^0 \rangle = 0$ stable at tree level (Foot et al 1992, Pleitez & Tonasse 1993).
- 2 No explicit LV but $\langle \sigma_1^0 \rangle \neq 0 \rightarrow$ a triplet Majoron: ruled out by the Z invisible width.
- 3 Explicit LV . $\langle \sigma_1^0 \rangle = 0$ at tree level, but $\langle \sigma_1^0 \rangle \neq 0$ by loop corrections (Pleitez, Tonasse 1993; Framton, Krastev & Liu 1994).
- 4 $\langle \sigma_1^0 \rangle = 0$ and no explicit LV . $\langle \sigma_2^0 \rangle = 0$ but loop corrections $\Rightarrow \langle \sigma_2^0 \rangle \neq 0$

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The scalar potential: three triplets

$$V(\eta, \rho, \chi) = V^{(2)} + V^{(3)} + V^{(4)},$$

where

$$\begin{aligned} V^{(2)} &= \sum_{x=\eta, \rho, \chi} \mu_X^2 X^\dagger X, \\ V^{(3)} &= \alpha \epsilon \eta \rho \chi \\ V^{(4)} &= a_1 (\chi^\dagger \chi)^2 + a_2 (\eta^\dagger \eta)^2 + a_3 (\rho^\dagger \rho)^2 + \chi^\dagger \chi (a_4 \eta^\dagger \eta + a_5 \rho^\dagger \rho) \\ &\quad + a_6 (\eta^\dagger \eta) (\rho^\dagger \rho) + a_7 (\chi^\dagger \eta) (\eta^\dagger \chi) + a_8 (\chi^\dagger \rho) (\rho^\dagger \chi) + a_9 (\eta^\dagger \rho) (\rho^\dagger \eta) \\ &\quad + a_{10} [(\chi^\dagger \eta) (\rho^\dagger \eta) + H.c.] \end{aligned} \tag{2}$$

$a_{10} \neq 0 \Rightarrow$ explicit LV .

three triplets $\rightarrow$ 2HDM+1 neutral singlet

$$\begin{aligned} V_{2HD+1HS} = & \mu_\eta^2 \Phi_1^\dagger \Phi_1 + \mu_\rho^2 \Phi_2^\dagger \Phi_2 + a_2 (\Phi_1^\dagger \Phi_1)^2 + a_3 (\Phi_2^\dagger \Phi_2)^2 \\ & + a_6 [(\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1)] \\ & + \alpha [(\Phi_1)_1(\Phi_2)_2 - (\Phi_1)_2(\Phi_2)_1] \chi^0 \\ & + \mu_\chi^2 \chi^{0*} \chi^0 + a_1 (\chi^{0*} \chi^0)^2 + \dots + H.c. \end{aligned} \quad (3)$$

... mean interactions involving χ^0 , H_2^{--} , ...

three triplets and the sextet $\rightarrow$ 3HDM

$$V(\eta, \rho, \chi, S) = V^{(2)} + V^{(3)} + V^{(4a)} + \dots + V^{(4e)},$$

$$\begin{aligned}
 V^{(2)} &= \sum_X \mu_X^2 \text{Tr}(X^\dagger X), X = \eta, \rho, \chi, S \\
 V^{(3)} &= \alpha_1 \epsilon \eta \rho \chi + \alpha_2 \chi^T S^\dagger \rho + \alpha_3 \eta^T S^\dagger \eta + \alpha_4 \text{Tr}(SSS) \\
 &\quad + \frac{\alpha_5}{3!} \epsilon_{ijk} \epsilon_{mnl} S_{im} S_{jn} S_{kl}, \\
 V^{(4a)} &= (a_1, \dots, a_9)(\eta, \rho, \chi) + a_{10}[(\chi^\dagger \eta)(\rho^\dagger \eta) + H.c.], \\
 V^{(4b)} &= b_1 \chi^\dagger S \chi \eta + b_2 \rho^\dagger S \rho \eta + b_3 \eta^\dagger S \chi \rho + H.c., \\
 V^{(4c)} &= c_1 \eta \eta S S + c_2 \chi \rho S S + H.c., \\
 V^{(4d)} &= d_1 (\chi^\dagger \chi) \text{Tr} S^\dagger S + d_2 (\chi^\dagger S S^\dagger \chi) + d_3 (\eta^\dagger \eta) \text{Tr}(S^\dagger S) \\
 &\quad + d_4 (\eta^\dagger S S^\dagger \eta) + d_5 (\rho^\dagger \rho) \text{Tr} S^\dagger S + d_6 (\rho^\dagger S S^\dagger \rho), \\
 V^{(e)} &= e_1 (\text{Tr} S^\dagger S)^2 + e_2 \text{Tr}(S^\dagger S S^\dagger S). \tag{4}
 \end{aligned}$$

L is conserved if $\Rightarrow \alpha_3, \alpha_4, \alpha_5, a_{10}, b_3, c_2 = 0$. This is possible if we impose the discrete symmetry

$$\begin{aligned}\rho &\rightarrow i\rho, \chi \rightarrow i\chi, \eta \rightarrow -\eta, S \rightarrow -S, \\ \Psi &\rightarrow i\Psi, Q_{iL} \rightarrow -iQ_{iL}, Q_{3L} \rightarrow Q_{3L}, \\ u_{\alpha R} &\rightarrow u_{\alpha R}, d_{\alpha R} \rightarrow id_{\alpha R}\end{aligned}$$

$\eta, \rho, \chi, S \Rightarrow$ three-Higgs-Doublet model plus fields T and the singlets χ^0 and H_2^{++} .

However, there are only two quartic non-Hermitian interactions among the three dublets:

$$\begin{aligned}V(\Phi_i)_{3HD} &= \sum_i \mu_i^2 \Phi_i^\dagger \Phi_i + \left[\sum_{i < j} m_{ij}^2 (\Phi_i^\dagger \Phi_j) + H.c. \right] \\ &+ \sum_{i < j} \lambda_{ij} (\Phi_i^\dagger \Phi_j) (\Phi_j^\dagger \Phi_i) + b_2 (\Phi_1^\dagger \Phi_3) (\Phi_1^\dagger \Phi_3) \\ &+ d_4 (\Phi_1^\dagger \Phi_3) (\Phi_1^\dagger \Phi_2) + H.c.\end{aligned}$$

Quarks:

two quark generations transform in different way from the third one:

$$Q'_{mL} = (d_m, -u_m, j_m)_L^T \sim (\mathbf{3}, \mathbf{3}^*, -1/3); \quad m = 1, 2$$

$$Q'_{3L} = (u_3, d_3, J)_L^T \sim (\mathbf{3}, \mathbf{3}, 2/3),$$

$$u_{\alpha R} \sim (\mathbf{3}, \mathbf{1}, 2/3), d_{\alpha R} \sim (\mathbf{3}, \mathbf{1}, -1/3), \quad \alpha = 1, 2, 3,$$

$$j_{mR} \sim (\mathbf{3}, \mathbf{1}, -4/3), J_R \sim (\mathbf{3}, \mathbf{1}, 5/3).$$

Leptons.

$$\text{m331: } \Psi_a = (\nu_a l_a l_a^c)_L^T \sim (1, \mathbf{3}, 0).$$

$$\text{HL331: } \Psi_{aL} = (\nu_a l_a^- E_a^+) \sim (1, \mathbf{3}, 0),$$

$$l_{aR}^- \sim (1, 1, -1), E_{aR}^+ \sim (1, 1, 1), \quad a = e, \mu, \tau.$$

Right-handed neutrinos may be added.

Predictions of the models

- 1 The number of generations are three: anomaly cancellation and asymptotic freedom (or at least a multiple of three)
- 2 The electroweak mixing angle $\sin^2 \theta_W < 1/4$
- 3 A Landau-like pole at energies of the order of few TeVs when $\sin^2 \theta_W = 1/4$
- 4 This pole stabilizes the electroweak scale
- 5 The asymmetric treatment of the third quark family which may be related to the heaviness of the top quark
- 6 The electric charge is quantized (anomaly cancellation and the fermion's masses arise through the Yukawa interactions independently of the existence of right-handed neutrinos. In the SM with three generations this occurs only if neutrinos are right-handed and a Majorana mass term is added)
- 7 The models have an almost automatic Peccei-Quinn symmetry

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Reduction to the standard model

$$v_\chi \rightarrow \infty \Rightarrow 331 \rightarrow 321$$

$$\begin{aligned} Z_1 &\rightarrow Z, & Z_2 &\rightarrow Z' \\ g_{V,A}^{Z_1}(s_W^2, \bar{v}_\eta^2, \bar{v}_\rho^2) &\rightarrow g_{V,A}^Z, \\ f_{V,A}^{Z_2}(s_W^2, \bar{v}_\eta^2, \bar{v}_\rho^2) &\rightarrow f_{V,A}^{Z'}(s_W^2), \end{aligned}$$

another non-trivial reduction to SM:

$$\begin{aligned} \rho_1 &= \frac{M_{Z_1}}{c_W M_W} = 1 \Rightarrow \\ v_\rho^2 &= \frac{1 - 4s_W^2}{2c_W^2} v_W^2 \Rightarrow v_\rho \approx 54 \text{ GeV} \\ v_\eta^2 &= \frac{1 + 2s_W^2}{2c_W^2} v_W^2 \Rightarrow v_\rho \approx 240 \text{ GeV} \end{aligned}$$

v_χ is a free parameter

$$\begin{aligned}
 -\mathcal{L}_Y &= \bar{Q}'_{mL}[G_{m\alpha}U'_{\alpha R}\rho^* + \tilde{G}_{m\alpha}D'_{\alpha R}\eta^*] \\
 &+ \bar{Q}'_{3L}[F_{3\alpha}U'_{\alpha R}\eta + \tilde{F}_{3\alpha}D'_{\alpha R}\rho] + H.c.
 \end{aligned}$$

$$\begin{aligned}
 M^U &= \begin{pmatrix} rG_{11} & rG_{12} & rG_{13} \\ rG_{21} & rG_{22} & rG_{23} \\ F_{31} & F_{32} & F_{33} \end{pmatrix} |v_\eta|, \\
 M^D &= \begin{pmatrix} r^{-1}\tilde{G}_{11} & r^{-1}\tilde{G}_{12} & r^{-1}\tilde{G}_{13} \\ r^{-1}\tilde{G}_{21} & r^{-1}\tilde{G}_{22} & r^{-1}\tilde{G}_{23} \\ \tilde{F}_{31} & \tilde{F}_{32} & \tilde{F}_{33} \end{pmatrix} |v_\rho|.
 \end{aligned}$$

$$V_L^q M^q M^{q\dagger} V_L^{q\dagger} = V_R^q M^{q\dagger} M^q V_R^{q\dagger} = (\hat{M}^q)^2, \quad q = U, D.$$

$$V_L^U = \begin{pmatrix} -0.00032 & 0.07163 & -0.99743 \\ 0.00433 & -0.99742 & -0.07163 \\ 0.99999 & 0.00434 & -0.00001 \end{pmatrix},$$

$$V_L^D = \begin{pmatrix} 0.00273 \rightarrow 0.00562 & (0.03 \rightarrow 0.03682) & -(0.99952 \rightarrow 0.99999) \\ -(0.19700 \rightarrow 0.22293) & -(0.97436 \rightarrow 0.97993) & -0.00001 \\ 0.97483 \rightarrow 0.98039 & -(0.19708 \rightarrow 0.22291) & -(0.00415 \rightarrow 0.00434) \end{pmatrix}$$

and the CKM matrix

$$|V_{CKM}| = \begin{pmatrix} 0.97385 \rightarrow 0.97952 & 0.20134 \rightarrow 0.22714 & 0.00021 \rightarrow 0.00042 \\ 0.20116 \rightarrow 0.22679 & 0.97307 \rightarrow 0.97869 & 0.04116 \rightarrow 0.04166 \\ 0.00849 \rightarrow 0.01324 & 0.03919 \rightarrow 0.04028 & 0.99914 \rightarrow 0.99999 \end{pmatrix} \quad (6)$$

Similar expressions for $V_R^{U,D}$.

$$\mathcal{L}_{Z'} = -\frac{g}{2\cos\theta_W} \sum_{q=U,D} [\bar{q}_L \gamma^\mu K_L^q q_L + \bar{q}_R \gamma^\mu K_R^q q_R] Z'_\mu, \quad (7)$$

where we have defined

$$K_L^q = V_L^q Y_L^q V_L^{q\dagger}, \quad K_R^q = V_R^q Y_R^q V_R^{q\dagger}, \quad q = U, D; \quad (8)$$

with $V_{L,R}^{U,D}$ given in Eqs. (5) and (??) and

$$Y_L^U = Y_L^D = -\frac{1}{2\sqrt{3}h(x)} \text{diag}[-2(1-2x), -2(1-2x), 1] \quad (9)$$

and

$$Y_R^U = -\frac{4x}{\sqrt{3}h(x)} \mathbf{1}_{3\times 3}, \quad Y_R^D = \frac{2x}{\sqrt{3}h(x)} \mathbf{1}_{3\times 3} \quad (10)$$

and $h(x) \equiv (1-4x)^{1/2}$, $x = \sin^2\theta_W$.

$$-\mathcal{L}_{qqh} = \sum_{q=U,D} \bar{q}_L \mathcal{K}^q q_R + \text{mass terms} + H.c.,$$

$$\mathcal{K}^D \approx$$

$$\begin{pmatrix} 10^{-4}\rho^0 - 10^{-6}\eta^0 & 10^{-4}\rho^0 - 10^{-5}\eta^0 & -10^{-4}\rho^0 + 10^{-5}\eta^0 \\ 10^{-6}\rho^0 + 10^{-4}\eta^0 & 10^{-5}\rho^0 + 10^{-3}\eta^0 & -10^{-6}\rho^0 + 10^{-2}\eta^0 \\ 10^{-6}\rho^0 - 10^{-5}\eta^0 & 10^{-6}\rho^0 - 10^{-3}\eta^0 & -10^{-6}\rho^0 + 0.011\eta^0 \end{pmatrix},$$

$$\mathcal{K}^U \approx$$

$$\begin{pmatrix} 0.0099\rho^0 - 10^{-6}\eta^0 & 0.00340\rho^0 - 10^{-5}\eta^0 & 0.0109\rho^0 - 10^{-5}\eta^0 \\ -0.13846\rho^0 + 10^{-7}\eta^0 & 0.0556\rho^0 + 10^{-6}\eta^0 & -0.1521\rho^0 - 10^{-6}\eta^0 \\ 1.9228\rho^0 - 10^{-11}\eta^0 & 0.8656\rho^0 - 10^{-10}\eta^0 & 2.3569\rho^0 - 10^{-10}\eta^0 \end{pmatrix}$$

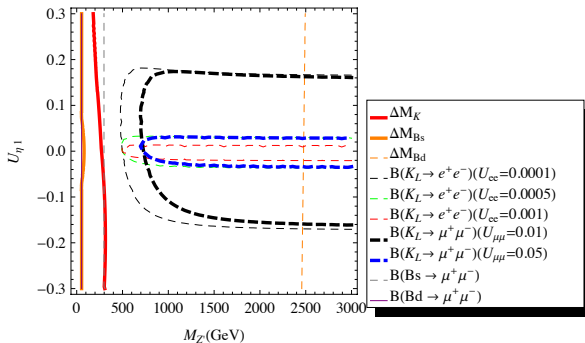


Figure: Allowed regions for $U_{\eta 1}$ and $M_{Z'}$. Taking into account the 125 GeV neutral scalar.

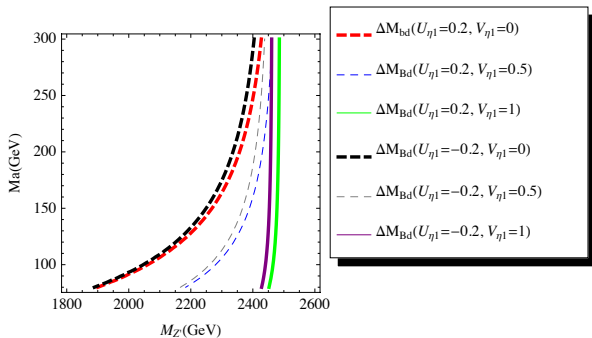


Figure: Allowed region for $M_{Z'}$ and M_A .

$$\begin{aligned}
 V(\eta, \rho, \chi) = & \sum_{x=\eta, \rho, \chi} \mu_X^2 X^\dagger X + \alpha \epsilon \eta \rho \chi + a_1 (\chi^\dagger \chi)^2 + a_2 (\eta^\dagger \eta)^2 \\
 & + a_3 (\rho^\dagger \rho)^2 + \chi^\dagger \chi (a_4 \eta^\dagger \eta + a_5 \rho^\dagger \rho) + a_6 (\eta^\dagger \eta) (\rho^\dagger \rho) \\
 & + a_7 (\chi^\dagger \eta) (\eta^\dagger \chi) + a_8 (\chi^\dagger \rho) (\rho^\dagger \chi) + a_9 (\eta^\dagger \rho) (\rho^\dagger \eta)
 \end{aligned}$$

HL331: three VEVs and α complex: four phases, one survives.

m331: four VEVs complex

The electric dipole moment of electron and neutron have been calculated in the context of the HL331 G. de Conto and Pleitez
arXiv:1408.6551

More hard violating phases in $V_L^{U,D}$ nine phases plus the
constrain $V_L^U V_L^{D\dagger} = V_{CKM}$ (Promberger)

$\langle S \rangle = 0$ at tree level (but there is a divergent contribution to $\langle \sigma_2^0 \rangle$ at one loop level),

$$\mathcal{L}_\nu = \bar{\Psi}_{aL} G_{ab}^{\nu\eta} \nu_{aR} \eta + \overline{(\nu_{aR})^c} (M_R)_{ab} \nu_{bR}$$

M_R is diagonal, and $M_R^{-1} = (1/M) \bar{M}_R$ where $\bar{M} = \text{diag}(M/m_{R1}, M/m_{R2}, 1)$ and $M \equiv m_{R3} > m_{R1}, m_{R2}$.

$$M_{ab}^\nu \approx -\frac{v_\eta^2}{2\bar{M}} G_{ab}^{\nu\eta} \bar{M}_R (G^{\nu\eta})^T, \quad M_{ab}^l = G_{ab}^{l\eta} \frac{v_\eta}{\sqrt{2}} + \frac{1}{\Lambda} G_{ab}^S v_\rho v_\chi.$$

$$|V_{PMNS}| \approx \begin{pmatrix} 0.826 & 0.548 & 0.130 \\ 0.506 & 0.618 & 0.602 \\ 0.249 & 0.563 & 0.788 \end{pmatrix}. \quad (11)$$

3-3-1 models have interesting features: FCNC in both quarks and lepton sectors, (partially) explain the number of generation, but many parameters have to be known before more realistic phenomenological studies are done: it is mandatory to consider vector, scalar and exotic charged fields altogether in order to have a realistic phenomenology. Until now most of the articles consider only one of these particles in a given process. As can be seen from Fig. 1 and 2, there is constructive and destructive interference between the Z' and the SM-like scalar h and one of the pseudoscalar A .

THANKS!

(below you can find some references)

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