

# A Second Higgs

# & the Cosmic *Matter* Asymmetry *Antimatter*

Work with G. Dorsch, S. Huber and K. Mimasu  
1405.5537 (Phys.Rev.Lett 113, 211802) & 1609.xxxxx  
+161x.xxxxx with T. Konstandin

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University of Sussex

US  
University of Sussex



Multi-Higgs, Lisbon, 09/09/2016



# BSM TO DO LIST

Dark Matter

Neutrino Masses

Matter-Antimatter Asymmetry → Baryogenesis

...

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SAKHAROV CONDITIONS *(for dynamical generation of baryon asymmetry)*

*B* Violation

*C/CP* Violation

Departure from Thermal Equilibrium



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SAKHAROV CONDITIONS (for dynamical generation of baryon asymmetry)

$$CP \sim \frac{\prod_{i \neq j}^{u,c,t} |m_i^2 - m_j^2| \times \prod_{i \neq j}^{d,s,b} |m_i^2 - m_j^2|}{T^{12}} \times J \sim 10^{-20}$$

B Violation ✓ **Sphalerons**

Kuzmin, Rubakov, Shaposhnikov, Phys. Lett. B155 (1985) 36

C/CP Violation ✗ **not enough**

Gavela, Hernandez, Orloff, Pene, Quimbay, Nucl. Phys. B430 (1994) 382

Departure from Thermal Equilibrium ✗ **not enough**

Kajantie, Laine, Rummukainen, Shaposhnikov, Phys. Rev. Lett. 77 (1996) 2887

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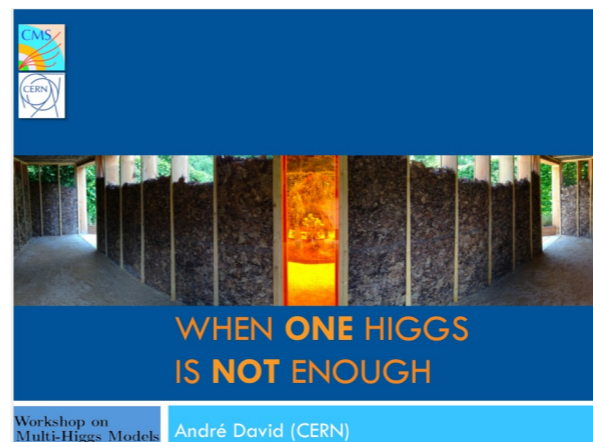
Departure from Thermal Equilibrium ✗ **not enough**

*Kajantie, Laine, Rummukainen, Shaposhnikov, Phys. Rev. Lett. 77 (1996) 2887*

→ EW Phase Transition Smooth CrossOver

→ No CPV in Scalar Sector

## The SM (1 Higgs) is not Enough



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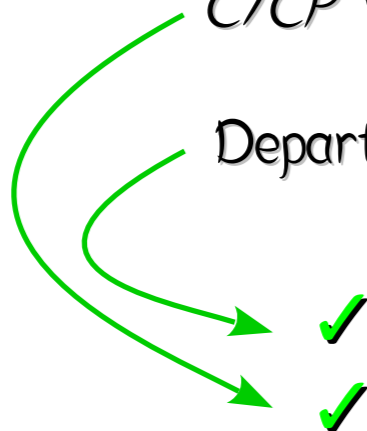
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## More Higgses Help!



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## More Higgses Help!



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# EW Phase Transition

*Universe Expands Adiabatically*  $\Rightarrow$  *Equilibrium Thermal Field Theory*  $\Rightarrow$  *Higgs Finite- $T$  Effective Potential*

$$V_{\text{eff}}(h, T) = V_0(h) + V_0^{\text{loop}}(h) + V_T(h, T)$$

Tree-level  
potential

Loop  
corrections

Thermal  
corrections



# EW Phase Transition

Universe Expands Adiabatically  $\Rightarrow$  Equilibrium Thermal Field Theory  $\Rightarrow$  Higgs Finite- $T$  Effective Potential

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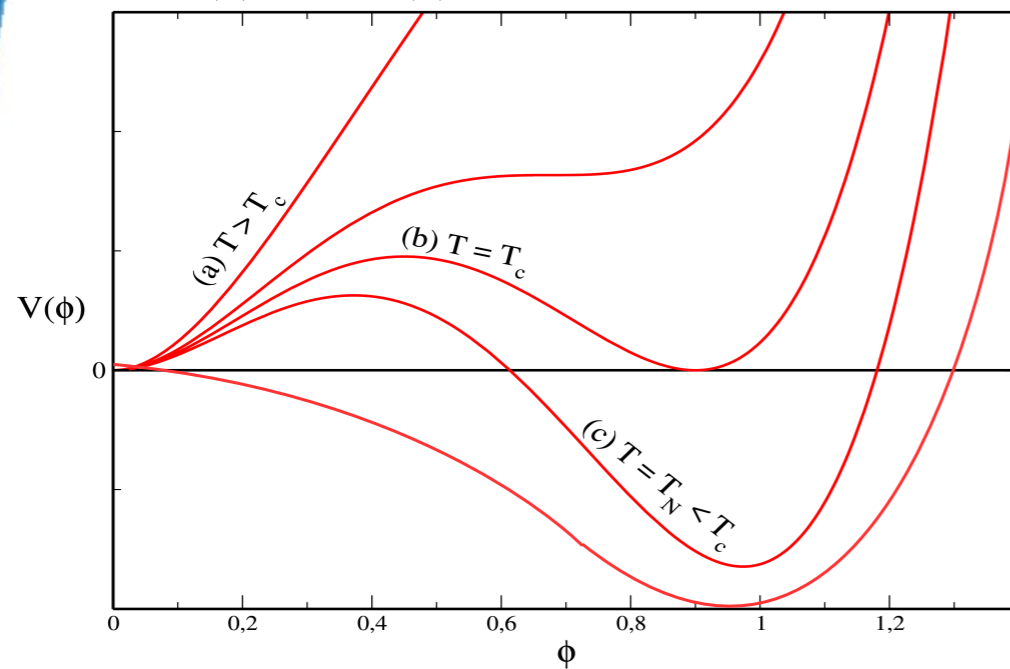
Loop  
corrections

Thermal  
corrections

$T \gg v$   
 $\kappa T^2 h^2$   
EW Symmetry  
Restoration

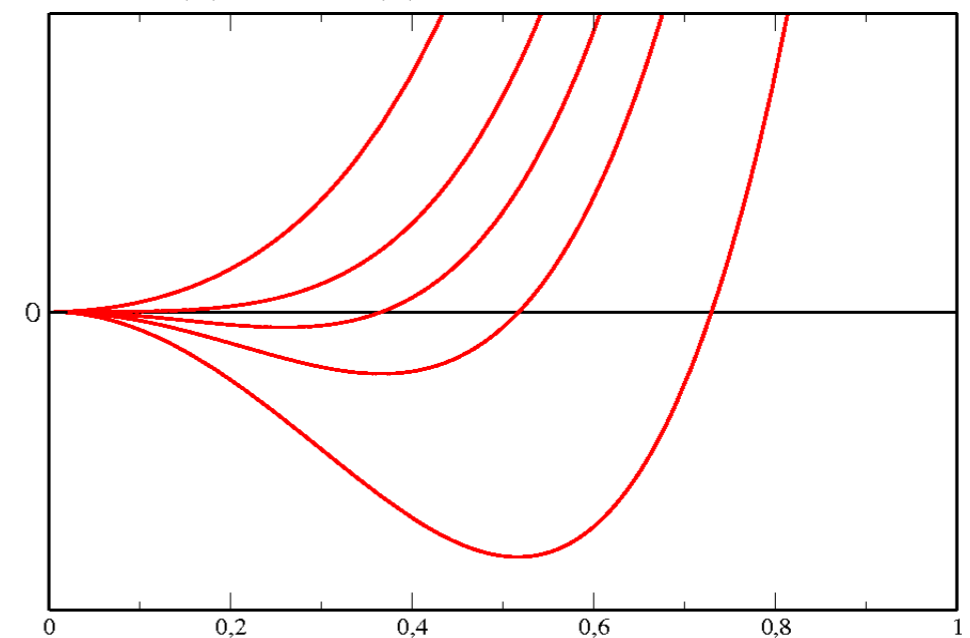
1<sup>st</sup> Order:

$\langle h \rangle = 0 \rightarrow \langle h \rangle = h(T)$  Discontinuous



2<sup>nd</sup> Order:

$\langle h \rangle = 0 \rightarrow \langle h \rangle = h(T)$  Continuous



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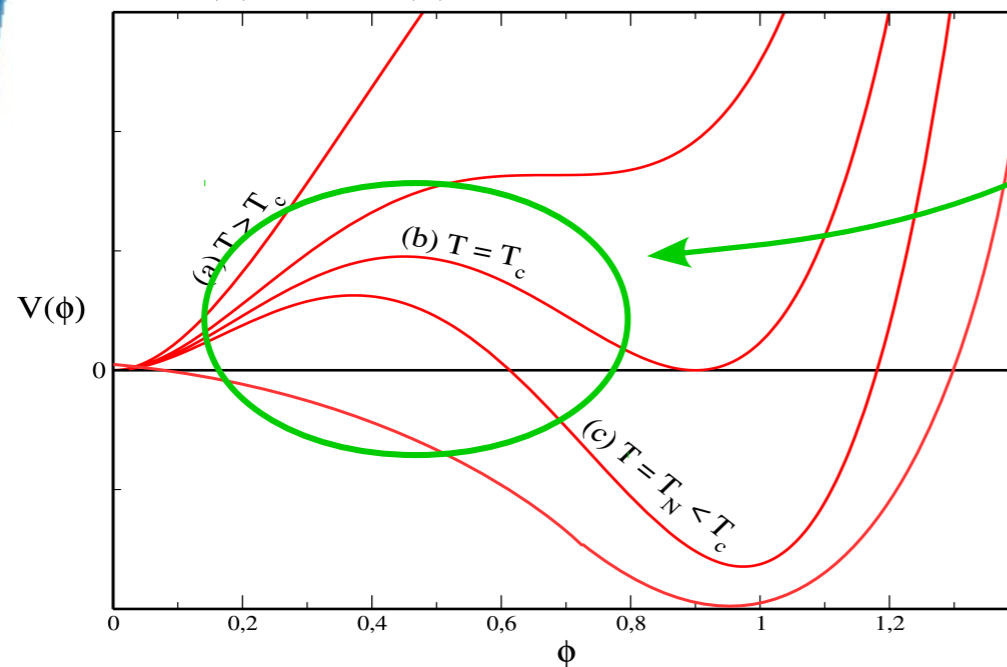
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$\Rightarrow$  **Thermal Effects** (In the SM, W, Z gauge bosons  $\rightarrow$  not sufficient)  
Add New BOSONS to generate a thermal barrier

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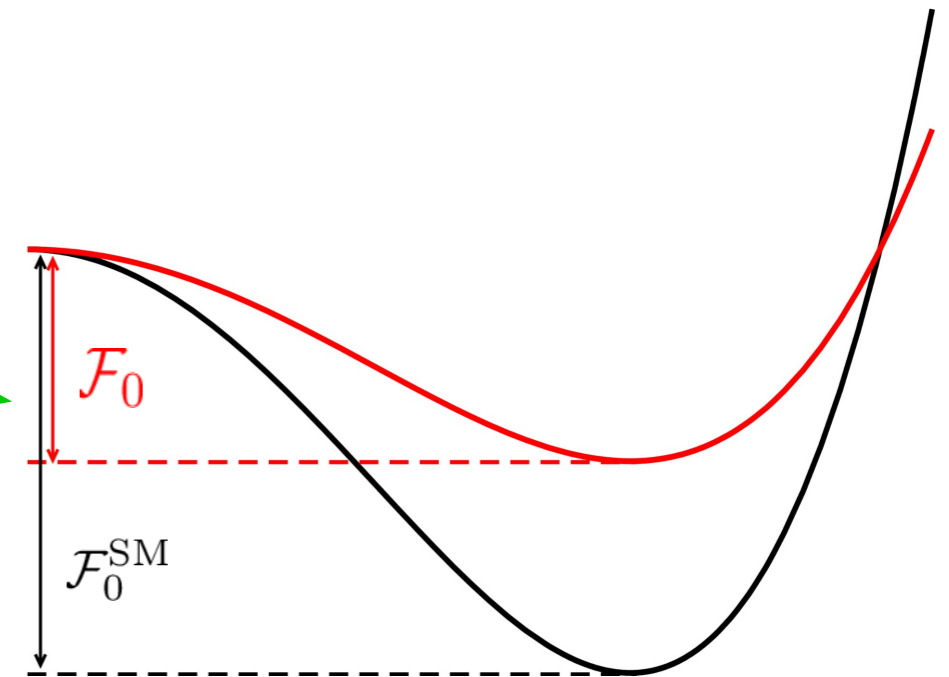
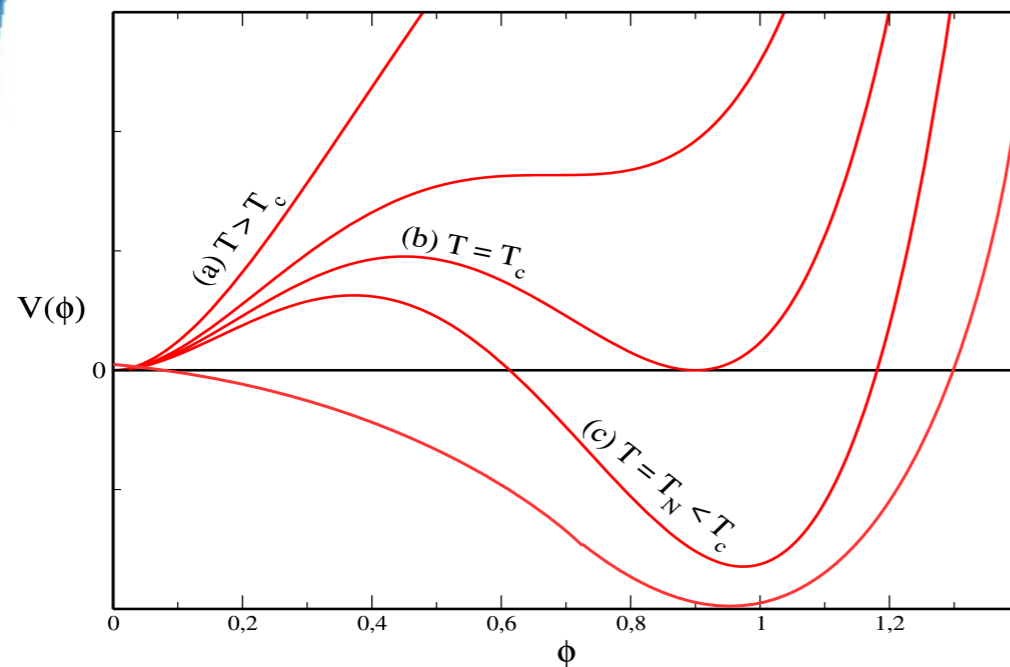
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Tree-level potential

Loop corrections

Thermal corrections

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 Add Particles whose loops reduce vacua energy difference.

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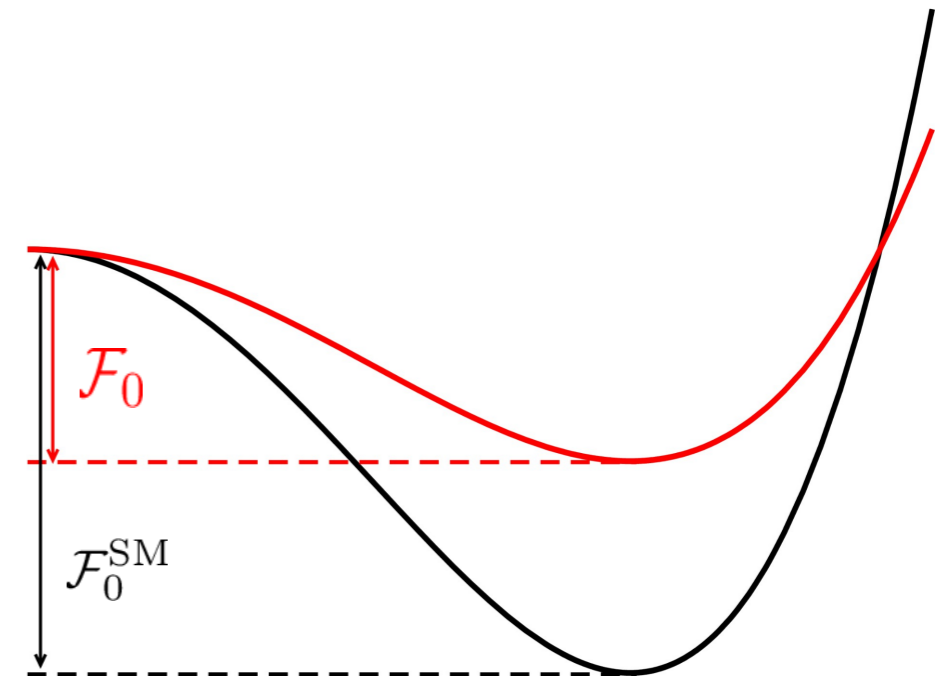
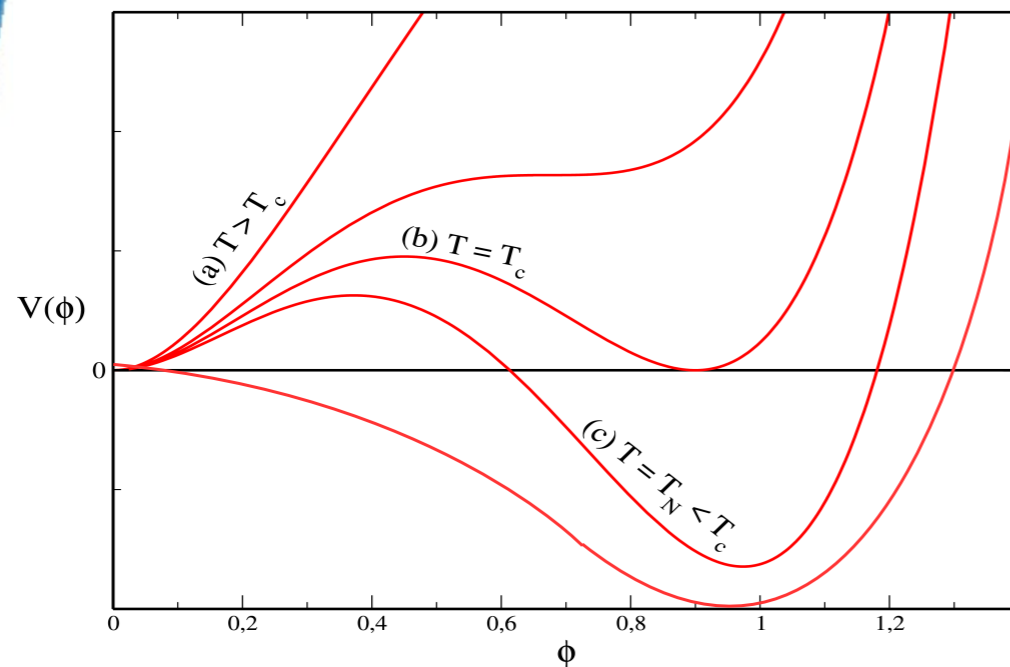
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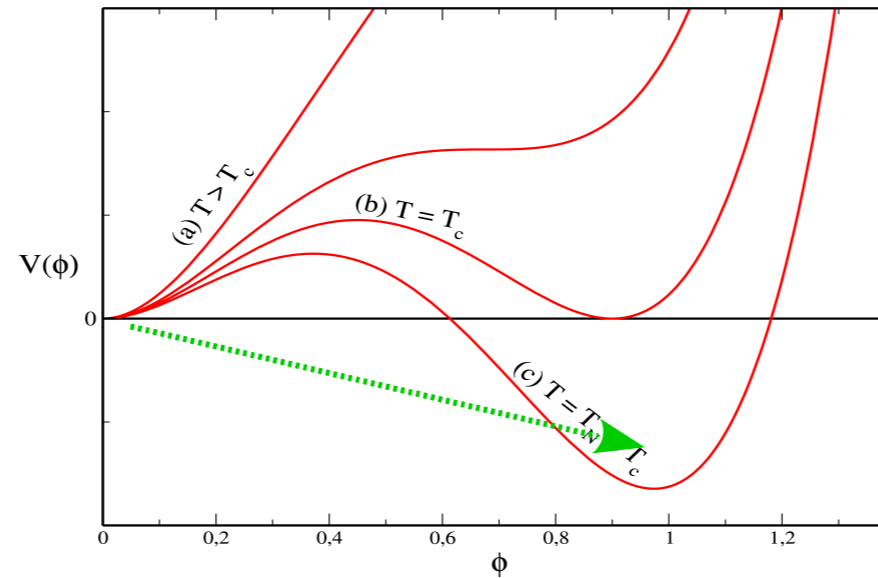
$\Rightarrow$  **Loop Effects**

Add Particles whose loops reduce vacua energy difference.

$\Rightarrow$  **Tree-level Effects**

Add scalars that modify the tree-level potential

# EW Phase Transition



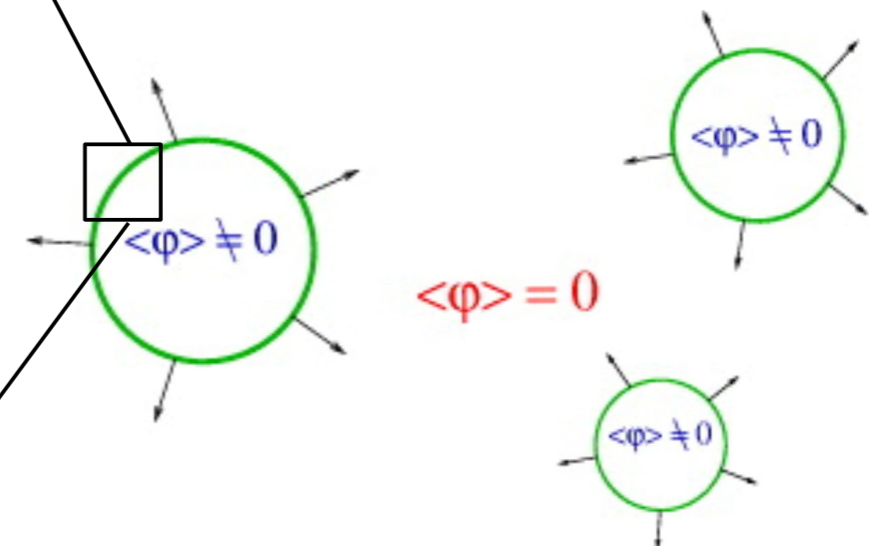
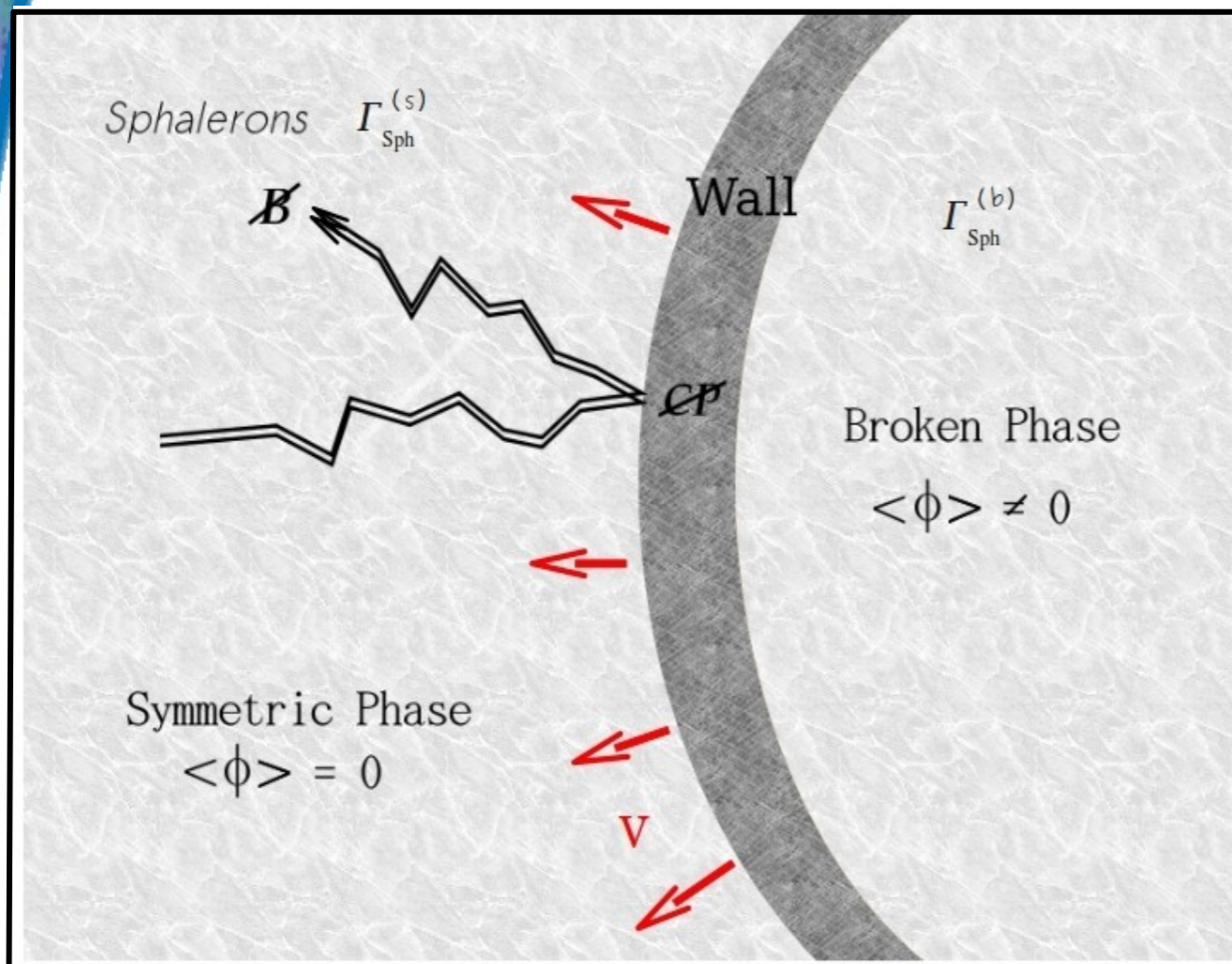
$$\Gamma_{\text{Sph}}^{(s)} \sim (\alpha_W T)^4$$

$$\Gamma_{\text{Sph}}^{(b)} \sim T^4 e^{-E_{\text{Sph}}/T}$$

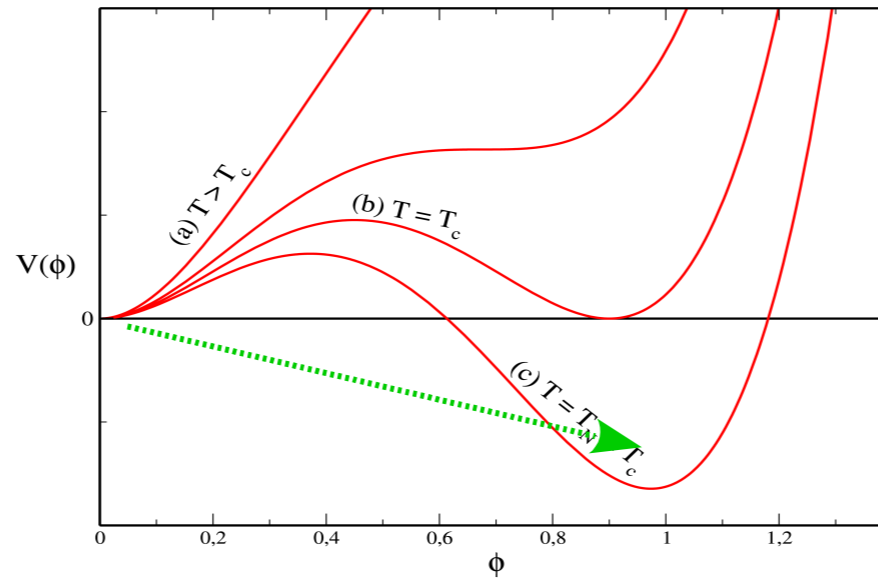
$$E_{\text{Sph}} = \kappa \frac{4\pi \langle \phi \rangle}{g}$$

Out of Equilibrium  
(SPHALERON SHUT-OFF)

$$\langle \phi \rangle / T > 1$$



# EW Phase Transition



Issues with Gauge Dependence

Patel, Ramsey-Musolf, JHEP **1107** (2011) 029

Garny, Konstandin, JHEP **1207** (2012) 189

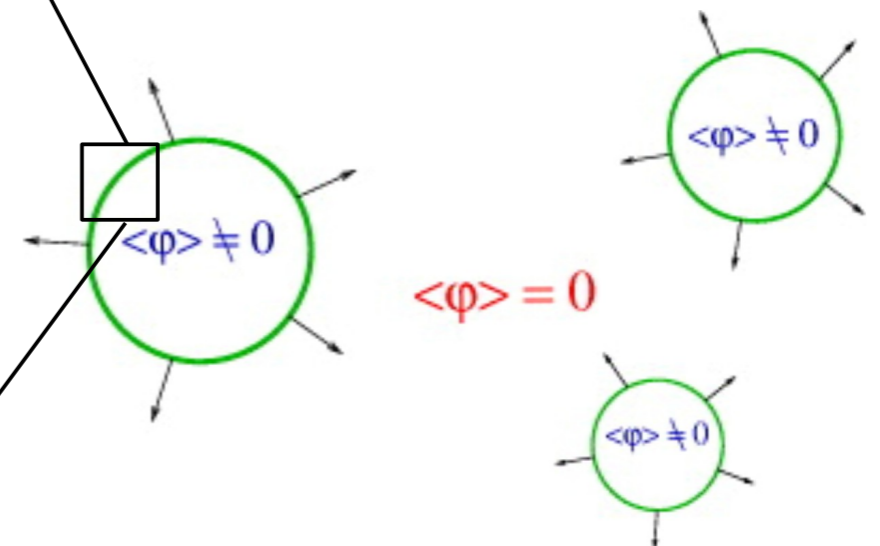
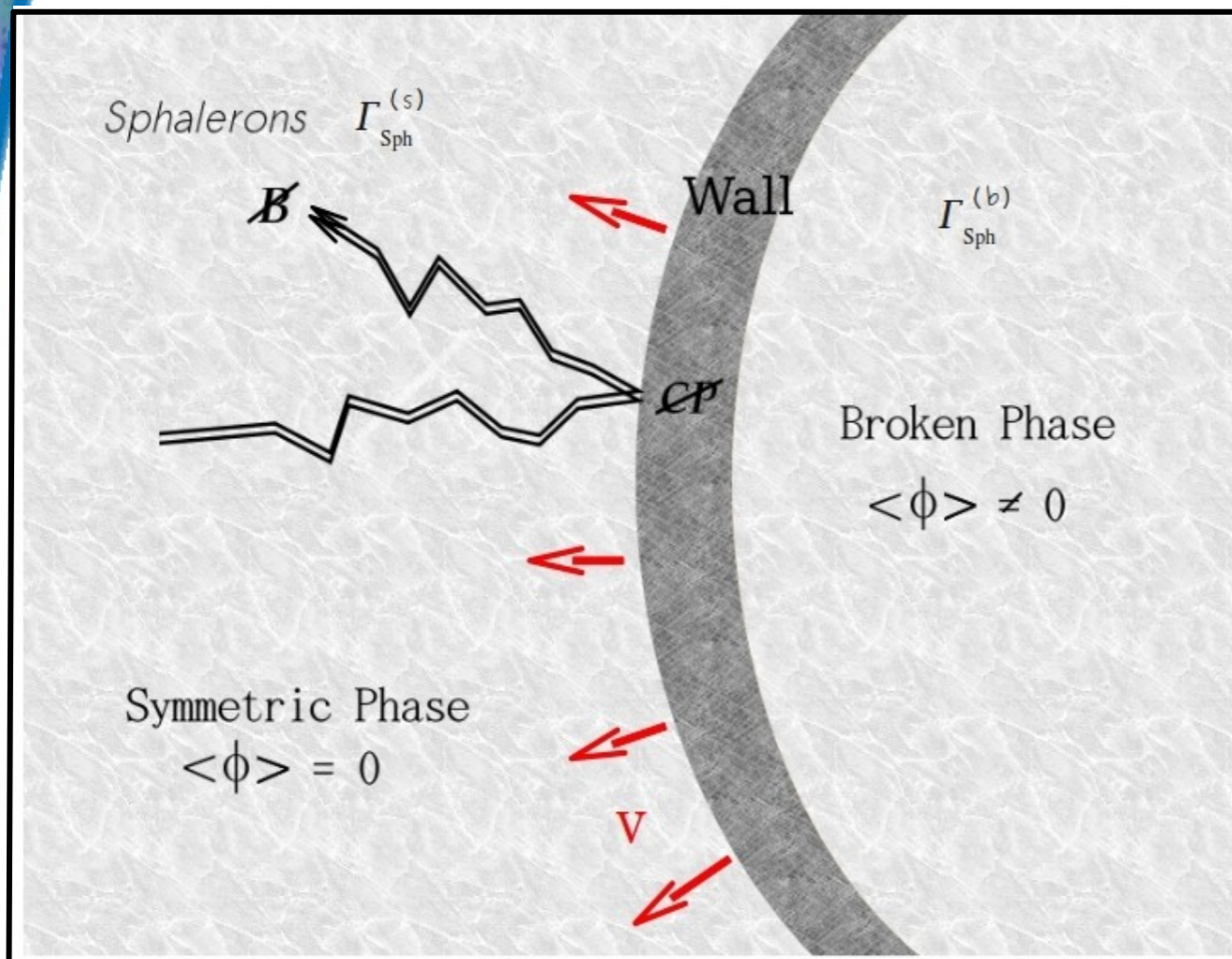
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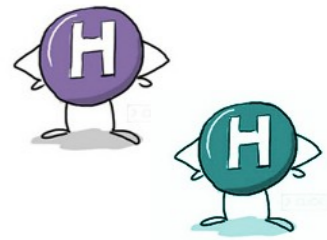
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# Baryogenesis with a Second Higgs



$$\begin{aligned}
 V(H_1, H_2) &= \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\
 &+ \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\
 &+ \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[ (H_1^\dagger H_2)^2 + \text{h.c.} \right]
 \end{aligned}
 \quad H_j = \begin{pmatrix} \phi_j^+ \\ \frac{v_j + h_j + i \eta_j}{\sqrt{2}} \end{pmatrix}$$

$$H^\pm = -s_\beta \phi_1^\pm + c_\beta \phi_2^\pm$$

$$h = -s_\alpha h_1 + c_\alpha h_2$$

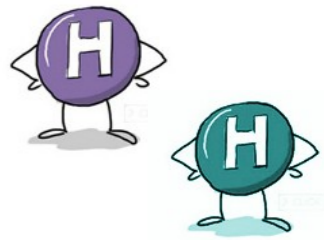
$$A_0 = -s_\beta \eta_1 + c_\beta \eta_2$$

$$H_0 = -c_\alpha h_1 - s_\alpha h_2$$

ALL Needed Ingredients for EW Baryogenesis: **Out-of-Equilibrium + CPV**

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## EW Phase Transition



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 H_0 &= -c_\alpha h_1 - s_\alpha h_2
 \end{aligned}$$

ALL Needed Ingredients for EW Baryogenesis: **Out-of-Equilibrium + CPV**

For simplicity, let's not consider CPV yet

BSM Parameters  $m_{H_0}$   $m_{A_0}$   $m_{H^\pm}$   $\tan\beta$   $\cos(\beta - \alpha)$   $\mu^2$

EWPO:  $m_{H^\pm} \simeq m_{H_0}$  OR  $m_{H^\pm} \simeq m_{A_0}$

Choice of  $H_j$  Couplings to Fermions not Relevant for EW Phase Transition

Type I, Type II ...



# Baryogenesis with a Second Higgs

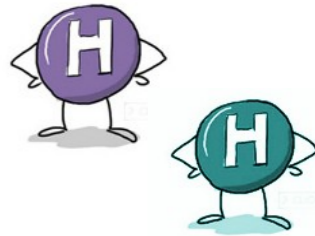
EW Phase Transition



Thermal + Loop + Tree-level Effects

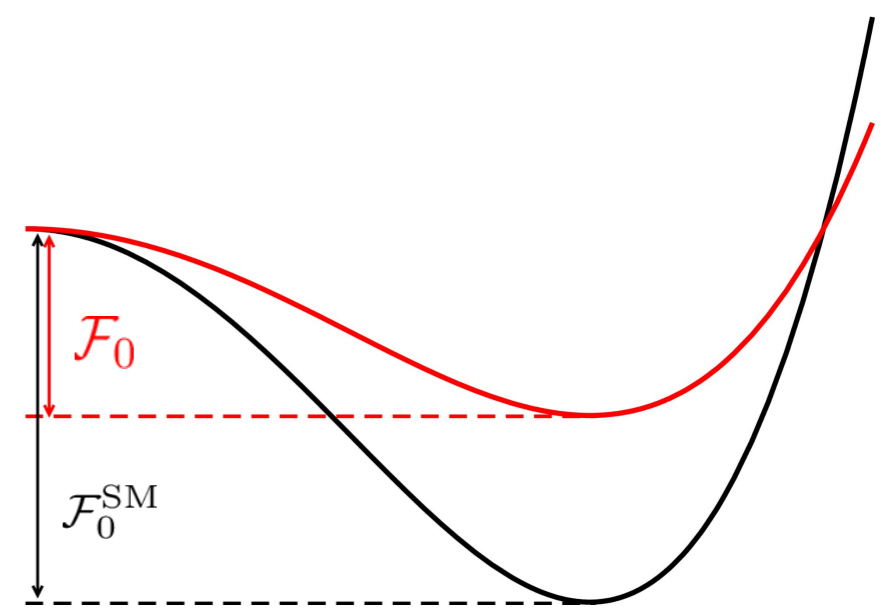
# Baryogenesis with a Second Higgs

## EW Phase Transition



Thermal + Loop + Tree-level Effects

$$\mathcal{F}_0 = V_0(v) + V_0^{\text{loop}}(v) - V_0(0) - V_0^{\text{loop}}(0)$$



[Journal of High Energy Physics](#)  
June 2016, 2016:5

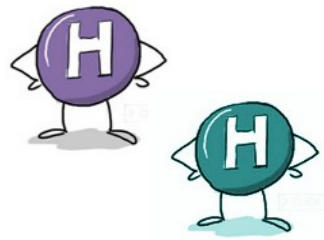
### Does zero temperature decide on the nature of the electroweak phase transition?

Christopher P.D. Harman , Stephan J. Huber

Huang, Kang, Shu, Wu, Yang, *Phys.Rev. D* **91** (2015) 025006  
Harman, Huber, *JHEP* **1606** (2016) 005

# Baryogenesis with a Second Higgs

## EW Phase Transition

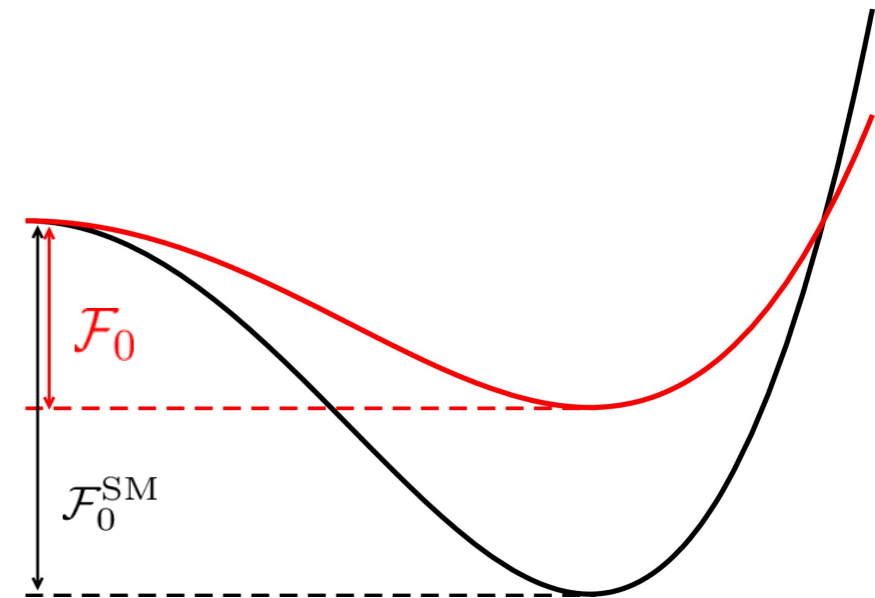


Thermal + Loop + Tree-level Effects

$$\mathcal{F}_0 = V_0(v) + V_0^{\text{loop}}(v) - V_0(0) - V_0^{\text{loop}}(0)$$



Huang, Kang, Shu, Wu, Yang, *Phys.Rev. D* **91** (2015) 025006  
 Harman, Huber, *JHEP* **1606** (2016) 005



$$[\mathcal{F}_0 - \mathcal{F}_0^{\text{SM}}]_{\text{tree}} = -\frac{v^2}{8} c_{\beta-\alpha}^2 (m_{H_0}^2 - m_h^2) < 0 \quad (\text{Assuming } m_h = 125 \text{ GeV})$$

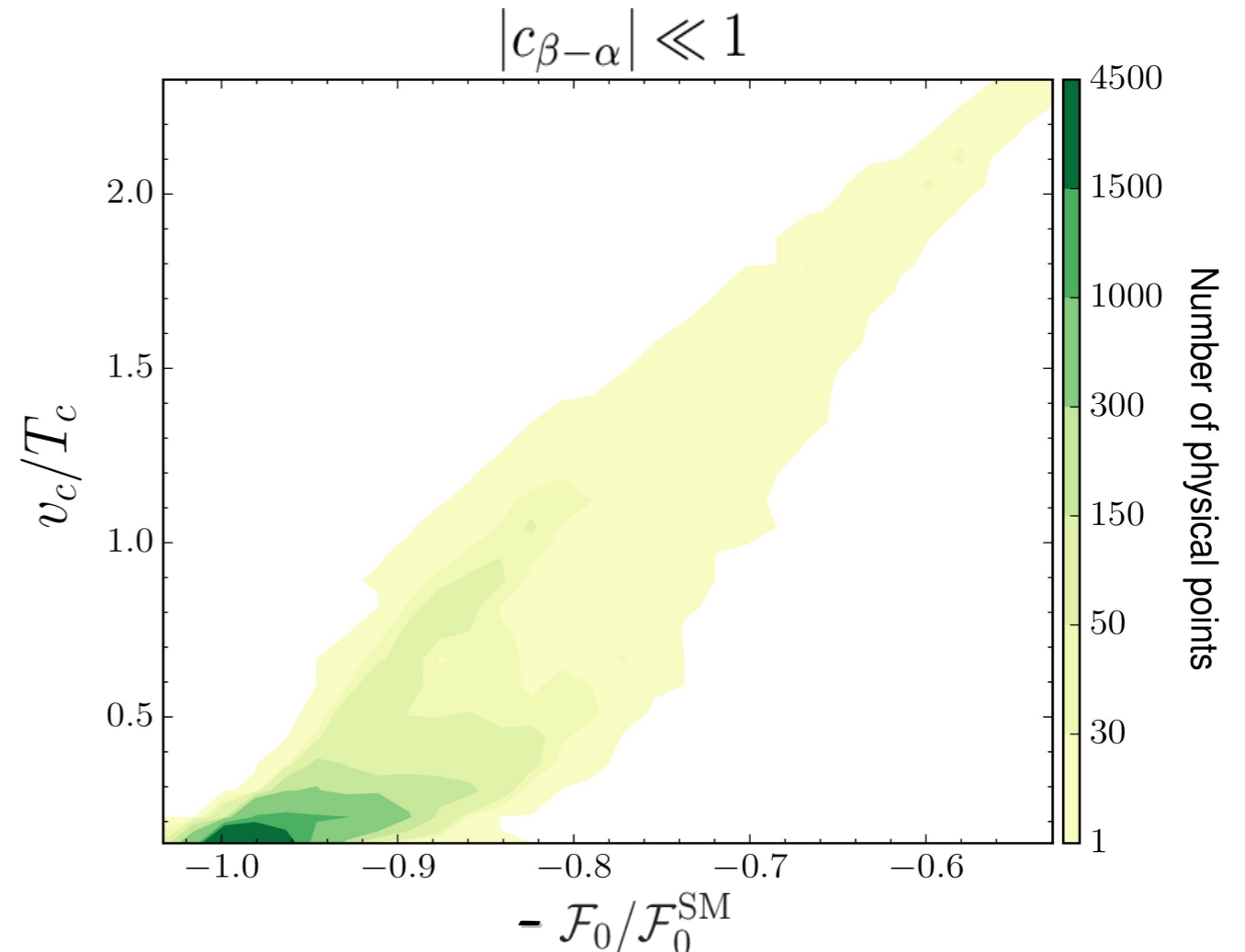
# Baryogenesis with a Second Higgs



## EW Phase Transition

LET'S LOOK (FIRST) AT THE ALIGNMENT LIMIT:

$$\mathcal{F}_0 - \mathcal{F}_0^{\text{SM}} = \frac{1}{64\pi^2} \left[ (2M^2 - m_h^2)^2 \left( \frac{3}{2} + \frac{1}{2} \log \left[ \frac{4m_{A_0} m_{H_0} m_{H^\pm}^2}{(2M^2 - m_h^2)^2} \right] \right) \right. \\ \left. + \frac{1}{2} (m_{A_0}^4 + m_{H_0}^4 + 2m_{H^\pm}^4) - (2M^2 - m_h^2) (m_{A_0}^2 + m_{H_0}^2 + 2m_{H^\pm}^2) \right] \quad M^2 = \mu^2 / (s_\beta c_\beta)$$



# Baryogenesis with a Second Higgs



## EW Phase Transition

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$$M^2 = \mu^2 / (s_\beta c_\beta)$$

Combine with (tree-level)

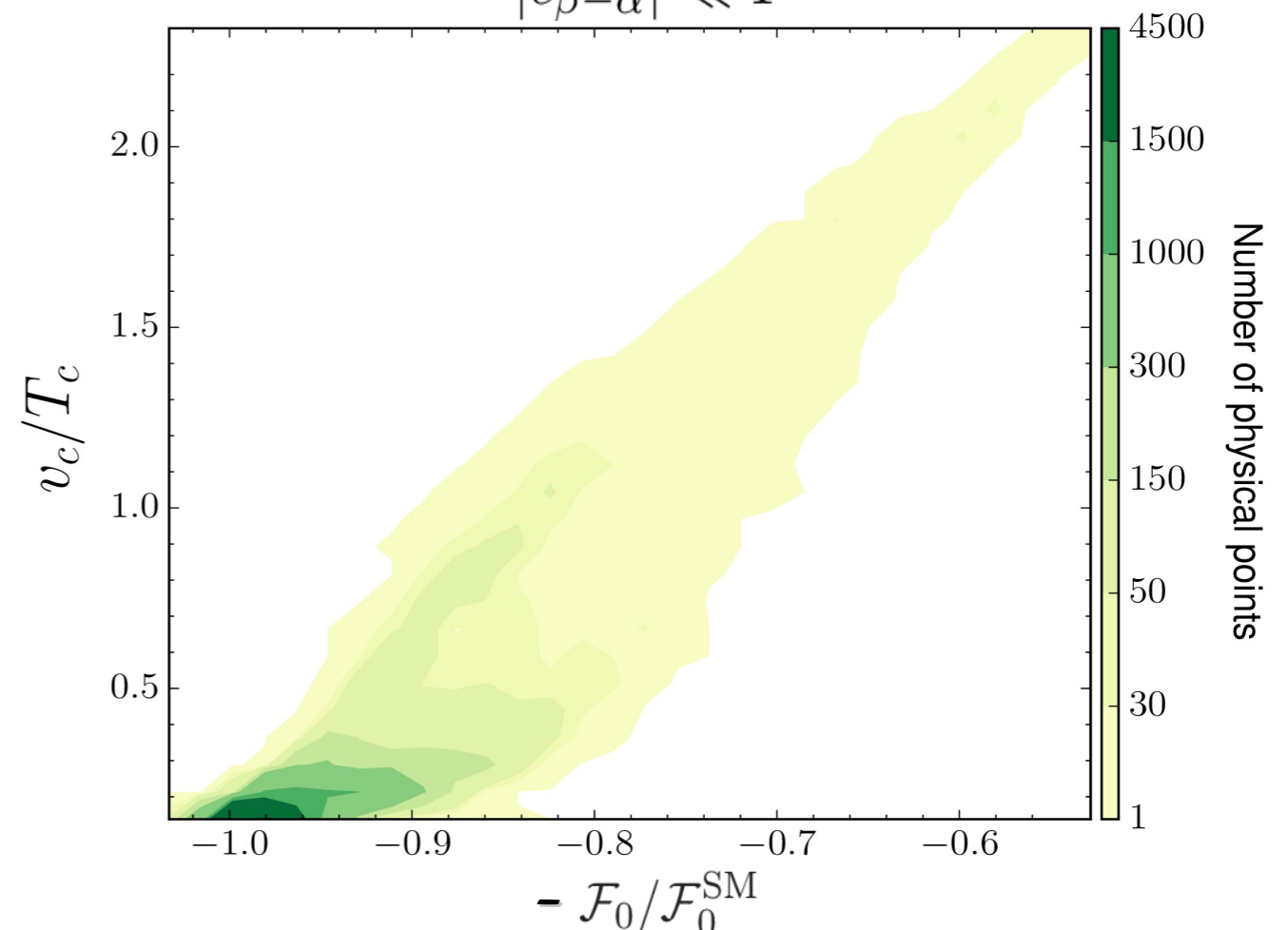
Vacuum Stability and Unitarity

[Gunion, Haber, Phys.Rev. D67 \(2003\) 075019](#)

[Ginzburg, Ivanov, Phys.Rev. D72 \(2005\) 115010](#)

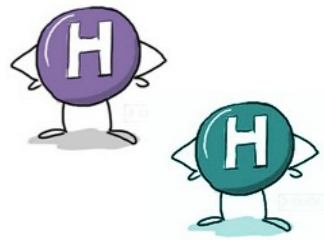
[Barroso, Ferreira, Ivanov, Santos, JHEP 1306 \(2013\) 045](#)

$$|c_{\beta-\alpha}| \ll 1$$



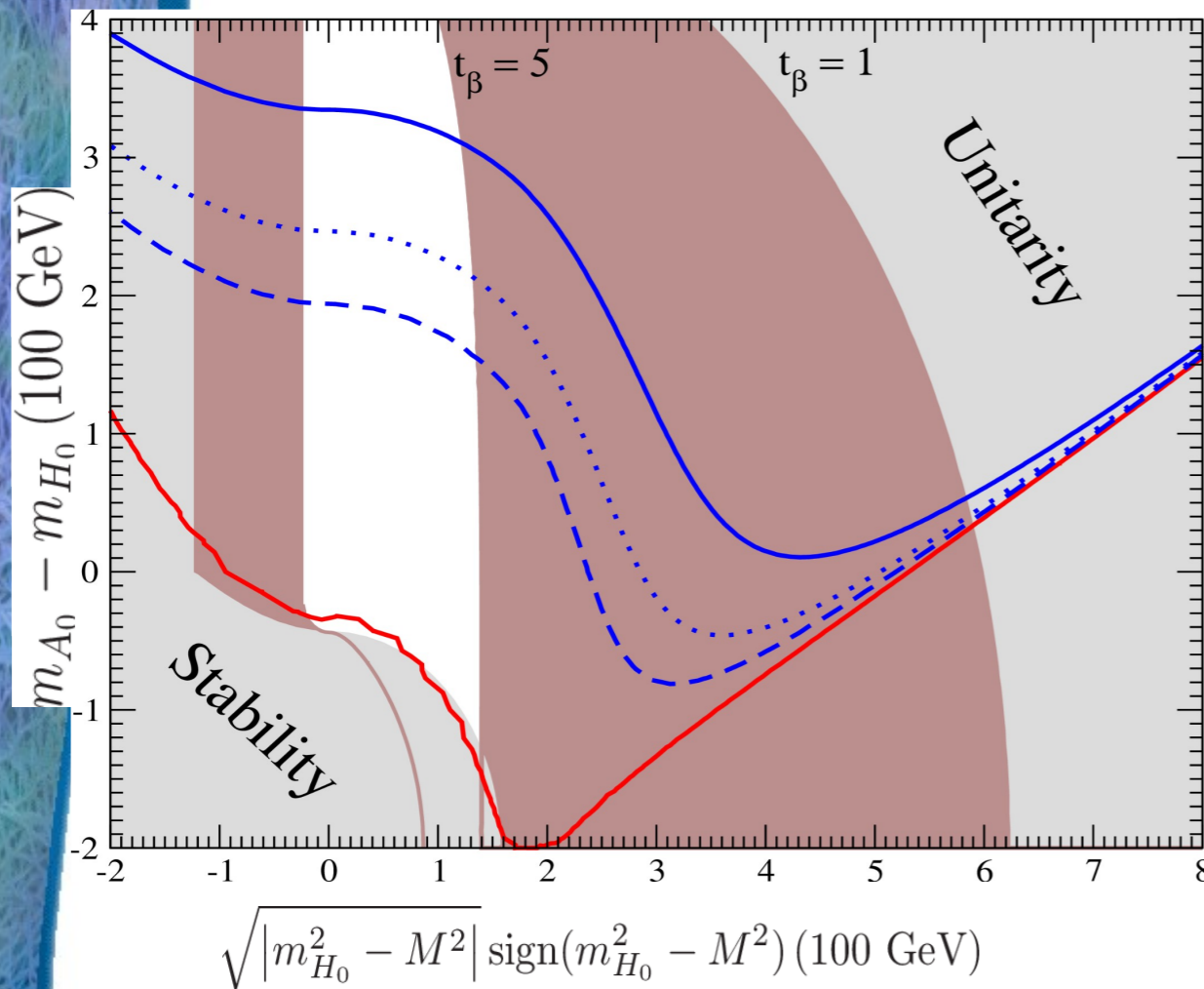
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## EW Phase Transition

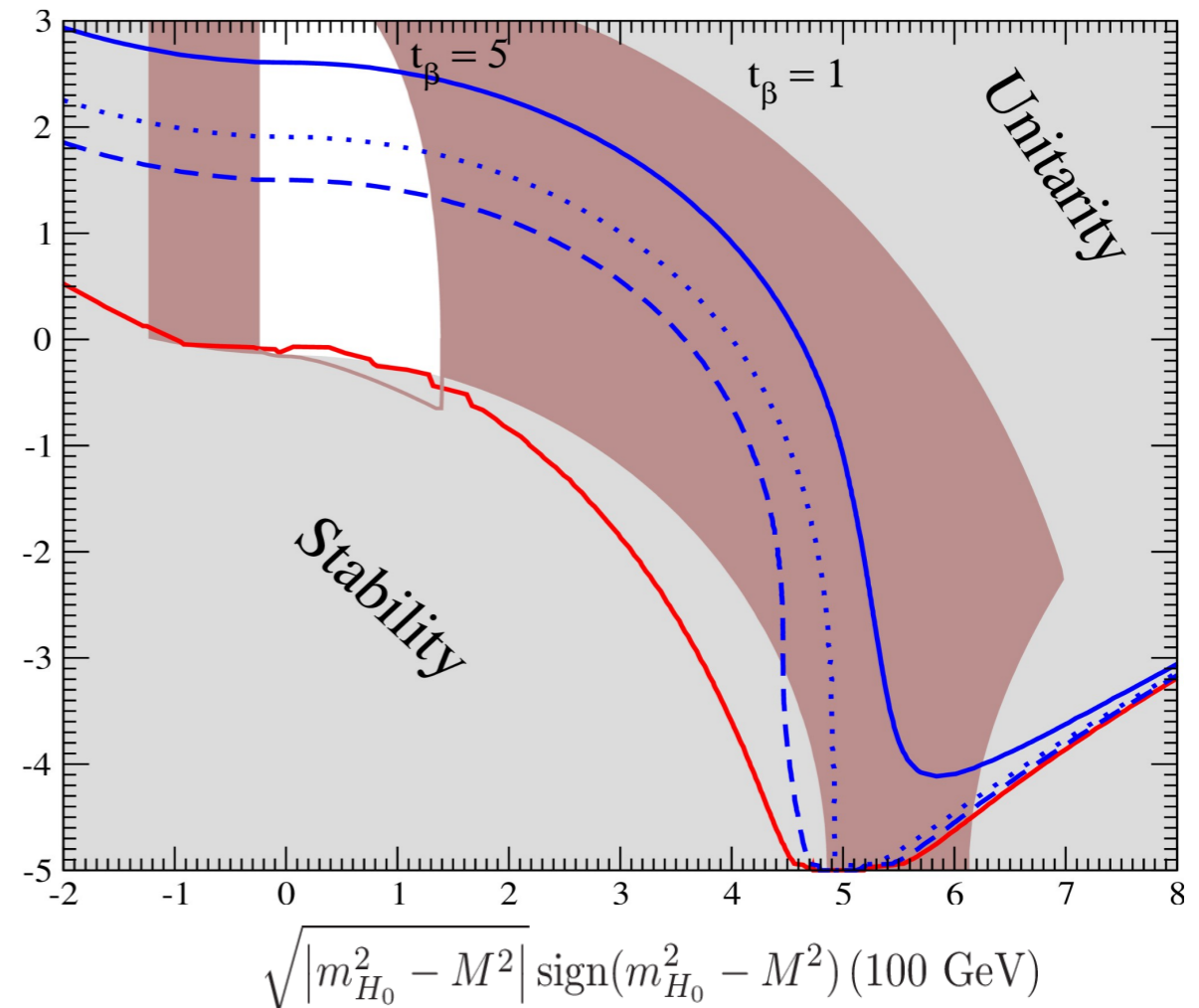


$$m_{H^\pm} \simeq m_{A_0}$$

$m_{H_0} = 200 \text{ GeV}$



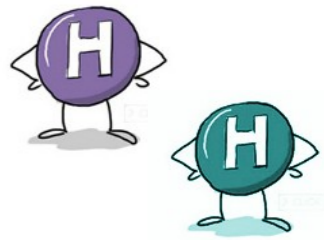
$m_{H_0} = 500 \text{ GeV}$



—  $\mathcal{F}_0/\mathcal{F}_0^{\text{SM}} = 1$    
 - - -  $\mathcal{F}_0/\mathcal{F}_0^{\text{SM}} = 0.8$    
 · · ·  $\mathcal{F}_0/\mathcal{F}_0^{\text{SM}} = 0.6$    
 —  $\mathcal{F}_0/\mathcal{F}_0^{\text{SM}} = 0$

Dorsch, Huber, Mimasu, JMN, arXiv:1609.xxxxx

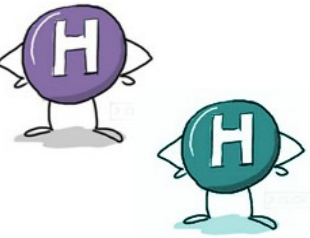
# Baryogenesis with a Second Higgs



EW Phase Transition

- ⇒ EW Phase Transition Favours 2HDM Alignment or  $m_{H_0} \sim m_h$
- ⇒ EW Phase Transition Favours  $m_{A_0} - m_{H_0} \sim v$  ( $> m_z$ )

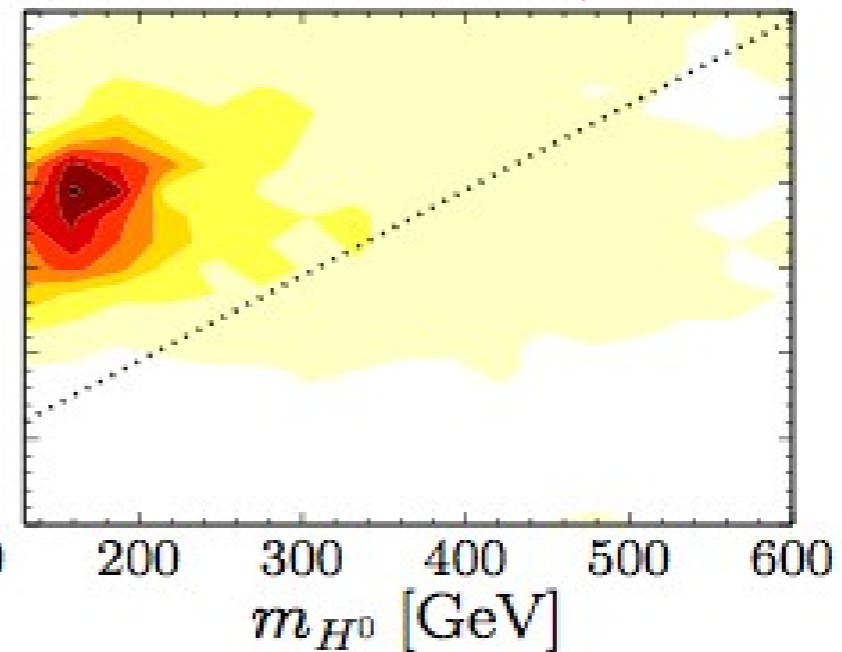
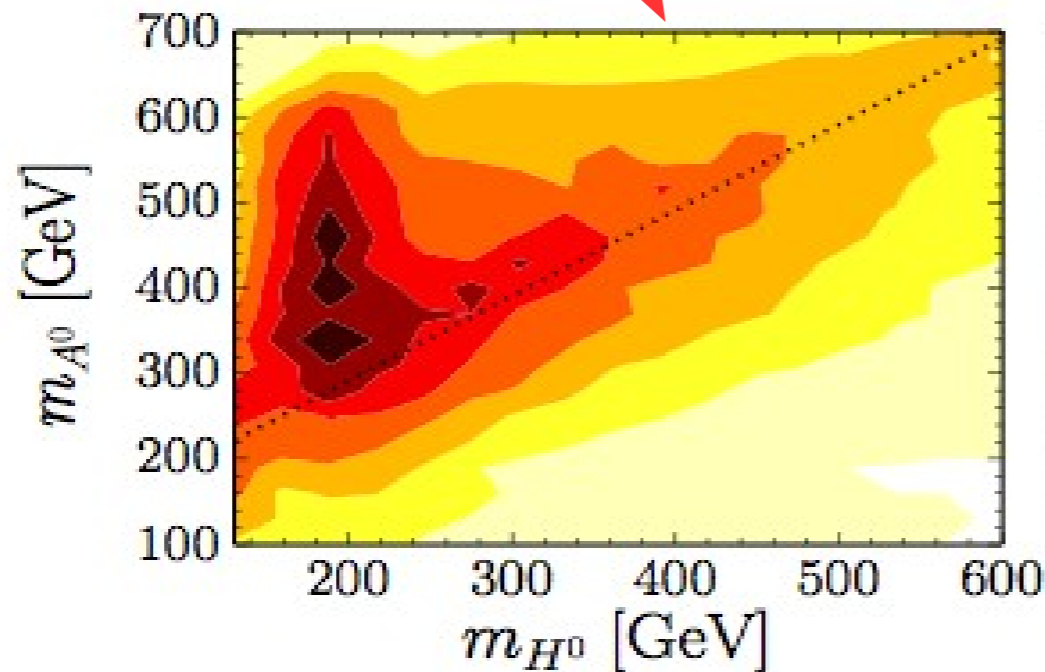
# Baryogenesis with a Second Higgs



## EW Phase Transition

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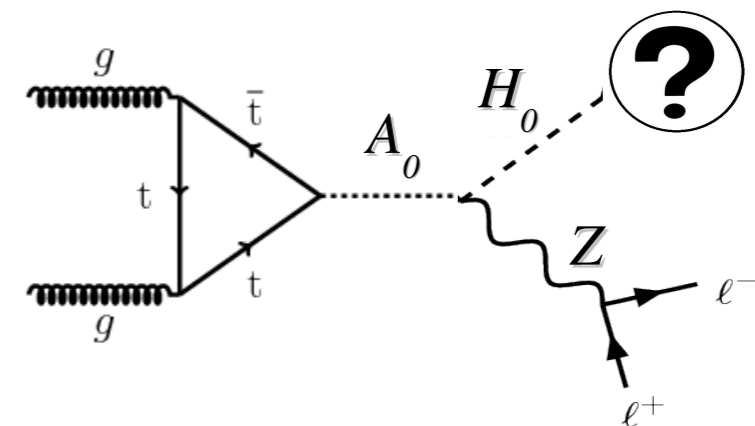
Physically Allowed  $\supset$  EW Phase Transition



EW Phase Transition SIGNATURE

$$A_0 \rightarrow H_0 Z$$

Dorsch, Huber, Mimasu, JMN, Phys. Rev. Lett. **113** (2014) 211802

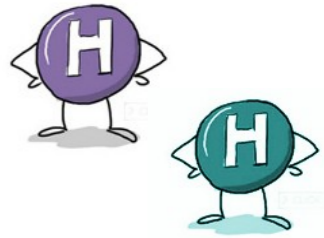


$\bar{b}b$   
or  
 $WW, ZZ$



# Baryogenesis with a Second Higgs @ LHC

EW Phase Transition



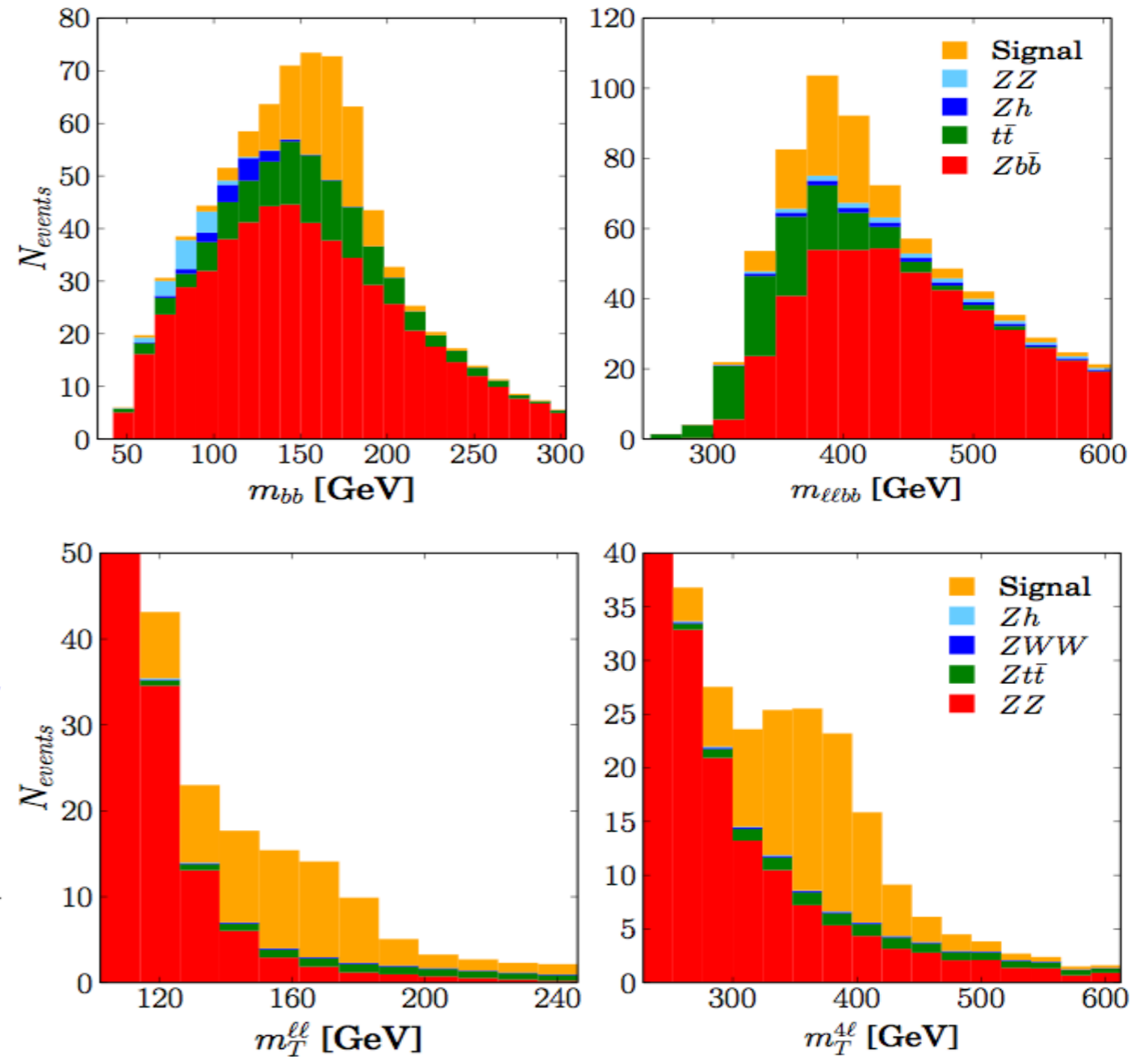
$$A_0 \rightarrow H_0 Z \rightarrow \bar{b}b \ell\ell$$

$$A_0 \rightarrow H_0 Z \rightarrow W^+W^- \ell\ell$$

$$(m_T^{\ell\ell})^2 = (\sqrt{p_{T,\ell\ell}^2 + m_{\ell\ell}^2} + \cancel{p}_T)^2 - (\vec{p}_{T,\ell\ell} + \vec{\cancel{p}}_T)^2$$

$$m_T^{4\ell} = \sqrt{p_{T,\ell'\ell'}^2 + m_{\ell'\ell'}^2} + \sqrt{p_{T,\ell\ell}^2 + (m_T^{\ell\ell})^2}$$

LHC Run II



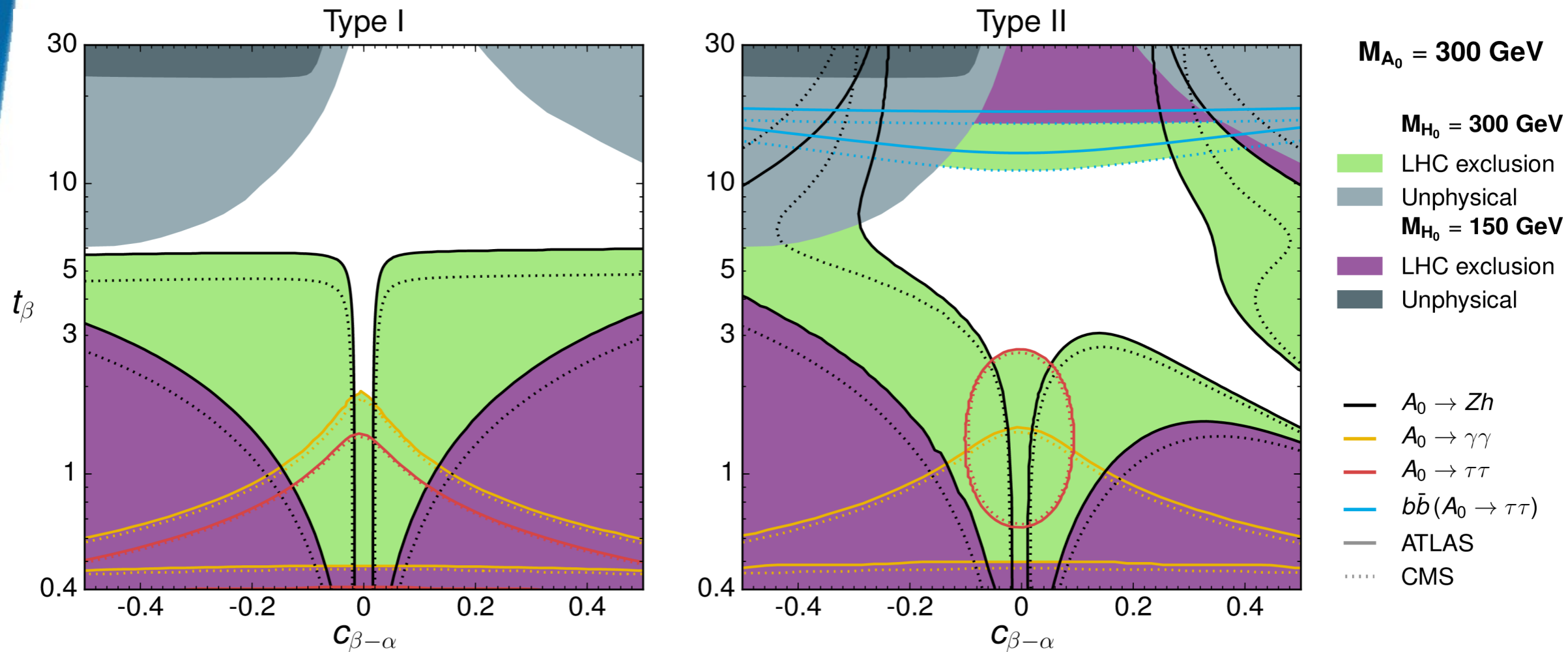
# Baryogenesis with a Second Higgs @ LHC

## EW Phase Transition



$A_0 \rightarrow H_0 Z$  channel Open

BR to other decay channels drastically reduced



Dorsch, Huber, Mimasu, JMN, Phys. Rev. D93 (2016) 115033

# CMS-PAS-HIG-15-001

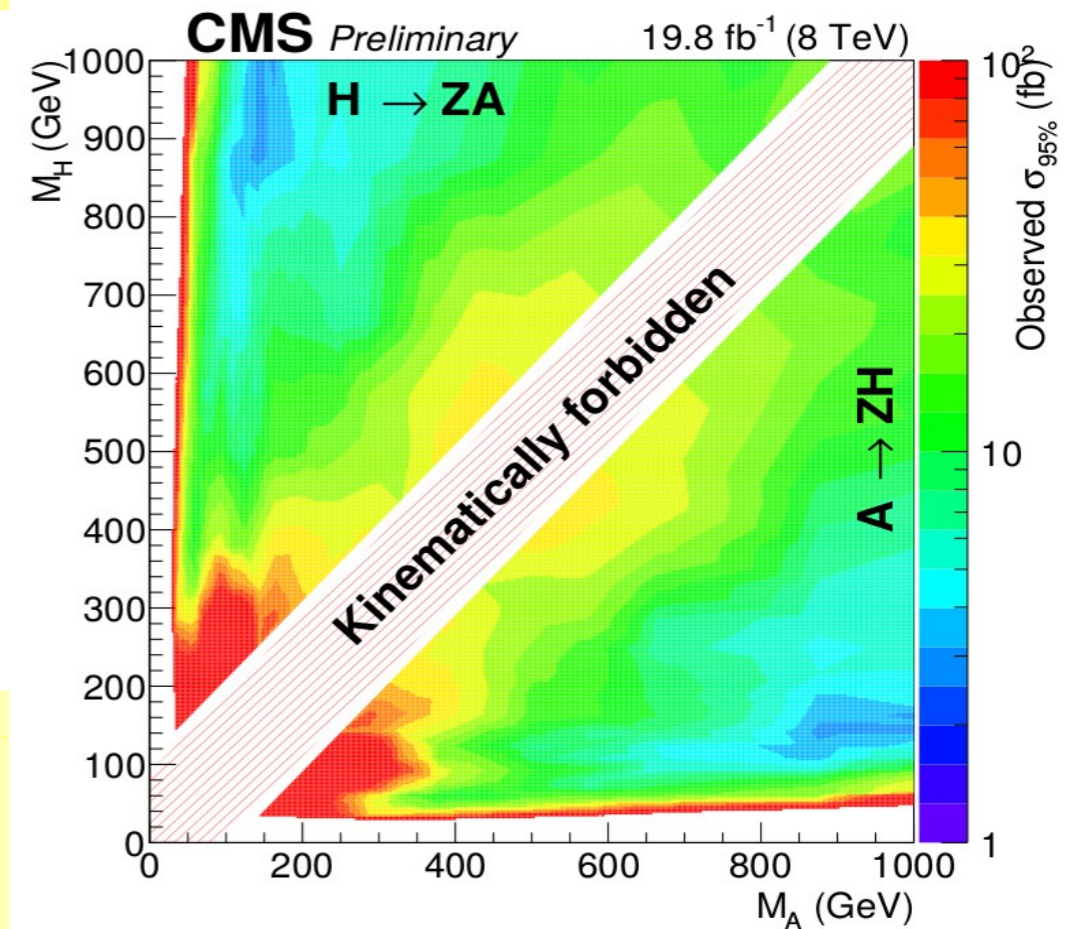
*Phys. Lett. B759 (2016) 369 (ArXiv:1603.02991)*

Search for H/A decaying into Z and A/H, with  $Z \rightarrow \ell\ell$  and  
 $A/H \rightarrow bb$  or  $A/H \rightarrow \tau\tau$

The CMS Collaboration

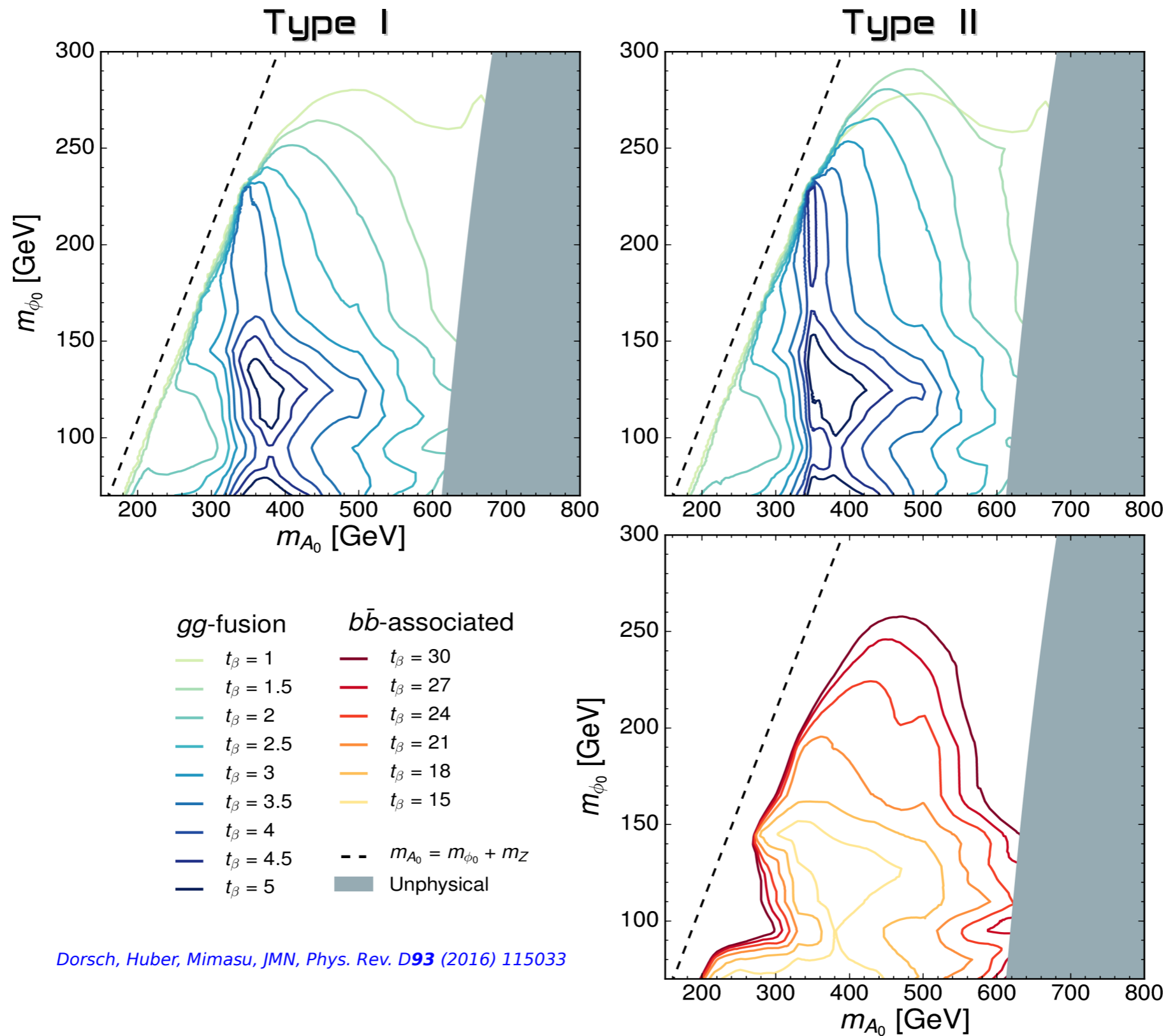
One important motivation for 2HDMs is that these models provide a way to explain the asymmetry between matter and anti-matter observed in the Universe [4, 5]. Another important motivation is Supersymmetry [6], which is a theory that falls in the broad class of 2HDMs. Axion models [7], which would explain how the strong interaction does not violate the CP symmetry, would give rise to an effective low-energy theory with two Higgs doublets. Finally, it has also been recently noted [8] that certain realizations of 2HDMs can accommodate the muon  $g - 2$  anomaly [9] without violating the present theoretical and experimental constraints.

In the most general case 14 parameters are necessary to describe the scalar sector in a 2HDM. However, only 6 free parameters remain once the so-called  $Z_2$  symmetry is imposed to suppress flavor changing neutral currents, in agreement with experimental observations, and the values of the mass of the recently discovered Higgs boson (125 GeV) and the electroweak vacuum expectation value (246 GeV) are assumed. The compatibility of a 125 GeV SM-like Higgs boson with 2HDMs is possible in the so-called alignment limit. In such a limit, one of the CP-even scalars,  $h$  or  $H$ , is identified with the 125 GeV Higgs boson and the condition  $\cos(\beta - \alpha) \approx 0$  or  $\sin(\beta - \alpha) \approx 0$  is satisfied, where  $\tan \beta$  and  $\alpha$  are, respectively, the ratio of the vacuum expectation values, and the mixing angle of the two Higgs doublets. A recent theoretical study [5] has shown that, in this limit, a large mass splitting ( $> 100$  GeV) between the A and H bosons would favor the electroweak phase transition that would be at the origin of the baryogenesis process in the early Universe, thus explaining the currently observed matter-antimatter asymmetry in the Universe. In such a scenario, the most frequent decay mode of the pseudoscalar A boson would be  $A \rightarrow ZH$ .



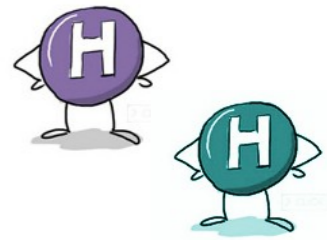
# Meaningful Constraints from LHC Run 1

(Assume 2HDM Alignment)



Dorsch, Huber, Mimasu, JMN, Phys. Rev. D93 (2016) 115033

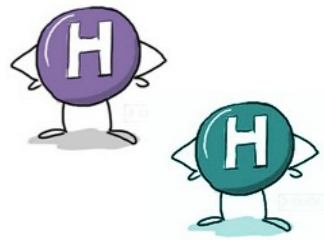
# Baryogenesis with a Second Higgs



Out-of-Equilibrium ✓  
(EW Phase Transition)

CPV: Baryogenesis vs EDMs

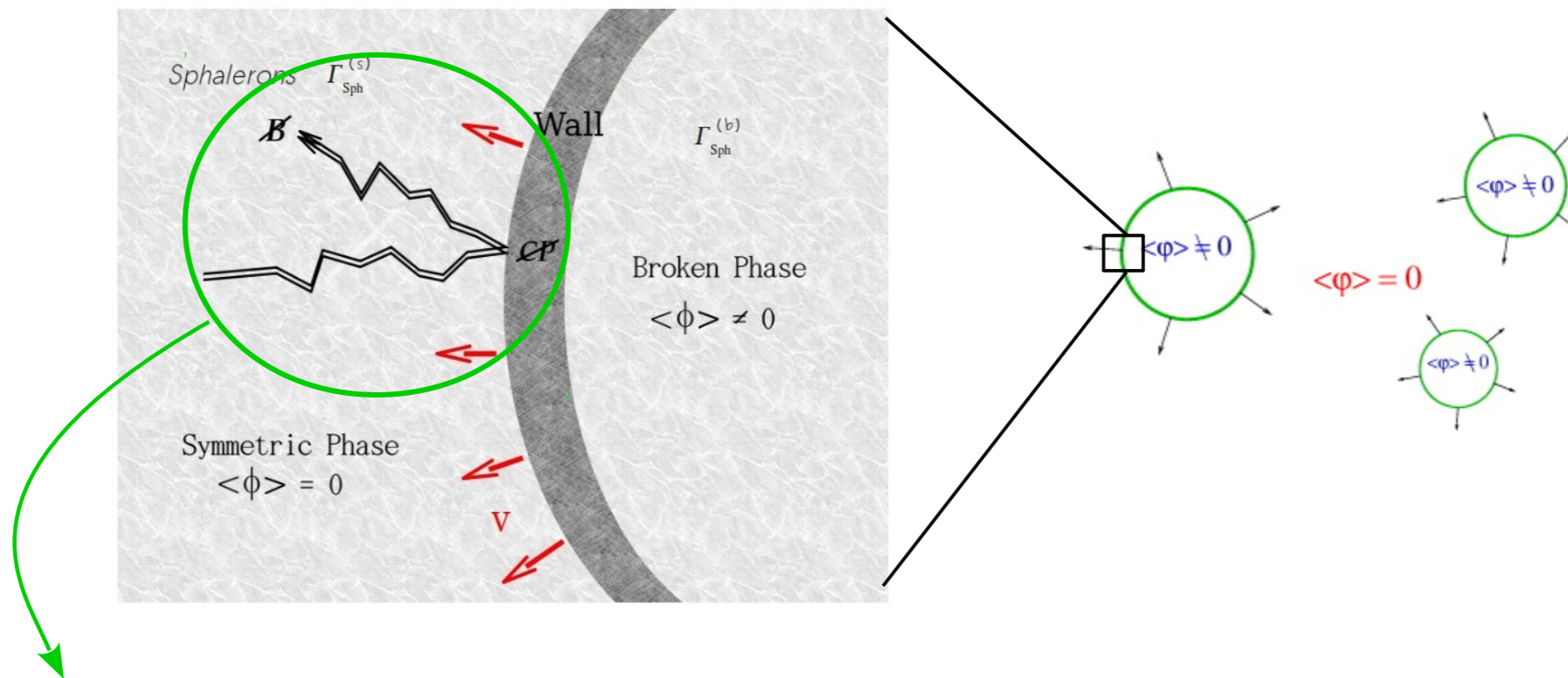
# Baryogenesis with a Second Higgs



Out-of-Equilibrium ✓  
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CPV: Baryogenesis vs EDMs

But before that... an Equally Important Issue

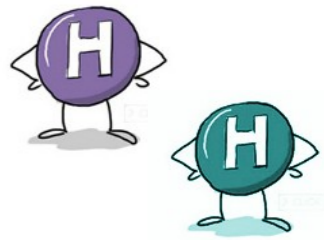


The Bubble Wall Velocity is Key for Successful Baryogenesis

$$\partial_\mu \partial^\mu \phi + \frac{\partial \mathcal{F}}{\partial \phi} - \mathcal{K}(\phi) = 0 \quad \mathcal{K}(\phi) = - \sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3 p}{(2\pi)^3 2E_i} \delta f_i(p)$$

$$F_\eta(v_w) = \int_{-\infty}^{\infty} dz \frac{d\phi}{dz} \sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3 p}{(2\pi)^3 2E_i} \delta f_i(p) \simeq \eta v_w$$

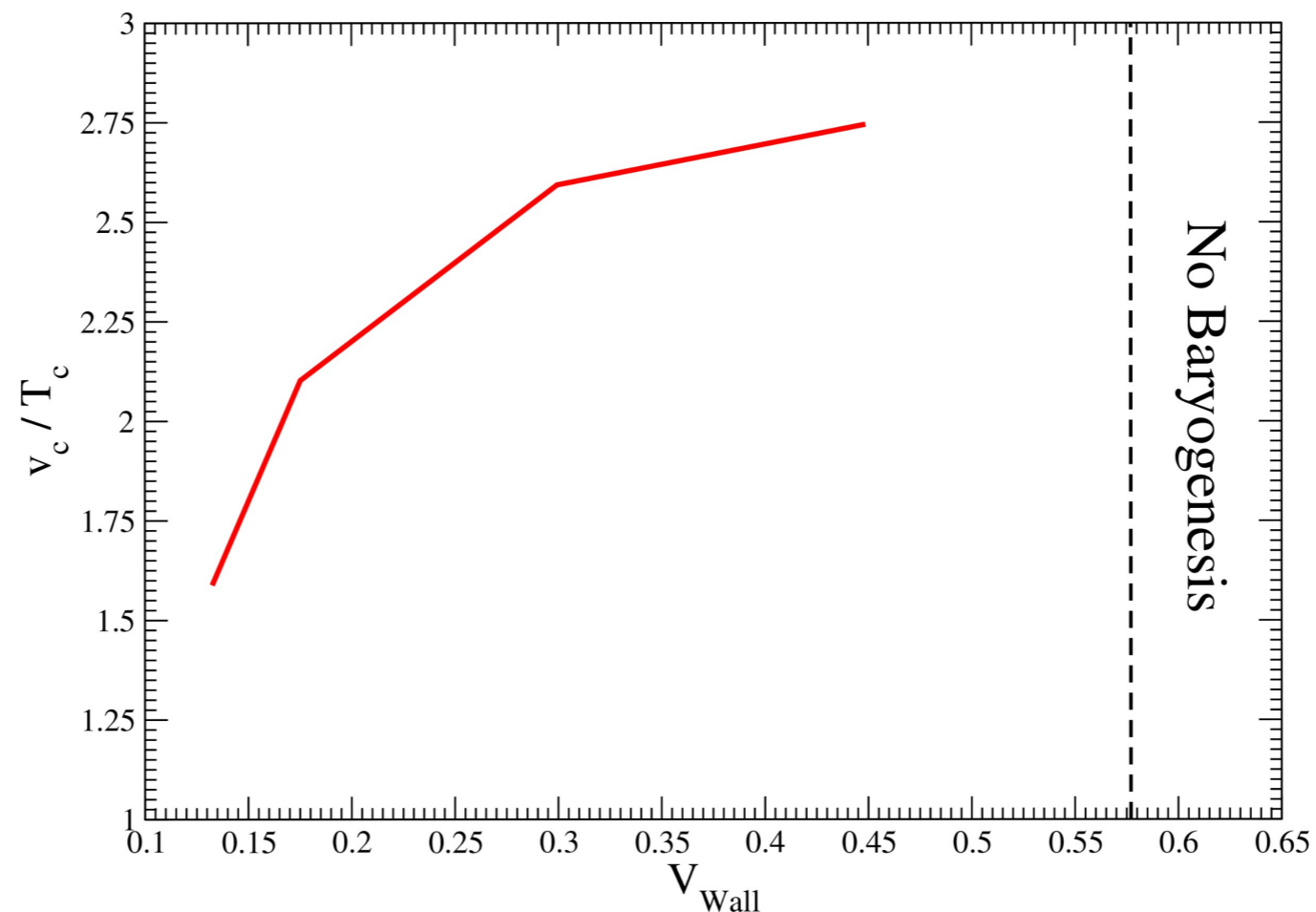
# Baryogenesis with a Second Higgs



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Towards a Global Picture of Baryogenesis in the 2HDM

Dorsch, Huber, Konstandin, JMN, arXiv:16xx.xxxxx

## Before I Conclude... (in 30s)

Introducing a new 2HDM Public Code → **py2HDM**  
(soon to be!) G. Dorsch, K. Mimasu, (JMN)

⇒ The Usual Things:

Theoretical Constraints: Vacuum Stability, Perturbativity, Unitarity. EWPO. Scalar BRs.

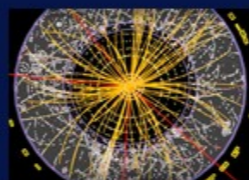
⇒ + CP Violation (Scalar BRs & EDMs) 2-loop Running of Quartics  
+ EW Phase Transition

Modular & "Pure-Execution" Modes

(since not everyone will be interested  
in computing EDMs or the strength of  
the EW Phase Transition!)

We would be very happy to get your feedback!





**KEEP  
CALM  
AND FIND THE  
HIGGS  
BOSON<sub>s</sub>**