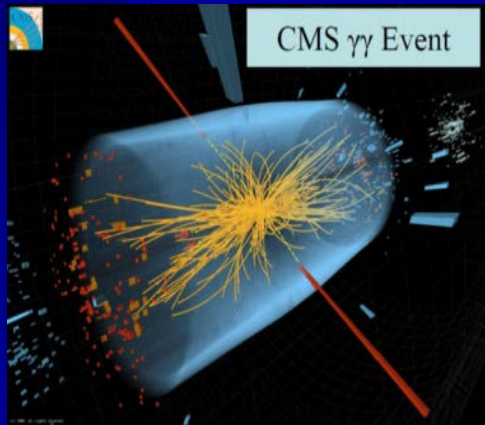
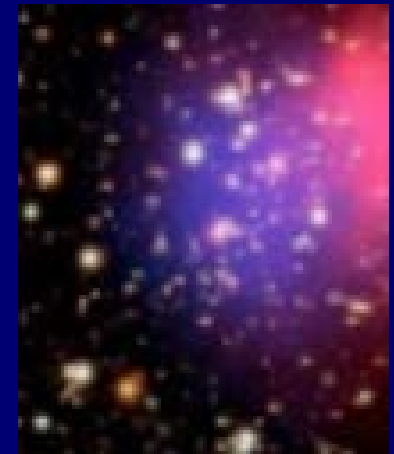


The Inert Doublet Model and not only



Maria Krawczyk
University of Warsaw



In coll. with I. Ginzburg, K. Kanishev, D.Sokołowska, B. Świeżewska, G. Gil, P.Chankowski, M. Matej, N. Darvishi, A. Ilnicka, T. Robens, L. Diaz-Cruz, C. Bonilla, S. Najjari

Inert Doublet Model (IDM) Ma, ... '78

- a model with two SU(2) doublets with Barbieri '06 an **exact** Z_2 symmetry (Lagrangian & vacuum)

Higgs & DM OK

- Evolution of Universe from EWs to Inert phase in one, two or three steps, with 1st or 2nd order phase transitions (*T2 evolution, Ginzburg, ... PRD2010*)

- Strong enough first-order phase transition needed for baryogenesis (*G. Gil Msc'2011, G. Gil, P. Chankowski, MK PL.B 2012*);

- Metastability of vacua in IDM (Chowdhury... Borah & Cline... *B. Świeżewska 2015*)

- IDM+complex singlet *Bonilla, DiazCruz, Sokołowska, Darvishi, MK'14*

- SM+complex singlet *Darvishi, MK'15, Darvishi 2016(baryogenesis)*

Higgs particle at LHC -summer 2016

ATLAS+CMS Run 1

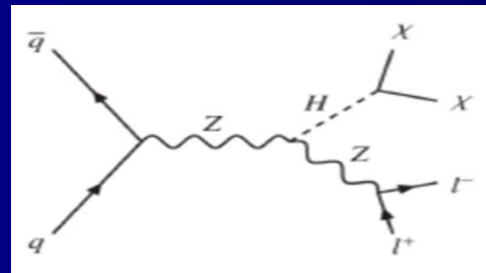
arXiv:1606.02266v1 [hep-ex] 7 Jun

SM-like scenario observed

- Mass 125.09 ± 0.24 GeV $ZZ \rightarrow 4l, \gamma\gamma$
- Total width < 23 MeV (95%CL); SM ~ 4 MeV
- Signal strengths $\mu = R = \sigma \times Br / (\sigma \times Br)_{SM}$; SM = 1

global	$1.09 \pm 0.11/0.10$	
$\gamma\gamma$	$1.14 \pm 0.19/0.18$	$\rightarrow R_{\gamma\gamma}$

- Invisible decay
BR < 0.34 (95% CL)
- Spin/CP $J^{CP} = 0^+$

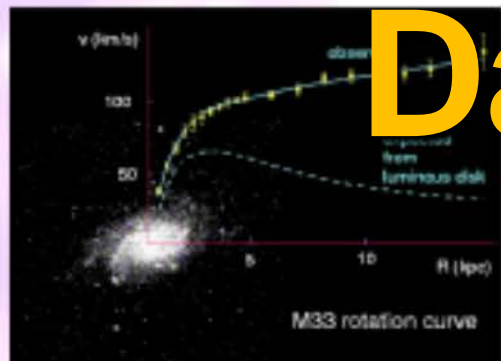


Rotation curves of galaxies

Gravitational lensing

Bullet cluster

Dark matter



Morsolli, Corfu 2014

Relic DM density

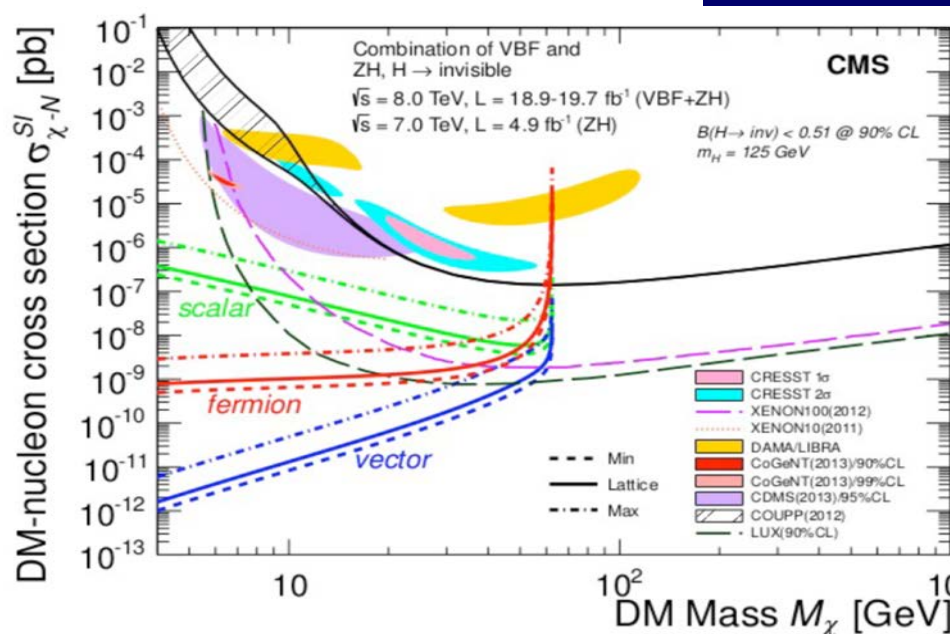
$$0.1018 < \Omega_{DM} h^2 < 0.1234$$

WMAP
3 σ

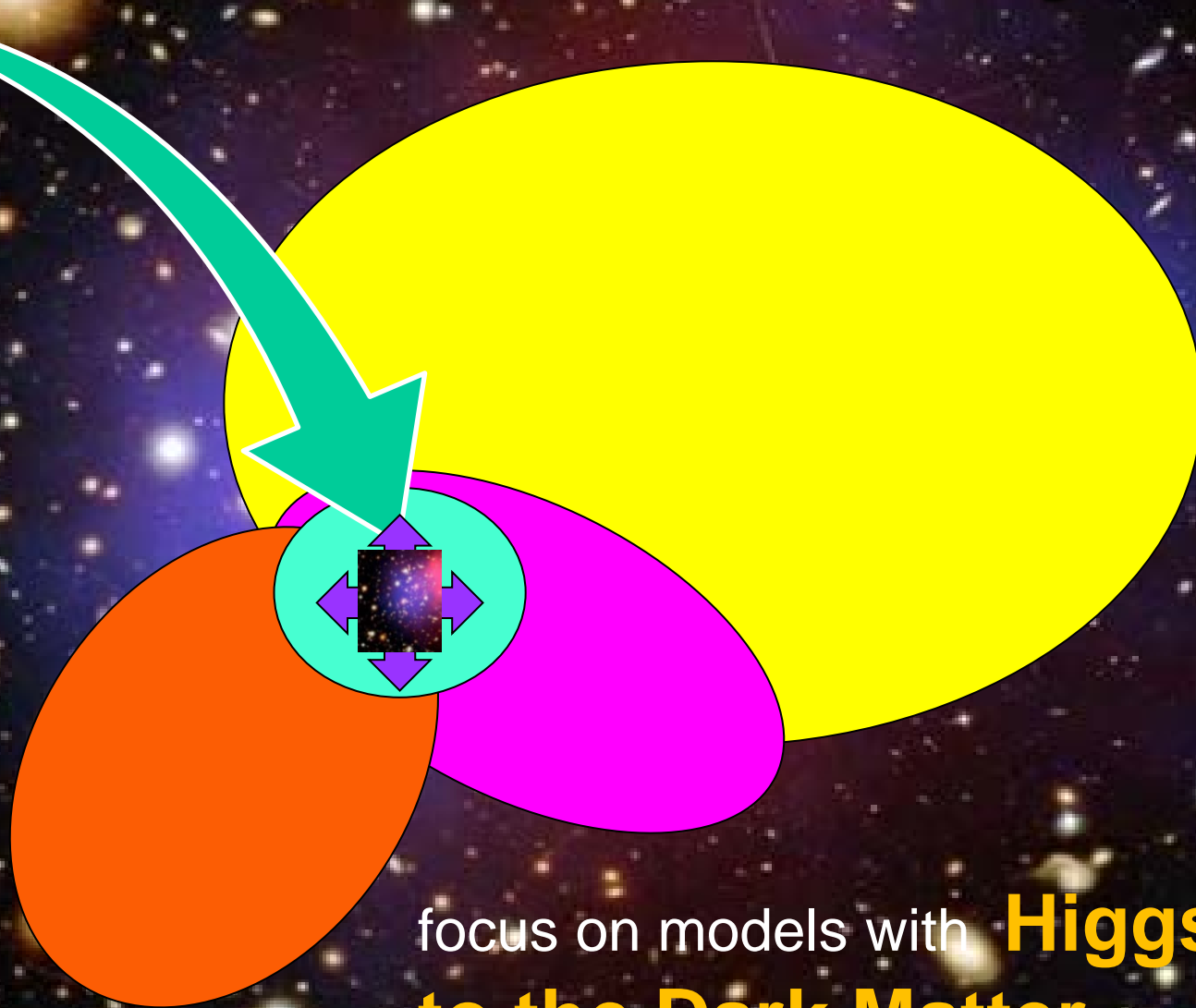
$$0.1118 < \Omega_{DM} h^2 < 0.128$$

PLANCK

Direct DM detection



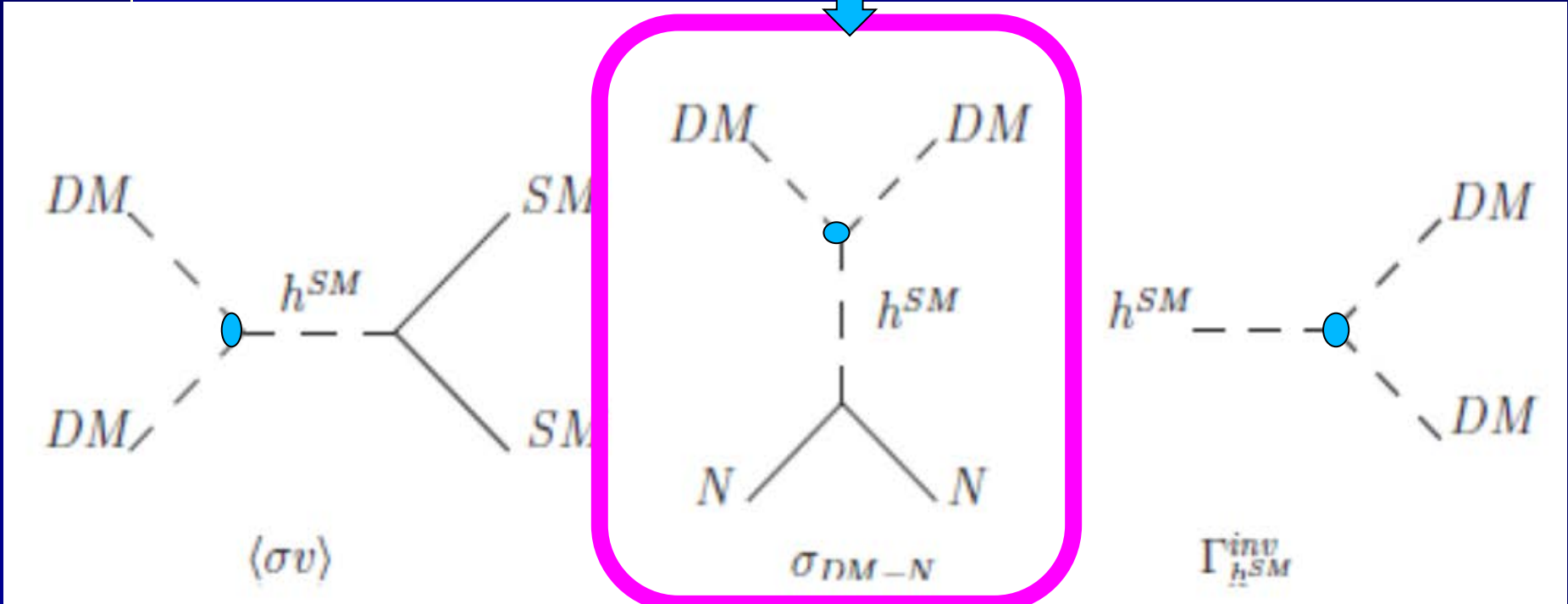
SM - like scenario in many models



focus on models with **Higgs Portal**
to the Dark Matter

Higgs portal with the SM-like h

direct detection



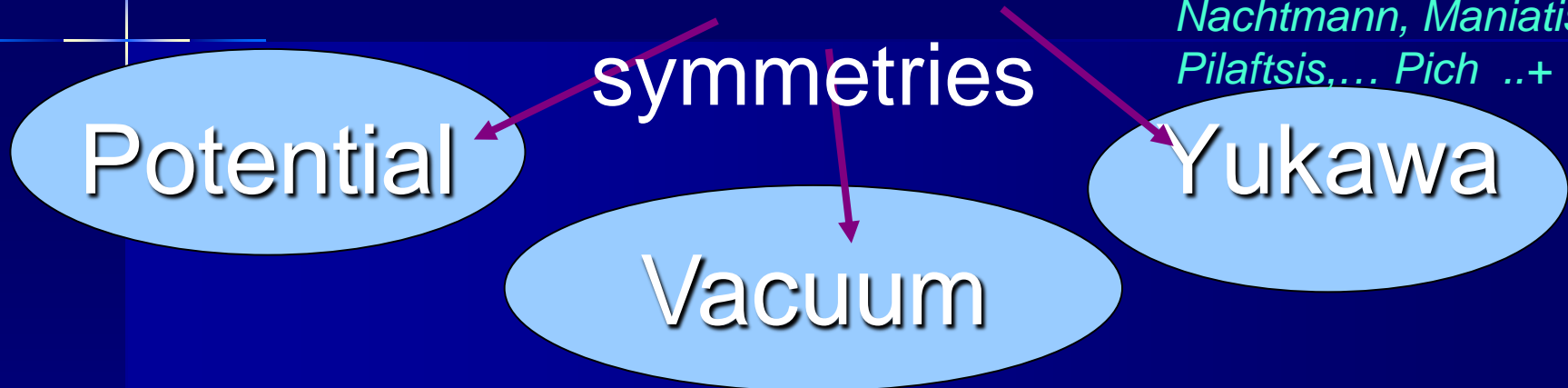
relic DM density

invisible decay

*T.D. Lee 1973
CP violation
in scalar sector*

*Branco, Rebelo, Ferreira
Silva, Lavoura, Sher, Ma
Haber, Gunion, Grimus
Ginzburg, MK, Osland,
Grzadkowski, Ivanov
Nachtmann, Maniatis,
Pilaftsis, ... Pich ..+*

2HDM's



Two Higgs Doublet Models

Two doublets of $SU(2)$ ($Y=1, \rho=1$) - Φ_1, Φ_2

Masses for $W^{+/-}, Z$, no mass for photon?

Fermion masses via Yukawa interaction –

various models: Model I, II, III, IV, X, Y, ...

5 scalars: 3 neutral and 2 charged

Z_2 symmetric Lagrangian of 2HDM

Potential $V =$

Branco, Rebelo ,85 (CP conserved)

$$\begin{aligned} & \frac{1}{2}\lambda_1(\Phi_1^\dagger\Phi_1)^2 + \frac{1}{2}\lambda_2(\Phi_2^\dagger\Phi_2)^2 - \frac{1}{2}m_{11}^2(\Phi_1^\dagger\Phi_1) - \frac{1}{2}m_{22}^2(\Phi_2^\dagger\Phi_2) \\ & + \lambda_3(\Phi_1^\dagger\Phi_1)(\Phi_2^\dagger\Phi_2) + \lambda_4(\Phi_1^\dagger\Phi_2)(\Phi_2^\dagger\Phi_1) + \frac{1}{2}[\lambda_5(\Phi_1^\dagger\Phi_2)^2 + \text{h.c.}] \end{aligned}$$
$$\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5$$

Z_2 symmetry transf.: $\Phi_1 \rightarrow \Phi_1$ $\Phi_2 \rightarrow -\Phi_2$

Yukawa interaction

Model I – one doublet Φ_1 couples to all fermions

Vacuum state ?
various possible

M. Krawczyk, Lisbon 2016

positivity (stability) constraints

$$\lambda_1 > 0, \quad \lambda_2 > 0, \quad R + 1 > 0, \quad R_3 + 1 > 0$$

$$\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5, \quad R = \lambda_{345}/\sqrt{\lambda_1\lambda_2}, \quad R_3 = \lambda_3/\sqrt{\lambda_1\lambda_2}.$$

Extrema (\rightarrow vacua)

Ma78, Velhinho, Santos, Barroso..94

Z_2 symmetry $\Phi_1 \rightarrow \Phi_1, \Phi_2 \rightarrow -\Phi_2$

new notation: $\Phi_1 \rightarrow \Phi_S$ & $\Phi_2 \rightarrow \Phi_D$ (**D symmetry**)

$$\langle \phi_S \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_S \end{pmatrix}, \quad \langle \phi_D \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} u \\ v_D \end{pmatrix}$$

v_S, v_D, u - real
 $v^2 = v_S^2 + v_D^2 + u^2$

----- **$u=0$** -----

EWs

Inert

Inert-like

Mixed (Normal, MSSM like)

EWs

I_1

I_2

M

$$v_D = v_S = 0$$

$$v_D = 0$$

$$v_S = 0$$

$$v_D, v_S \neq 0$$

----- **$u \neq 0$** -----

Charge Breaking

CB

$$v_D = 0$$

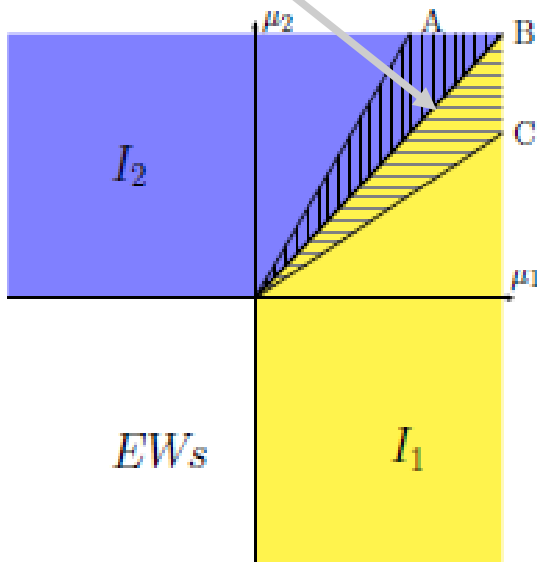
Phase diagrams for D-sym. V

$$\mu_1 = \frac{m_{11}^2}{\sqrt{\lambda_1}}, \quad \mu_2 = \frac{m_{22}^2}{\sqrt{\lambda_2}}$$

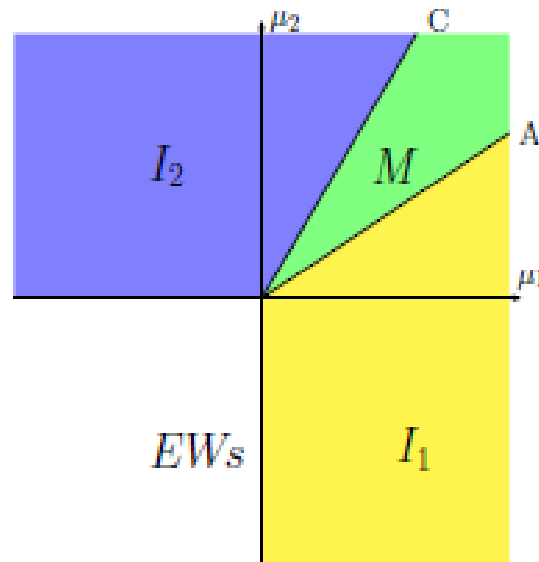
$$\mathcal{E}_{I_1} - \mathcal{E}_M = \frac{(m_{11}^2 \lambda_{345} - m_{22}^2 \lambda_1)^2}{8\lambda_1^2 \lambda_2 (1 - R^2)}$$

$$R = \lambda_{345} / \sqrt{\lambda_1 \lambda_2}$$

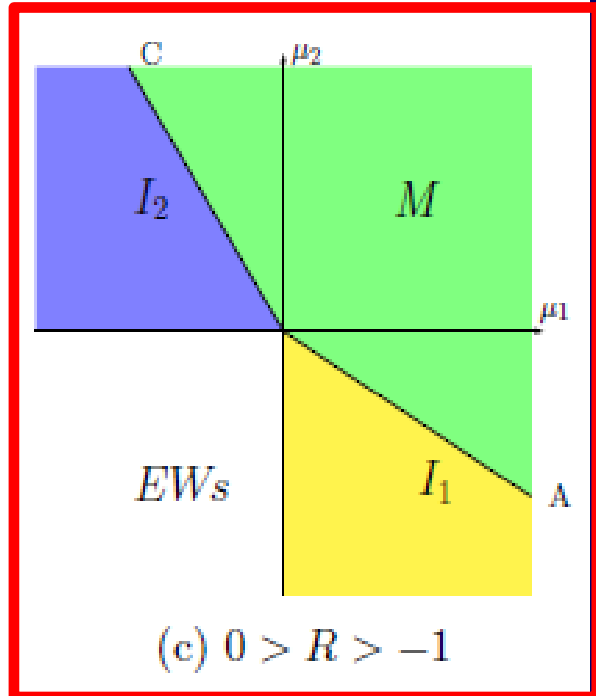
coexistence of I_1 and I_2 minima



(a) $R > 1$



(b) $1 > R > 0$



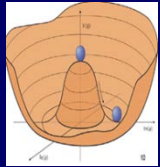
(c) $0 > R > -1$

Inert (I_1) vacuum
for $M_h = 125$ GeV \rightarrow fixed μ_1

here $\lambda_{345} < 0$!

Inert Doublet Model

Φ_S as in SM (BEH)



$$\Phi_S = \begin{pmatrix} \phi^+ \\ \frac{v+h+i\zeta}{\sqrt{2}} \end{pmatrix}$$

Higgs boson h (SM-like)

Φ_D – no vev

$$\Phi_D = \begin{pmatrix} H^+ \\ H+iA \end{pmatrix}$$

(no Higgses!)

4 scalars H^+, H^-, H, A

no interaction with fermions

D symmetry $\Phi_S \rightarrow \Phi_S \quad \Phi_D \rightarrow -\Phi_D$ exact

▸ D parity

▸ only Φ_D has odd D-parity

▸ the lightest scalar stable - DM candidate (H)

▸ (Φ_D dark doublet with dark scalars)

Inert case - masses

- SM-like Higgs scalar h

$$M_h^2 = m_{11}^2 = \lambda_1 v^2 = (125 \text{ GeV})^2$$

- Dark particles D

$$M_{H^+}^2 = -\frac{m_{22}^2}{2} + \frac{\lambda_3}{2} v^2$$

m_{22}^2 arbitrary,

so if large negative \rightarrow
 H, H^+, A heavy, degenerate

 H – dark matter

$\lambda_5 < 0$ and $\lambda_{45} < 0$

$$M_H^2 = -\frac{m_{22}^2}{2} + \frac{\lambda_3 + \lambda_4 + \lambda_5}{2} v^2$$

$$M_A^2 = -\frac{m_{22}^2}{2} + \frac{\lambda_3 + \lambda_4 - \lambda_5}{2} v^2$$

Testing Inert Doublet Model

❖ Theoretical constraints

vacuum stability,

perturbative unitarity

condition for Inert vacuum

Ma'2006, Barbieri 2006, Dolle, Su, Gorczyca(Świeżewska), MSc T2011, 1112.4356, ...5086, ..1305. Posch 2011, Arhrib..2012, Chang, Stal ..2013....

$$\frac{m_{11}^2}{\sqrt{\lambda_1}} \geq \frac{m_{22}^2}{\sqrt{\lambda_2}}$$

Swiezewska

❖ Detailed study of

- the SM-like h

❖ Study of dark scalars $D = (\mathbf{H}, A, H^+, H^-)$

- the dark scalars D in pairs!

D couple to $V = W/Z$ (eg. $AZH, H^- W^+ H$), not DVV !

Quartic selfcouplings D^4 proportional to λ_2

Couplings with Higgs:

$$hHH \sim \lambda_{345}$$

$$h H^+ H^- \sim \lambda_3$$

13

LHC – Higgs H_{125} data \rightarrow h (IDM)

Direct couplings to W/Z and fermions - as in SM

Loop coupling to gg – as in SM

Loop coupling to $\gamma\gamma$, $Z\gamma$ – extra H+ contribution

Total width – extra contributions $h \rightarrow HH, AA, H+H-$

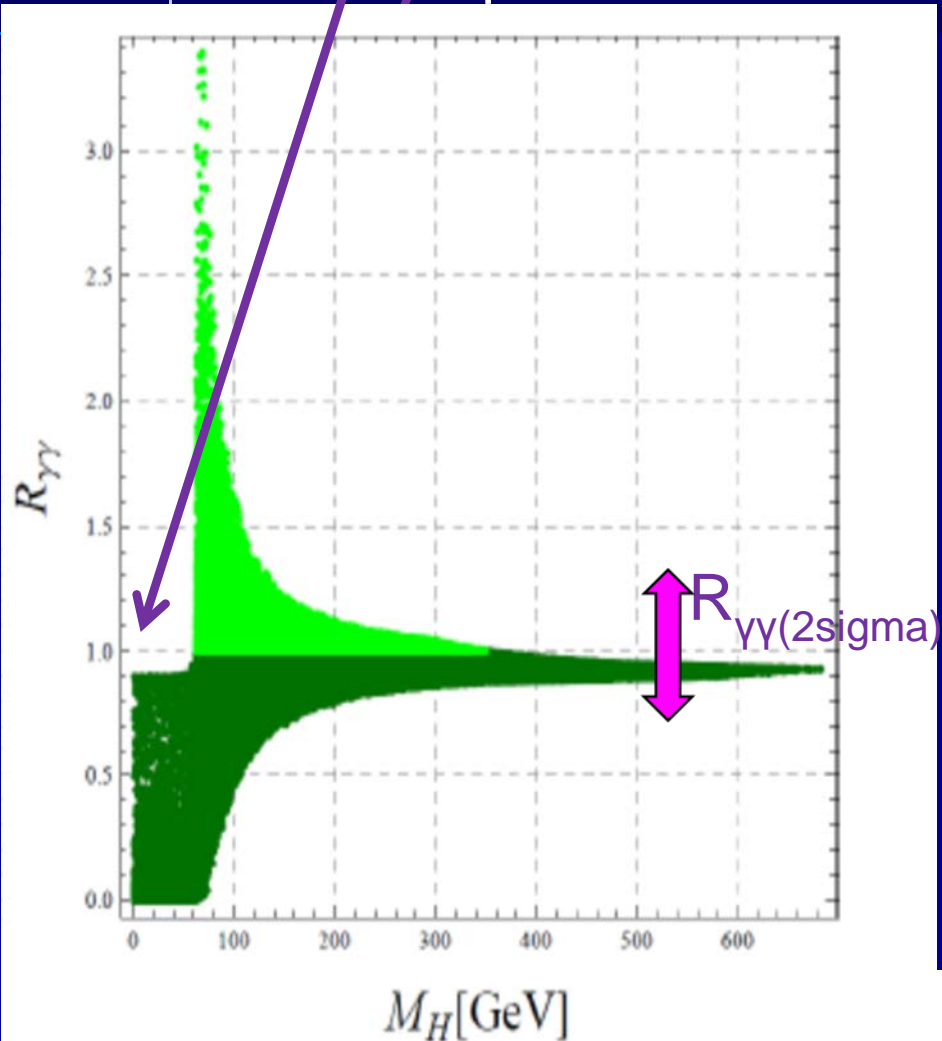
Invisible decay $h \rightarrow HH$

$$R_{\gamma\gamma} = \frac{\Gamma_{\gamma\gamma}}{\Gamma_{\gamma\gamma}^{\text{SM}}} \frac{\Gamma_{\text{tot}}^{\text{SM}}}{\Gamma_{\text{tot}}}$$

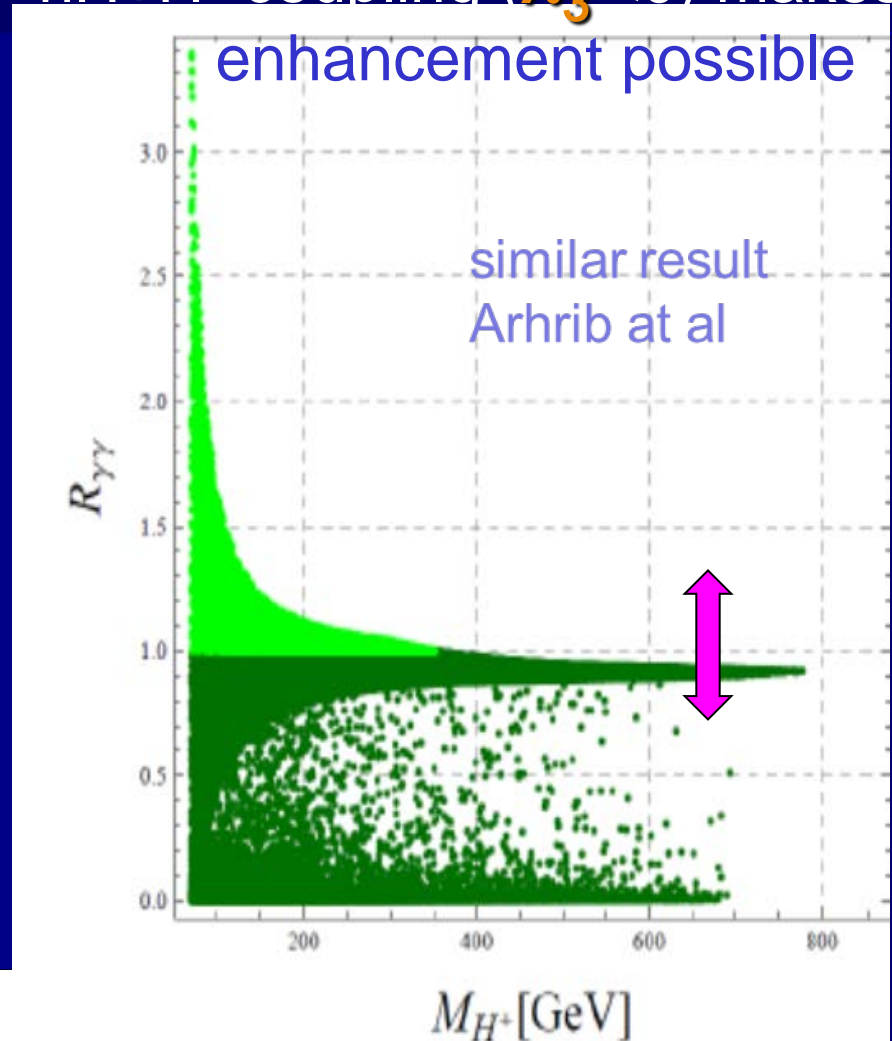
invisible decays important : $R_{\gamma\gamma} > 1$
only if DM mass above 62.5 GeV

$R_{\gamma\gamma}$ as a function of mass H, H^+

Invisible decays makes enhancement impossible



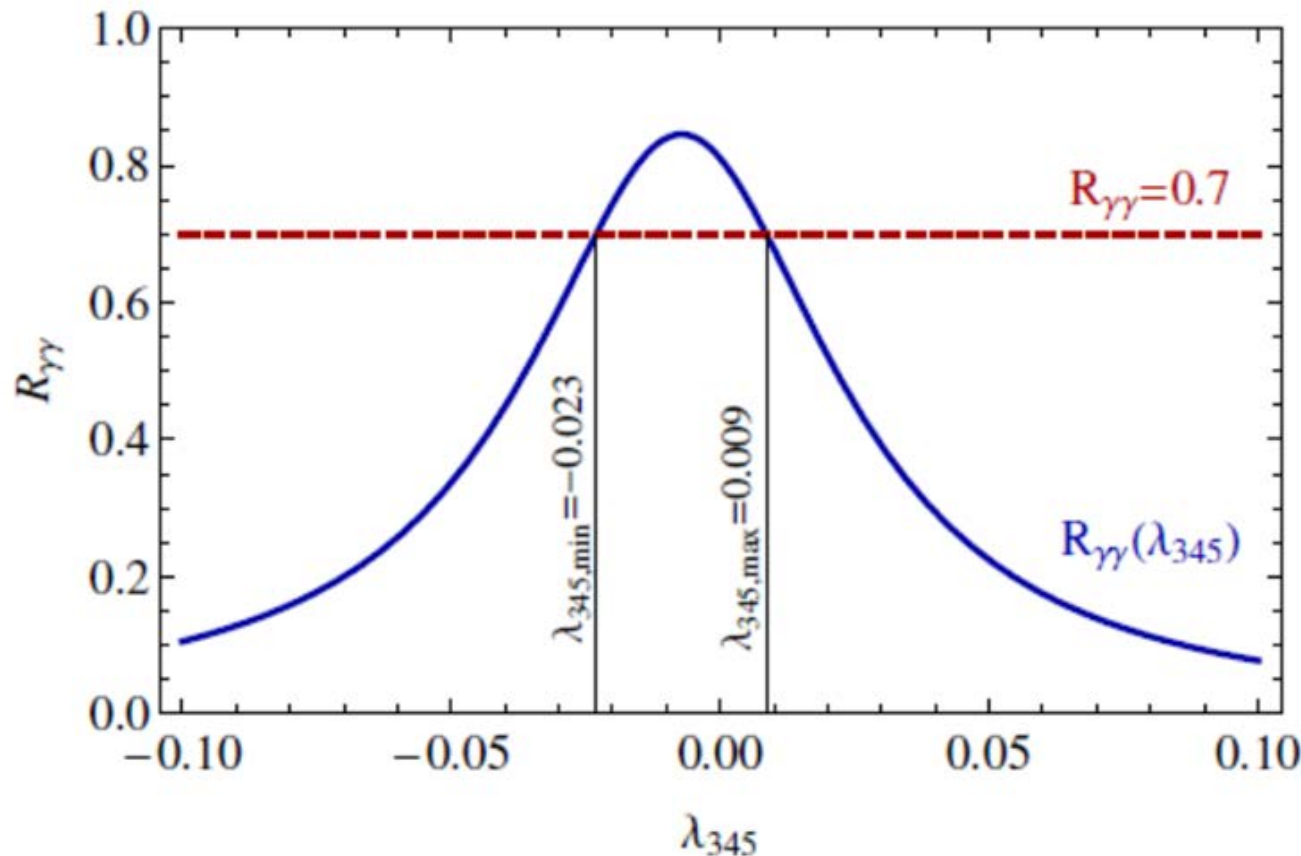
Light H^+ with proper sign of hH^+H^- coupling ($\lambda_3 < 0$) makes enhancement possible



Signal strength $R_{\gamma\gamma}$ if below 1 sensitive to hHH coupling

2013

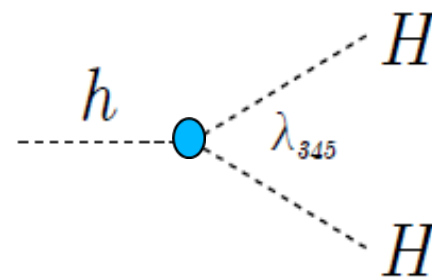
ATLAS : $R_{\gamma\gamma} = 1.65 \pm 0.24(\text{stat})_{-0.18}^{+0.25}(\text{syst})$,
CMS : $R_{\gamma\gamma} = 0.79_{-0.26}^{+0.28}$



ATLAS+CMS 2016
 1.14 ± 0.19

Invisible h decay \rightarrow coupling hHH

- $h \rightarrow HH$ – invisible decay (H is stable)
- augmented total width of the Higgs boson, $\Gamma(h \rightarrow HH) \sim \lambda_{345}^2$

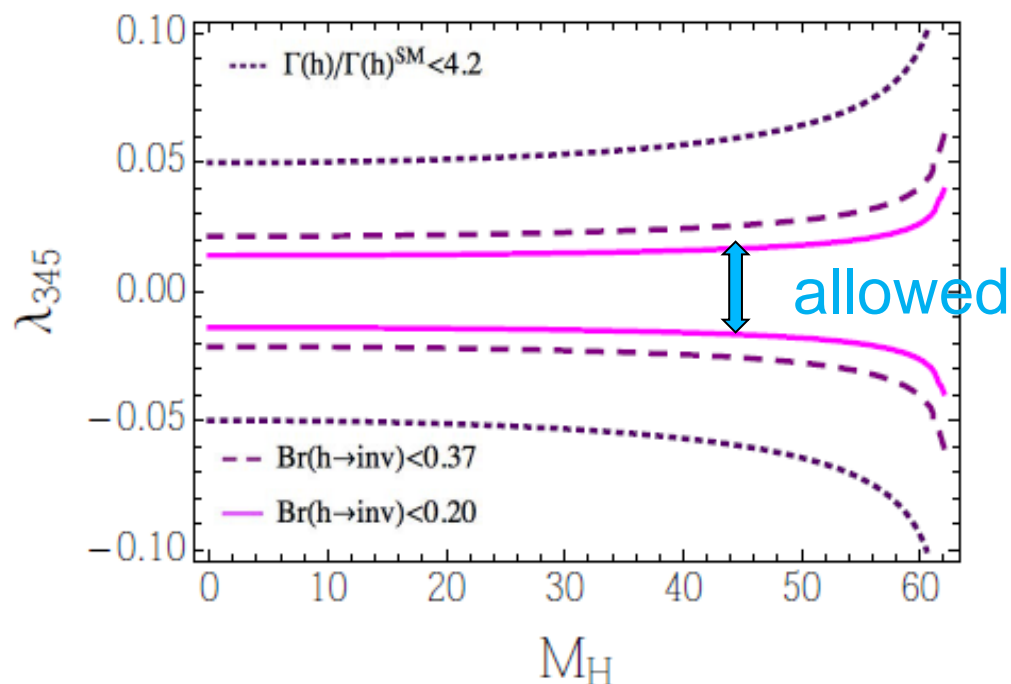


LHC:

- $\text{Br}(h \rightarrow \text{inv}) < 37\%$,
- $\Gamma(h)/\Gamma(h)^{\text{SM}} < 4.2$

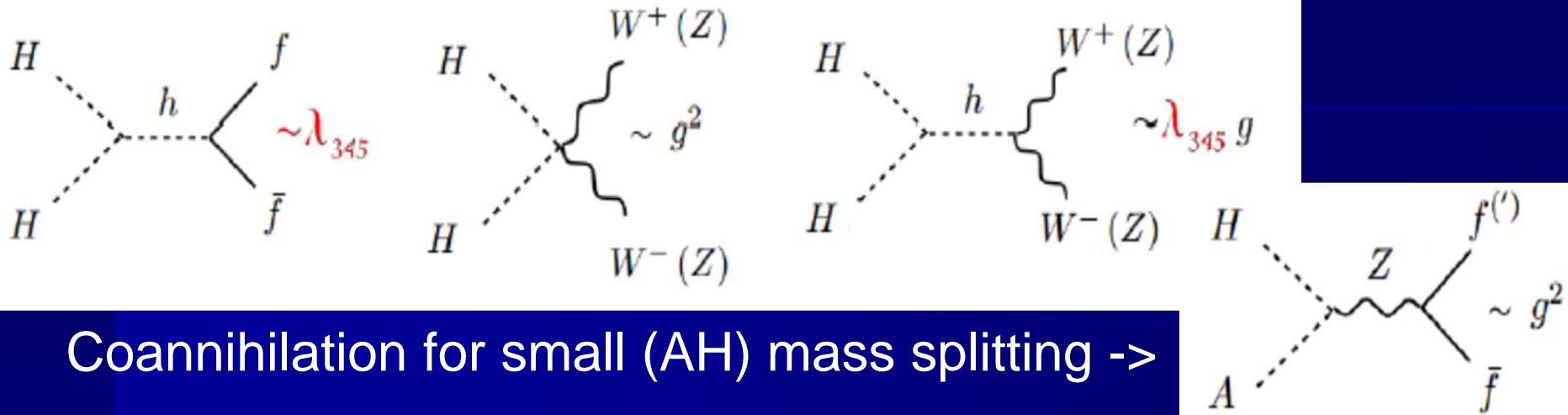
global fit:

- $\text{Br}(h \rightarrow \text{inv}) \lesssim 20\%$



[G. Bélanger, B. Dumont, U. Ellwanger, J. F. Gunion, S. Kraml, PLB 723 (2013) 340; ATLAS-CONF-2014-010; 2014 CMS-PAS-HIG-14-002]

Relic DM density \rightarrow H (IDM)



Coannihilation for small (AH) mass splitting \rightarrow

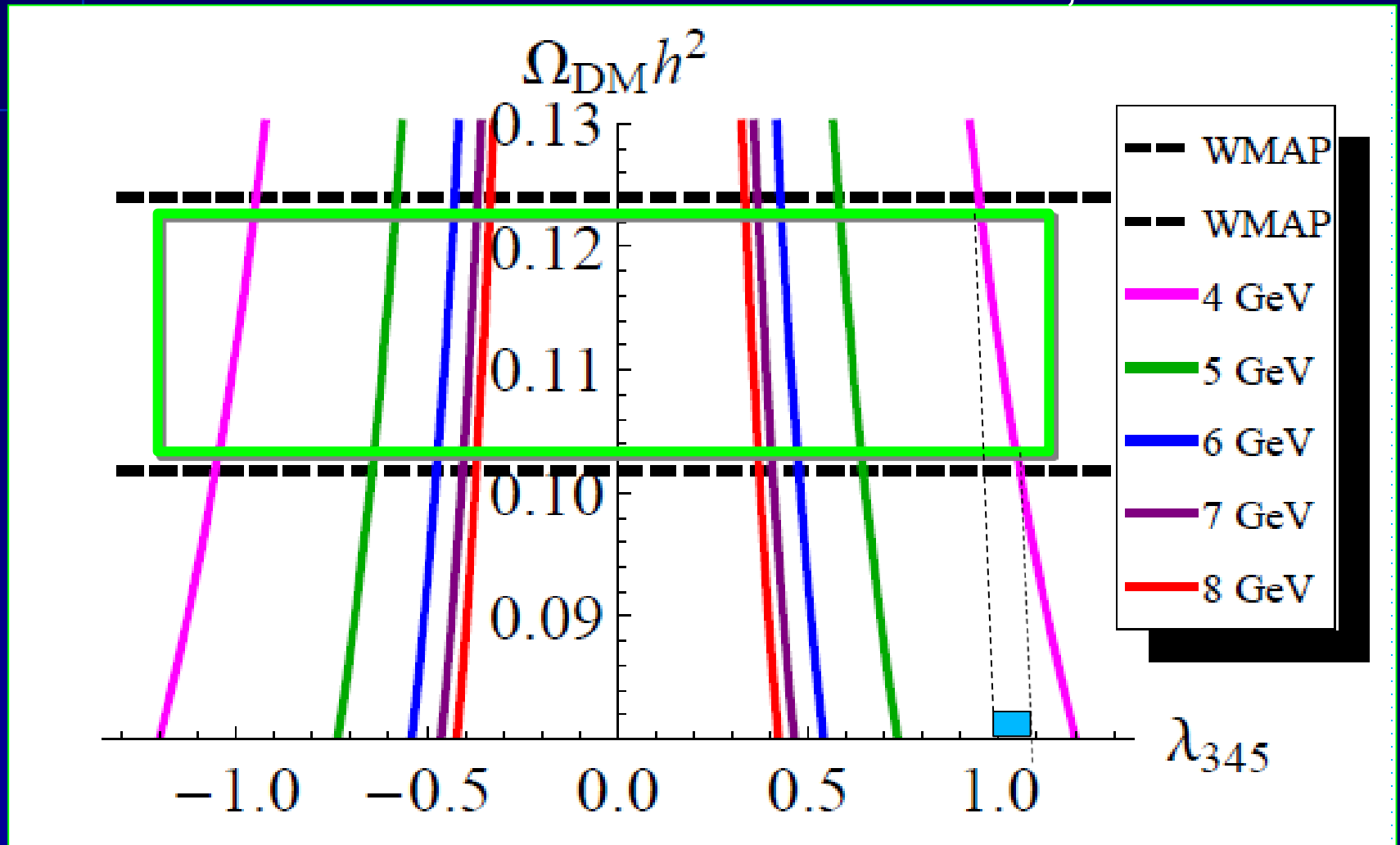
WMAP

$$\Omega_{DM} h^2 = 0.1126 \pm 0.0036.$$

- low DM mass $M_H \lesssim 10$ GeV, $g_{HHh} \sim \mathcal{O}(0.5)$
- medium DM mass $M_H \approx (40 - 160)$ GeV, $g_{HHh} \sim \mathcal{O}(0.05)$
- high DM mass $M_H \gtrsim 500$ GeV, $g_{HHh} \sim \mathcal{O}(0.1)$

WMAP window for light H (DM)

D. Sokołowska, 2013

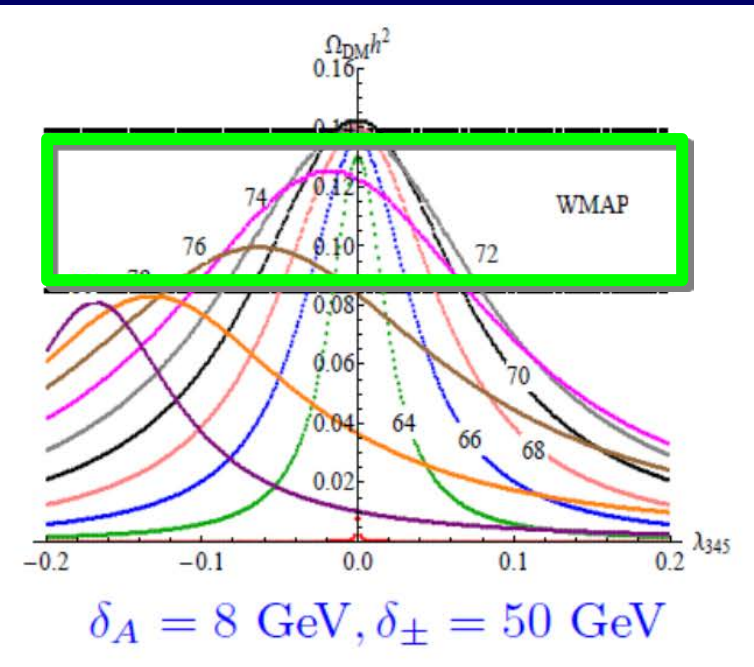
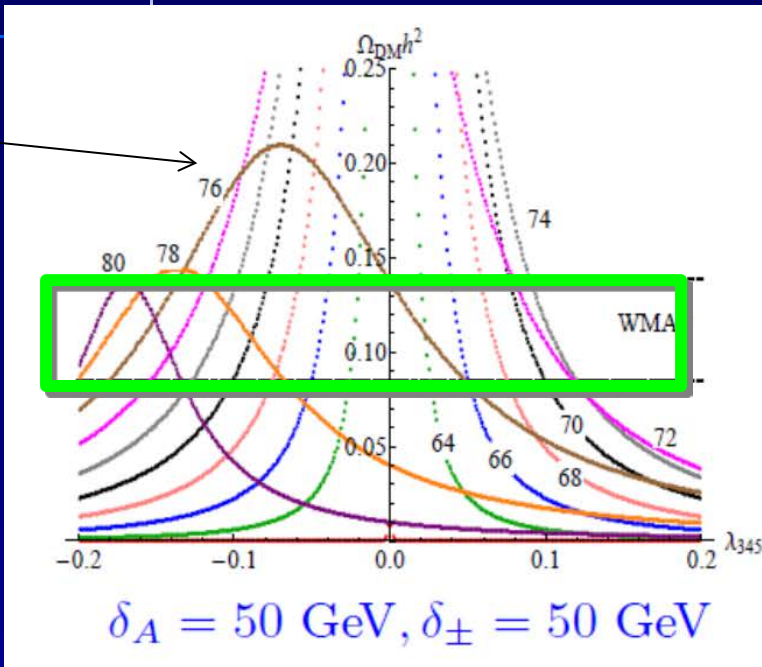


using MicrOmegas

Relic density for DM with mass > 64 GeV

D. Sokołowska, 2013

$$M_{A,H^\pm} = M_H + \delta_{A,\pm}$$



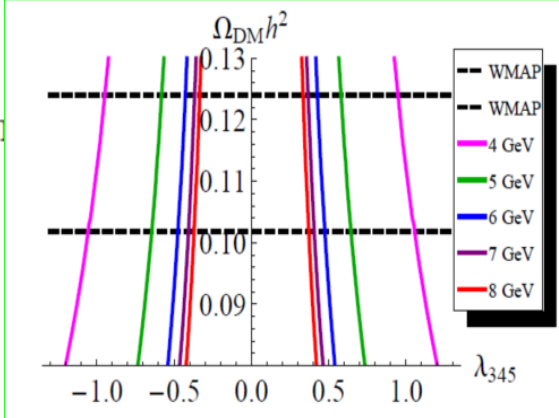
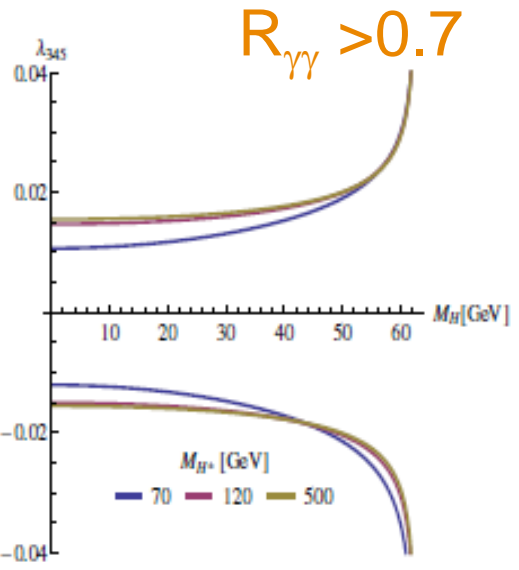
above 76 GeV asymmetry due to annihilation to gauge bosons

Low mass H – excluded by LHC!

Constraining Inert Dark Matter by $R_{\gamma\gamma}$ and WMAP data

[M. Krawczyk, D. Sokołowska, P. Swaczyna, BŚ, arXiv:1305.6266 [hep-ph], JHEP 2013]

$M_H \lesssim 10 \text{ GeV}$, $M_A \approx M_{H^\pm} \approx 100 \text{ GeV}$
 $h \rightarrow AA$ channel closed, $h \rightarrow HH$ channel open



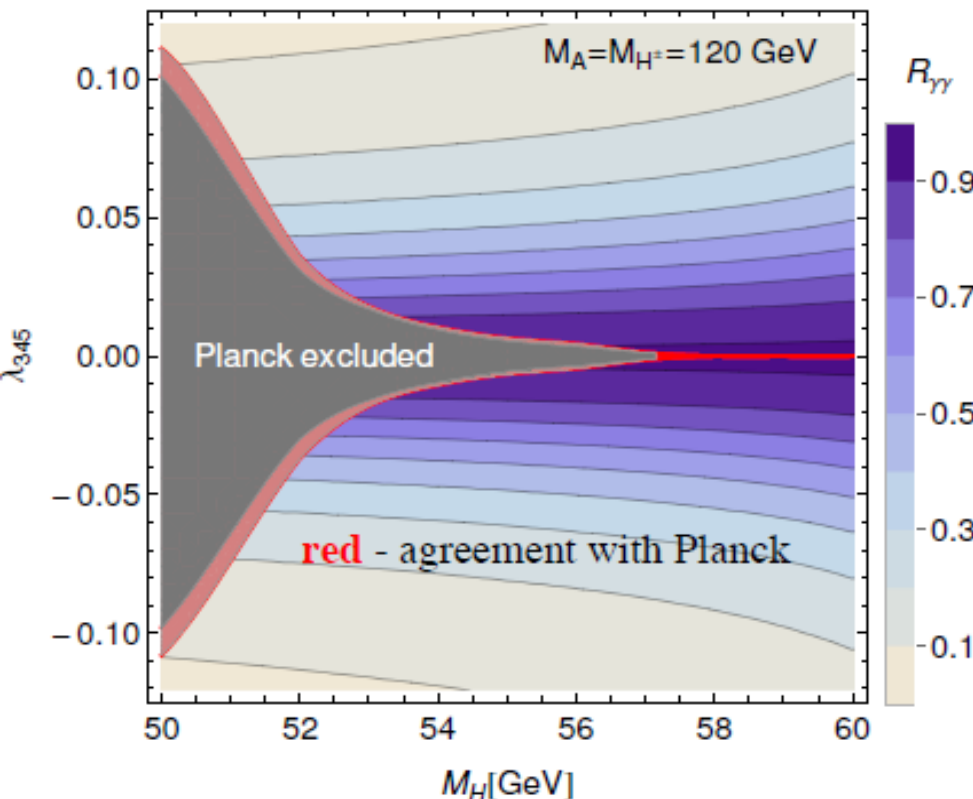
- Proper relic density
 $0.1018 < \Omega_{DM} h^2 < 0.1234 \Rightarrow |\lambda_{345}| \sim \mathcal{O}(0.5)$
- CDMS-II reported event:
 $M_H = 8.6 \text{ GeV} \Rightarrow |\lambda_{345}| \approx (0.35 - 0.41)$
- $R_{\gamma\gamma} > 0.7 \Rightarrow |\lambda_{345}| \lesssim 0.02 \Rightarrow$

Low DM mass excluded

Using PLANCK data

[Planck update: D. Sokołowska, P. Swaczyna, 2014]

$h \rightarrow HH$ open



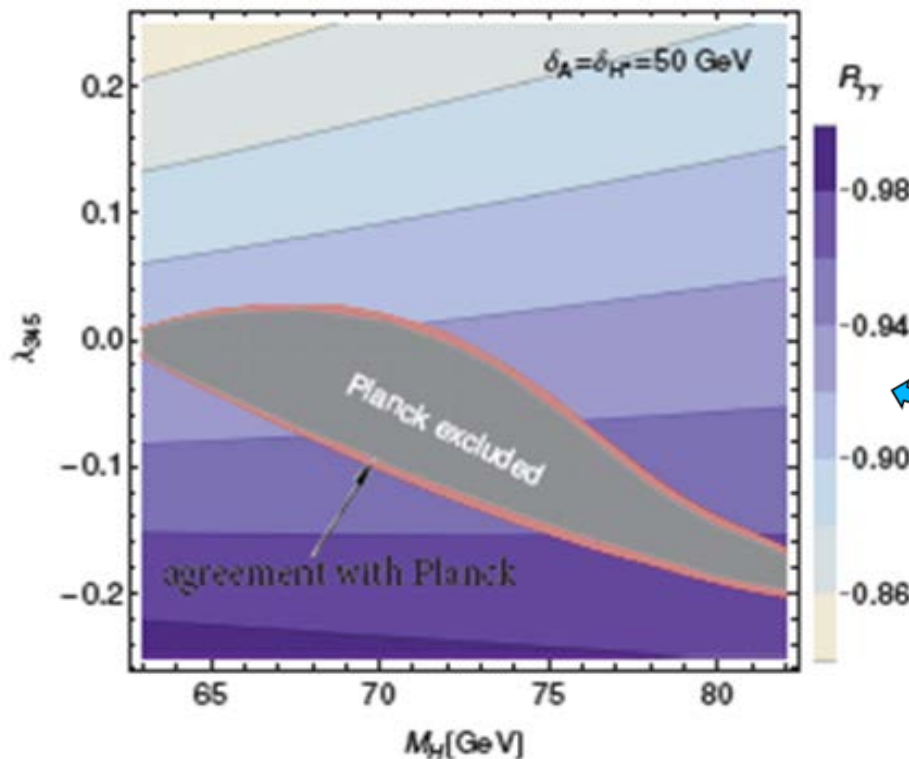
$50 \text{ GeV} < M_H < M_H/2, M_A = M_{H^\pm} = 120 \text{ GeV}$

- light DM ($M_H < 10 \text{ GeV}$)
 \Rightarrow excluded
- intermediate DM 1
($50 \text{ GeV} < M_H < M_H/2$)
 $\Rightarrow M_H > 53 \text{ GeV}$
- intermediate DM 2
($M_H/2 < M_H \lesssim 82 \text{ GeV}$)
 $\Rightarrow R_{\gamma\gamma} < 1$
- heavy DM
($M_H > 500 \text{ GeV}$)
 $\Rightarrow R_{\gamma\gamma} \approx 1$

Using PLANCK data

[Planck update: D. Sokołowska, P. Swaczyna, 2014]

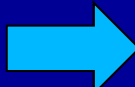
$h \rightarrow HH$ open



- light DM ($M_H < 10$ GeV)
 \Rightarrow excluded
- intermediate DM 1
($50 \text{ GeV} < M_H < M_H/2$)
 $\Rightarrow M_H > 53$ GeV
- intermediate DM 2
($M_H/2 < M_H \lesssim 82$ GeV)
 $\Rightarrow R_{\gamma\gamma} < 1$
- heavy DM
($M_H > 500$ GeV)
 $\Rightarrow R_{\gamma\gamma} \approx 1$

Full scan for IDM (2015) → T. Robens

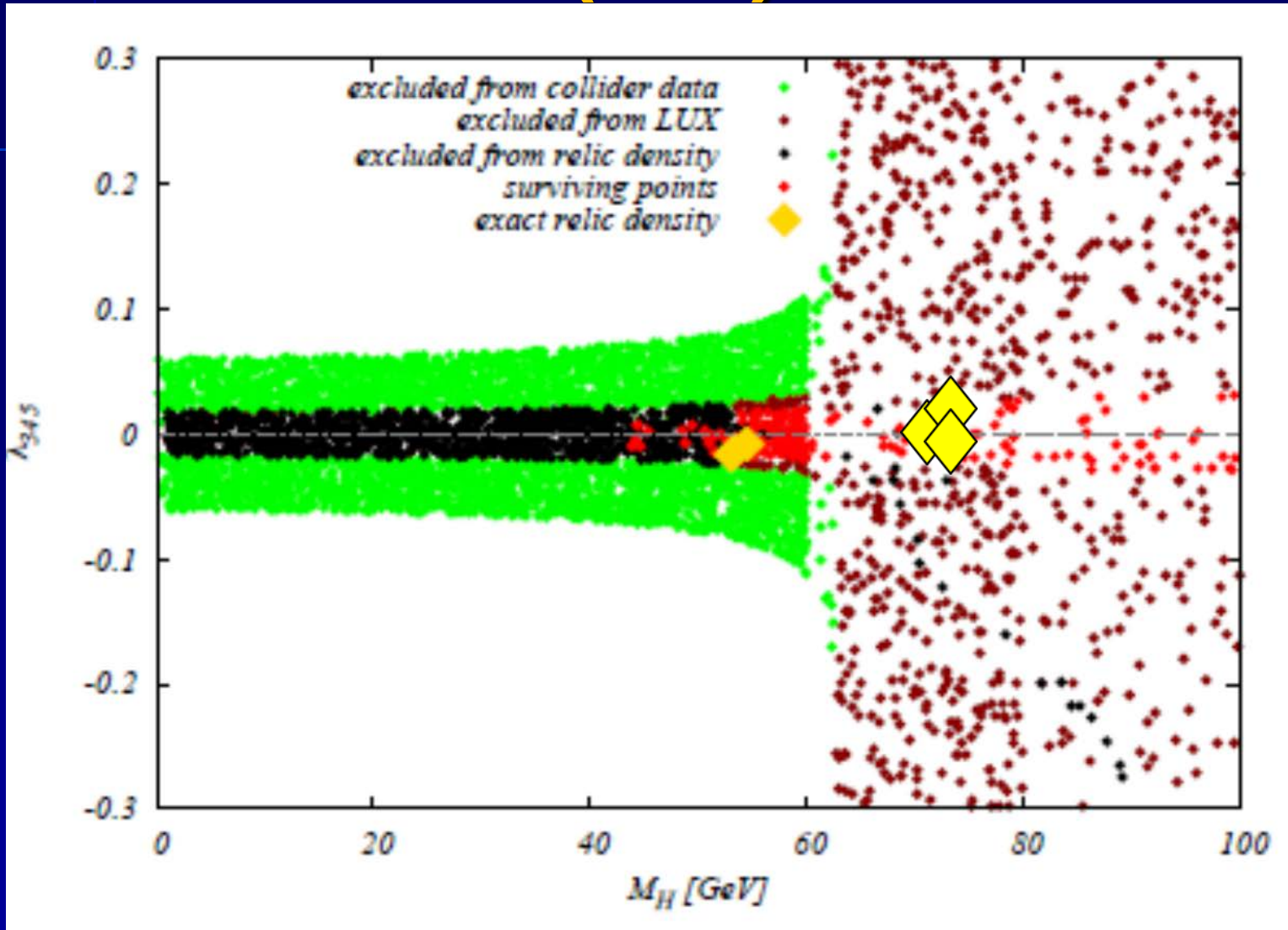
A. Inicka, T. Robens, MK Phys.Rev. D93 (2016)

- Theor. constraints –
stability of the potential (positivity), pert.unitarity,
condition for the Inert vacuum
- STU (from 2014)
- Higgs signal/Higgs bounds
- Lifetime of H^\pm ($< 10^{-7}$ s to decay inside detector)
- Relic density Planck $\Omega < 0.1241$ (95% CL)
- Direct detection LUX
- --> scan over M_H up to 1 TeV
-  Benchmarks

+LEP constraints
h total width
W/Z total width

Low mass H (DM)

1505.04734, 1508.01671



exact relic density 

Limit on mass of DM: $M_H > 45$ GeV !

Benchmarks for LHC II

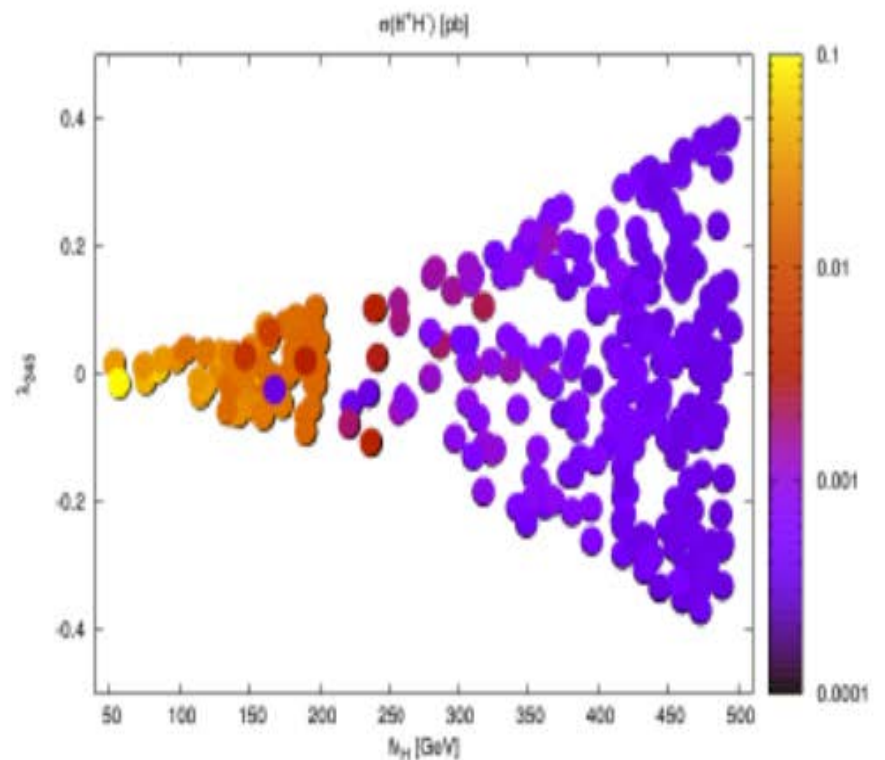
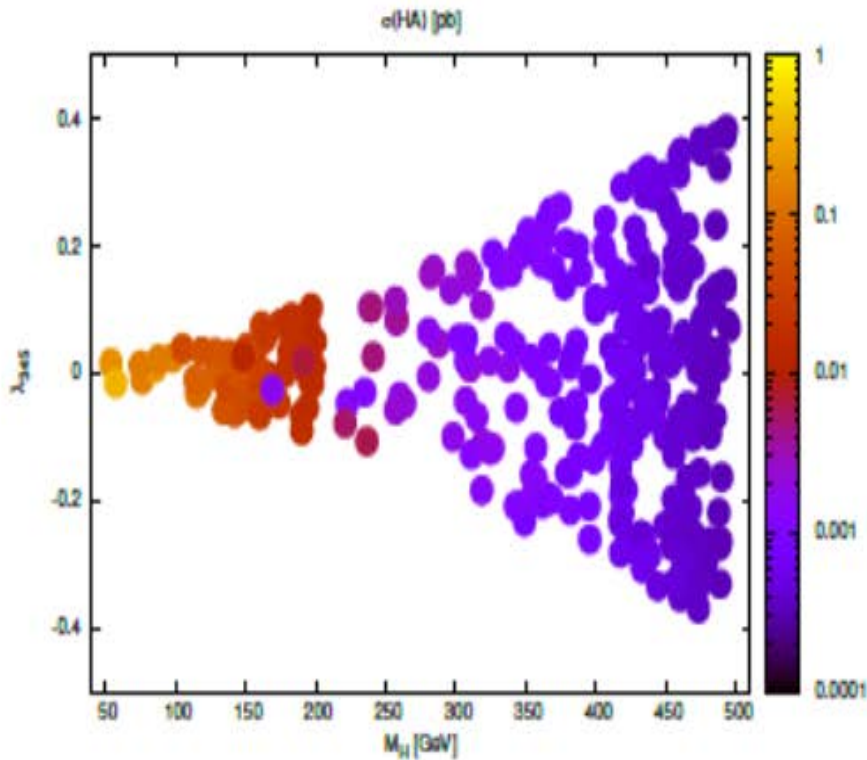
$pp \rightarrow HA: \leq 0.03 \text{ pb},$
 $pp \rightarrow H^\pm H: \leq 0.03 \text{ pb},$
 $pp \rightarrow H^\pm A: \leq 0.015 \text{ pb},$
 $pp \rightarrow H^+ H^-: \leq 0.01 \text{ pb},$
 $pp \rightarrow AA: \leq 0.0015 \text{ pb}.$

→ T. Robens talk

λ_{345}

HA

H+H-



$A \rightarrow Z H = 100\%$

$M_{H^\pm} - M_H > 100 \text{ MeV}$

M_H


Cross section in pb, mass in GeV

IDM at ILC/CLIC

S.Najjari, M.Hashemi, M.Krawczyk,
AF. Zarnecki JHEP 1602 (2016) 187,
arXiv:1512.01175

- ▶ $e^+e^- \rightarrow H^+H^-$
- ▶ $e^+e^- \rightarrow H A$

0.5 and 1 TeV for 500 fb^{-1}

Taking low mass benchmarks BP from our scan (difficult for LHC) reconstruction of dark particle masses with accuracy 100 MeV possible 

BP	m_H	m_A	m_{H^\pm}
$BP1$	57.5	113	123
$BP2$	85.5	111	140
$BP3$	128	134	176

	theo.	123	140	176
m_{H^\pm}	0.5	117.1 ± 7.2	136.8 ± 10.0	167.4 ± 14.6
	1	112.7 ± 4.5	131.4 ± 5.0	172.2 ± 5.5
	theo.	57.5	85.5	128
m_H	0.5	58.5 ± 3.6	88.9 ± 6.5	127.2 ± 9.5
	1	53.0 ± 2.1	81.5 ± 3.1	129.1 ± 1.2
	theo.	113	111	134
m_A	0.5	113.9 ± 3.7	114.3 ± 6.6	133.1 ± 9.6
	1	104.6 ± 2.2	105.0 ± 3.2	134.8 ± 1.3

Main decays
 $H^+ \rightarrow W^+ + H$
 $A \rightarrow Z H$

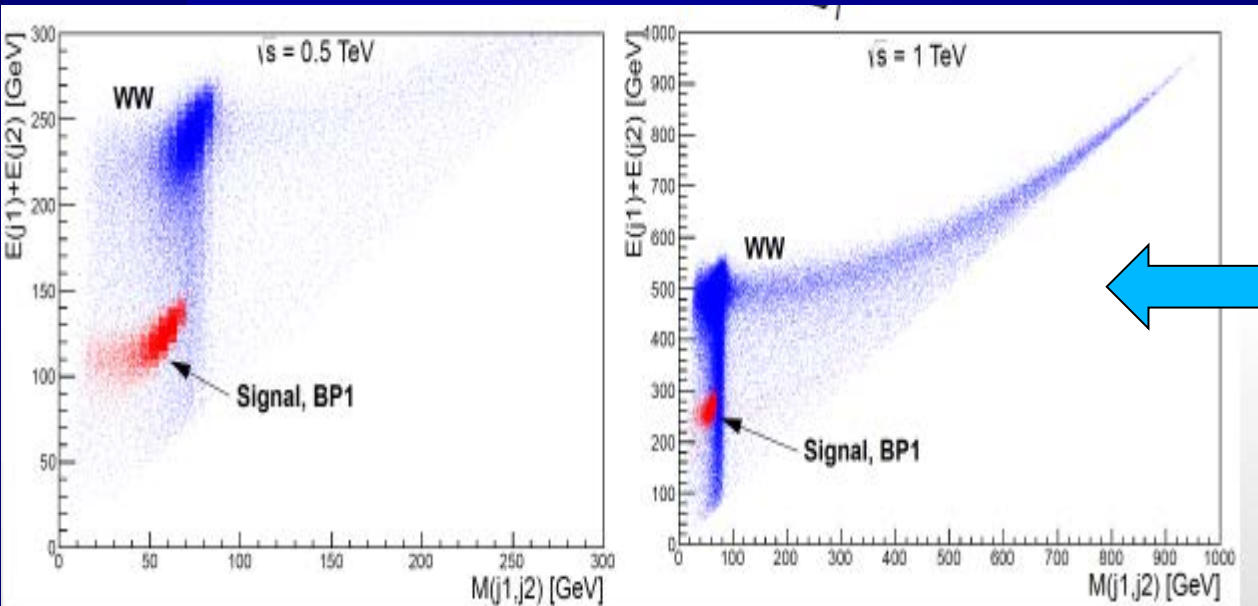
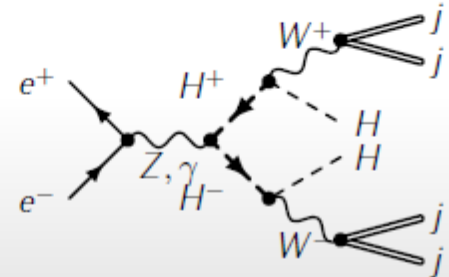
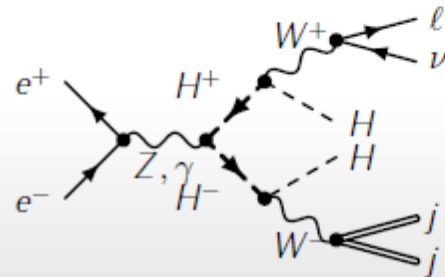
Also .. M. Aoki ... PL B725 (2013) 302

M. Krawczyk, Lisbon 2016

Idea I. Ginzburg J.Mod.Phys. 5 (2014) 1036

H+H-

Process	Signal			Background			
	BP1	BP2	BP3	WW	ZZ	Z+jets	$t\bar{t}$
σ [fb] @ 500 GeV	82.2	70.9	44.6	7807	583	16790	595
σ [fb] @ 1 TeV	28.1	27.3	25.3	3180	233	4304	212

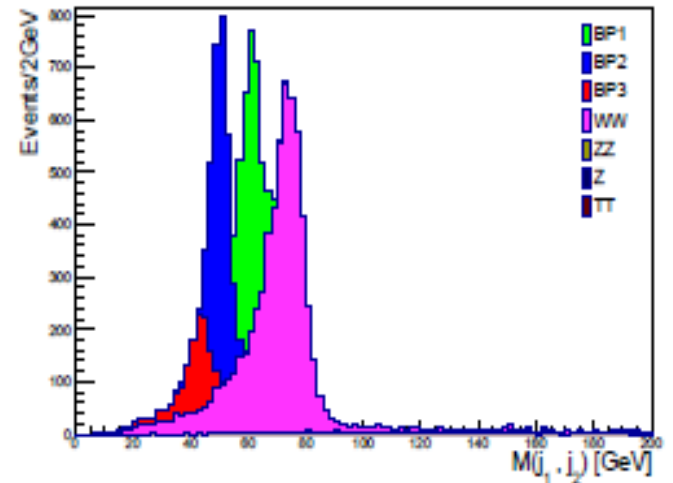
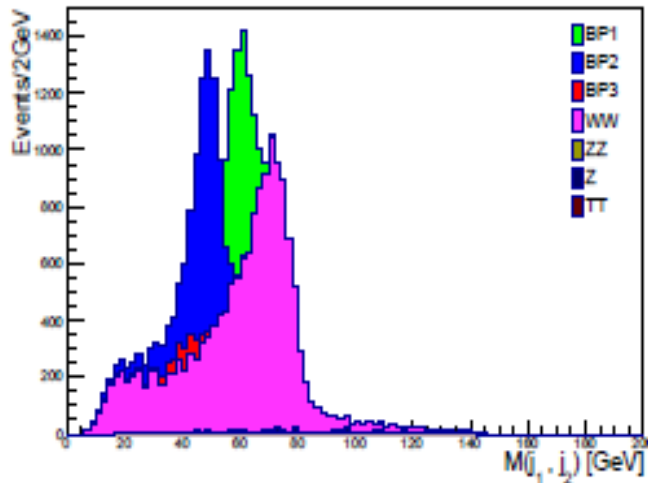
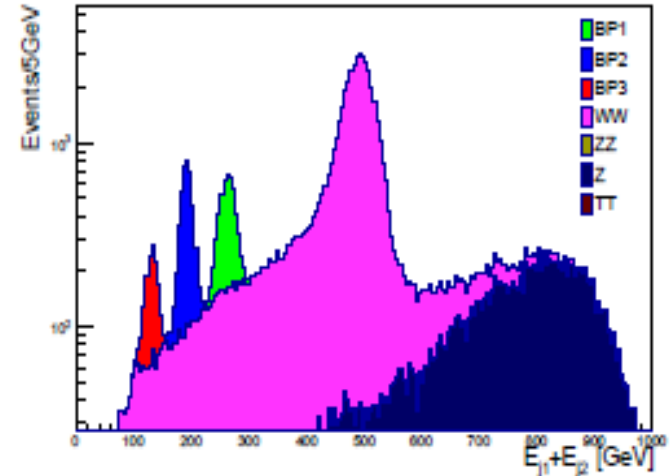
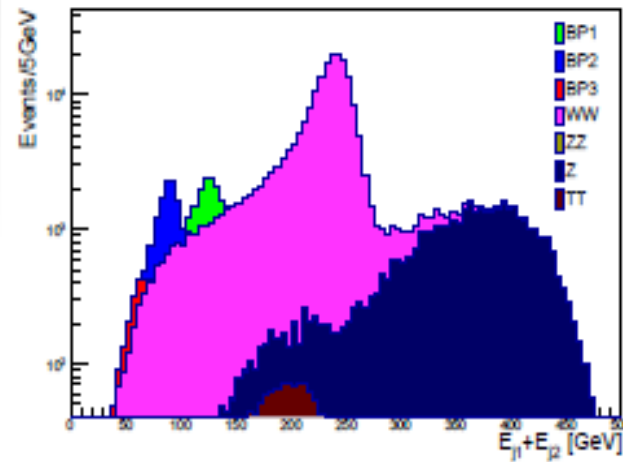


Sum of energies of two jets and their invariant mass for semileptonic events
 → cuts on sum of energies

background WW dominated by on-shell W

H+H-

Sum of the energies (up) and invariant mass (down) of two jets in semileptonic final state at $\sqrt{s} = 0.5$ TeV (left) and 1 TeV (right).

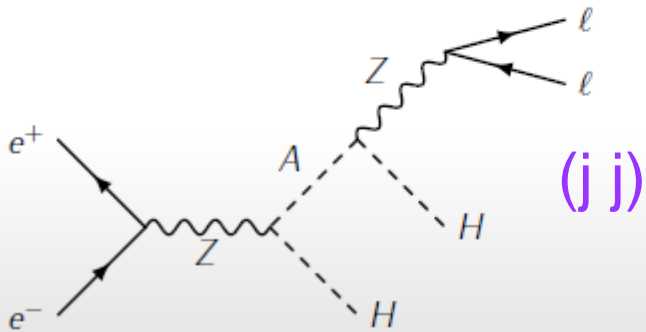


after
applying
cuts on
sum of
energies

150(350) GeV

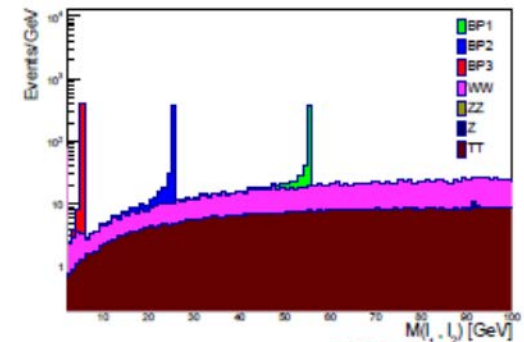
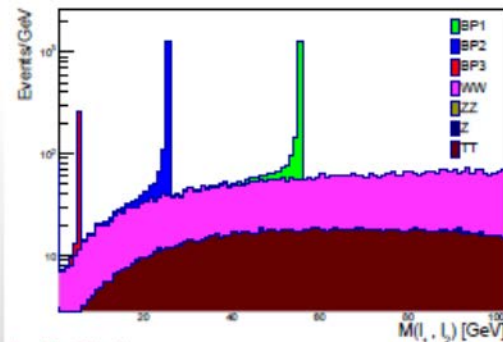
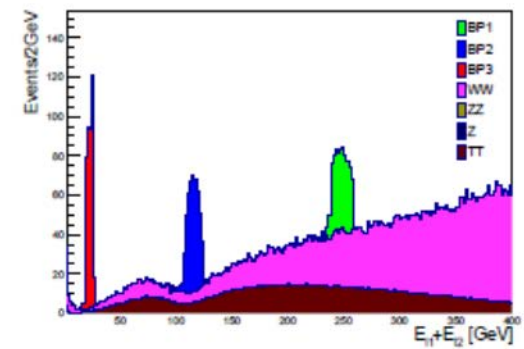
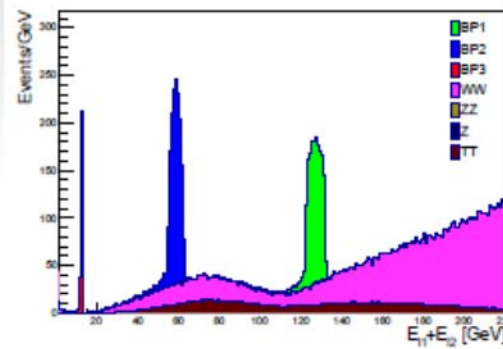
HA

	$\sqrt{s} = 0.5 \text{ TeV}$			$\sqrt{s} = 1 \text{ TeV}$		
Process	$e^+e^- \rightarrow AH$			$e^+e^- \rightarrow AH$		
Benchmark point	BP1	BP2	BP3	BP1	BP2	BP3
Cross section [fb]	45	42.9	34.2	12.5	12.4	11.8



Saereh Najjari

Sum of the energies (up) and invariant mass (down) of two lepton in leptonic final state at $\sqrt{s} = 0.5 \text{ TeV}$ (left) and 1 TeV (right).



Saereh Najjari

IDM @ Linear Coll

IDMS Bonilla, Diaz-Cruz, Darvishi, Sokolowska, MK – J.Phys. G43 (2016) 065001

- IDM + extra neutral complex singlet χ with complex vev
- SM-like doublet - singlet interaction \rightarrow CP violation in the scalar sector containing 3 neutral Higgses: $h1$ (SM-like), $h2$, $h3$
- Dark doublet as before $\rightarrow H$ is good DM candidate
- Small modifications of $h1$ couplings to SM particles

Potential IDMS

Φ_S

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v + \phi_1 + i\phi_6) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(\phi_4 + i\phi_5) \end{pmatrix},$$

Φ_D

$$\chi = \frac{1}{\sqrt{2}}(we^{i\xi} + \phi_2 + i\phi_3).$$

Z_2 : $\Phi_1 \rightarrow \Phi_1$, $\Phi_2 \rightarrow -\Phi_2$, SM fields \rightarrow SM fields, $\chi \rightarrow \chi$.

$$V = -\frac{1}{2} \left[m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 \right] + \frac{1}{2} \left[\lambda_1 \left(\Phi_1^\dagger \Phi_1 \right)^2 + \lambda_2 \left(\Phi_2^\dagger \Phi_2 \right)^2 \right]$$

$$+ \lambda_3 \left(\Phi_1^\dagger \Phi_1 \right) \left(\Phi_2^\dagger \Phi_2 \right) + \lambda_4 \left(\Phi_1^\dagger \Phi_2 \right) \left(\Phi_2^\dagger \Phi_1 \right) + \frac{\lambda_5}{2} \left[\left(\Phi_1^\dagger \Phi_2 \right)^2 + \left(\Phi_2^\dagger \Phi_1 \right)^2 \right]$$

$$- \frac{m_3^2}{2} \chi^* \chi + \lambda_{s1} (\chi^* \chi)^2 + \Lambda_1 (\Phi_1^\dagger \Phi_1) (\chi^* \chi)$$

$$- \frac{m_4^2}{2} (\chi^{*2} + \chi^2) + \kappa_2 (\chi^3 + \chi^{*3}) + \kappa_3 [\chi (\chi^* \chi) + \chi^* (\chi^* \chi)].$$

IDM

with softly broken U(1)

$U(1)$: $\Phi_1 \rightarrow \Phi_1$, $\Phi_2 \rightarrow \Phi_2$, $\chi \rightarrow e^{i\alpha} \chi$.

Higgs sector – $h_1(125 \text{ GeV}), h_2, h_3$

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}$$

$$R = R_1 R_2 R_3 = \begin{pmatrix} c_1 c_2 & c_3 s_1 - c_1 s_2 s_3 & c_1 c_3 s_2 + s_1 s_3 \\ -c_2 s_1 & c_1 c_3 + s_1 s_2 s_3 & -c_3 s_1 s_2 + c_1 s_3 \\ -s_2 & -c_2 s_3 & c_2 c_3 \end{pmatrix}$$

$$h_1 = c_1 c_2 \phi_1 + (c_3 s_1 - c_1 s_2 s_3) \phi_2 + (c_1 c_3 s_2 + s_1 s_3) \phi_3,$$

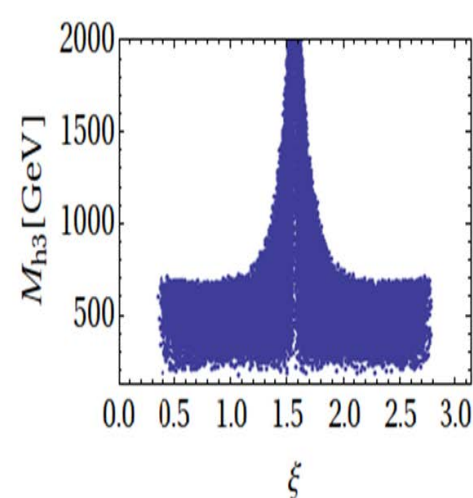
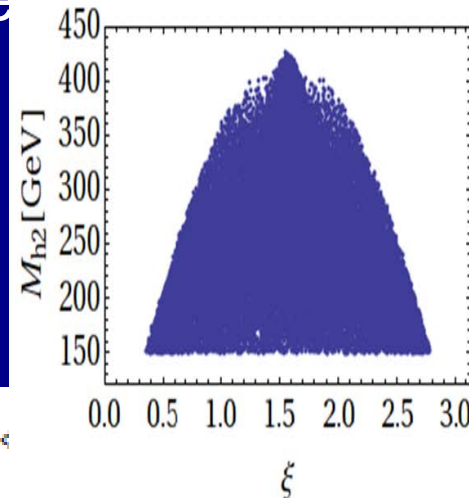
$$R_{11} = R_{11}^{-1} = c_1 c_2 \sim 1$$

$M_{h_1} \sim 125 \text{ GeV}$ $w=300-1000 \text{ GeV}$

$$M_{h_3} > M_{h_2} > 150 \text{ GeV.}$$

$$\kappa_{2,3} = w \rho_{2,3},$$

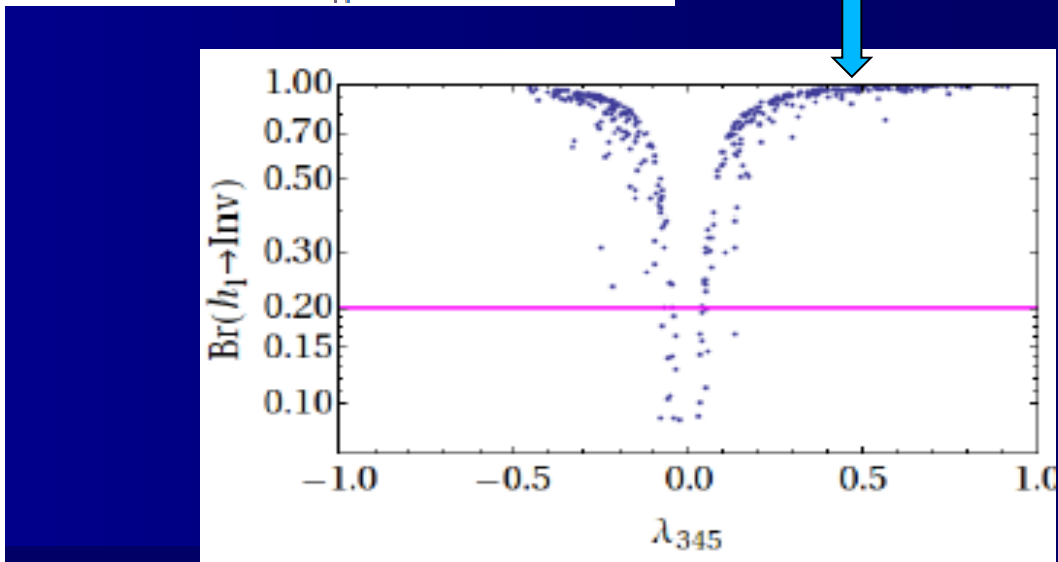
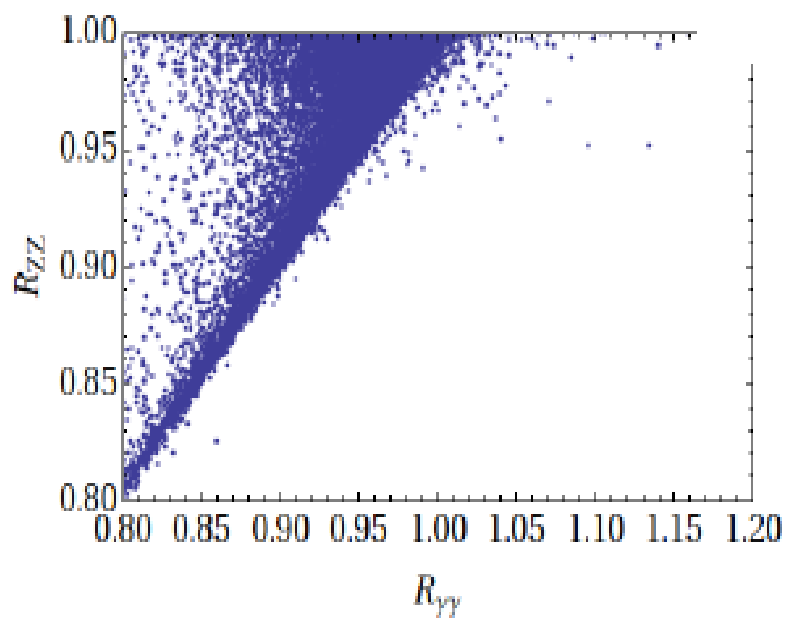
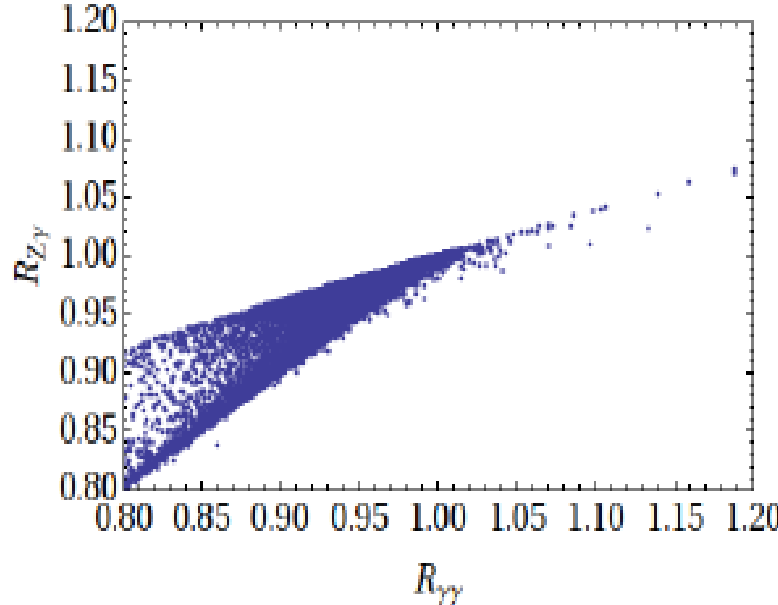
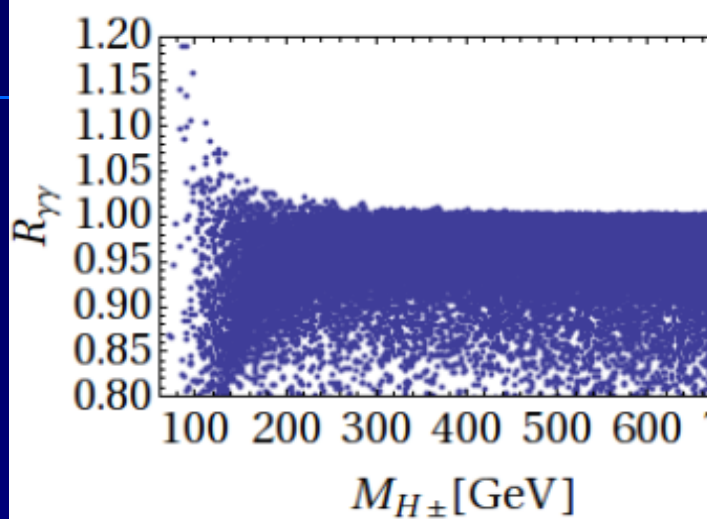
$$-1 < \Lambda_1 < 1, \quad 0 < \lambda_{s1} < 1, \quad -1 < \rho_{2,3} < 1, \quad 0 < \xi < 3$$



IDMS - $h_1(125 \text{ GeV})$

$$\Gamma(h_1 \rightarrow XX) = R_{11}^2 \Gamma(\phi_{SM} \rightarrow XX)$$

$XX = gg, VV^*$



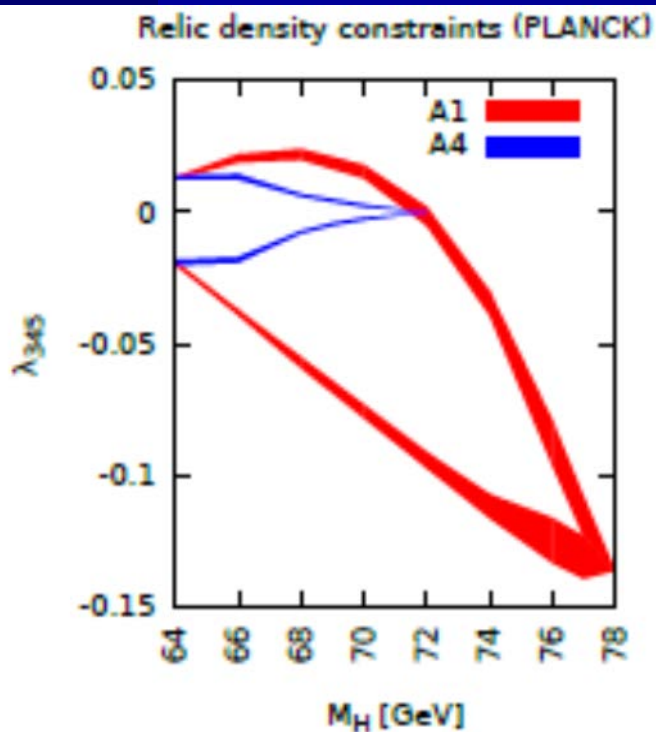
Br $h_1 \rightarrow \text{inv}$

Dark sector – Higgs portals via h_1, h_2, h_3 possible !

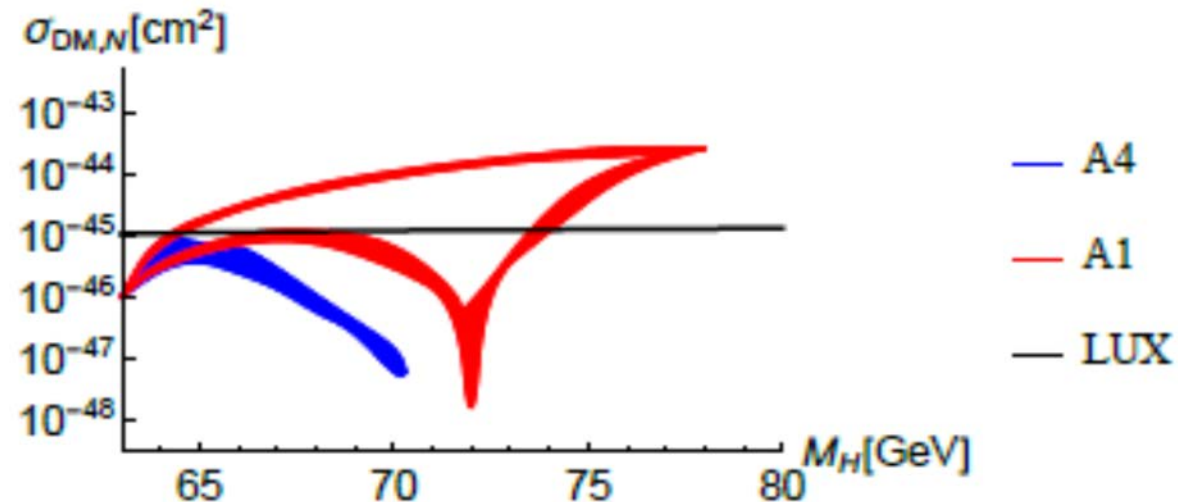
We proposed benchmarks – eg.

A1: $M_{h_1} = 124.83\text{GeV}$, $M_{h_2} = 194.46\text{GeV}$, $M_{h_3} = 239.99\text{GeV}$,

A4: $M_{h_1} = 125.36\text{GeV}$, $M_{h_2} = 149.89\text{GeV}$, $M_{h_3} = 473.95\text{GeV}$.



for A4 – Higgs portal via h_2



Summary

- SM-like scenario – still valid (Sept. 2016)
- IDM is a very natural extension of the SM
 - SM doublet → one Higgs SM-like h
 - Dark doublet → 4 scalars (two charged)
one stable ($H=DM$)
- IDM in agreement with LHC data and relic density + direct detection LUX data,
 $M_H > 45$ GeV
- Tests at LHC II , ILC/CLIC ?
- IDMS better ?
- Higgs is a sensitive probe of DM !

SM+complex singlet

Darvishi, Sokolowska, MK

1512.06437 (APP B47 2016)

Darvishi, MK 1603.00598

- Model
- Spontaneous CP violation
- Strong 1st order phase transition

- Baryogenesis with vector-like quarks (iso-doublet)

Darvishi, 1608.02820

CP violation in the Standard Model with a complex singlet

Constrained Potential
SM+CS \rightarrow cSMCS

N, Darvishi, MK

$$V = -\frac{1}{2}m_{11}^2\Phi^\dagger\Phi + \frac{1}{2}\lambda(\Phi^\dagger\Phi)^2 - \frac{m_s^2}{2}\chi^*\chi + \lambda_s(\chi^*\chi)^2 + \Lambda(\Phi^\dagger\Phi)(\chi^*\chi)$$

$$-\frac{m_4^2}{2}(\chi^{*2} + \chi^2) + \kappa_2(\chi^3 + \chi^{*3}) + \kappa_3[\chi(\chi^*\chi) + \chi^*(\chi^*\chi)] + \kappa_4\Phi^\dagger\Phi(\chi + \chi)$$

- ▶ Parameters $\rightarrow m_{11}^2, m_s^2, m_4^2, \lambda_s, \lambda, \Lambda, \kappa_2, \kappa_3, \kappa_4$.
If real \rightarrow **No explicit CP violation**
- ▶ Vacua with spontaneous CP violation
 $\langle\Phi\rangle = \frac{1}{\sqrt{2}}v$ and $\langle\chi\rangle = \frac{1}{\sqrt{2}}we^{i\xi}$
- ▶ **Spontaneous CP violation** \rightarrow relevant parameters $m_4^2, \kappa_2, \kappa_3, \kappa_4$
- ▶ Positivity conditions

similar study
Branco ...98

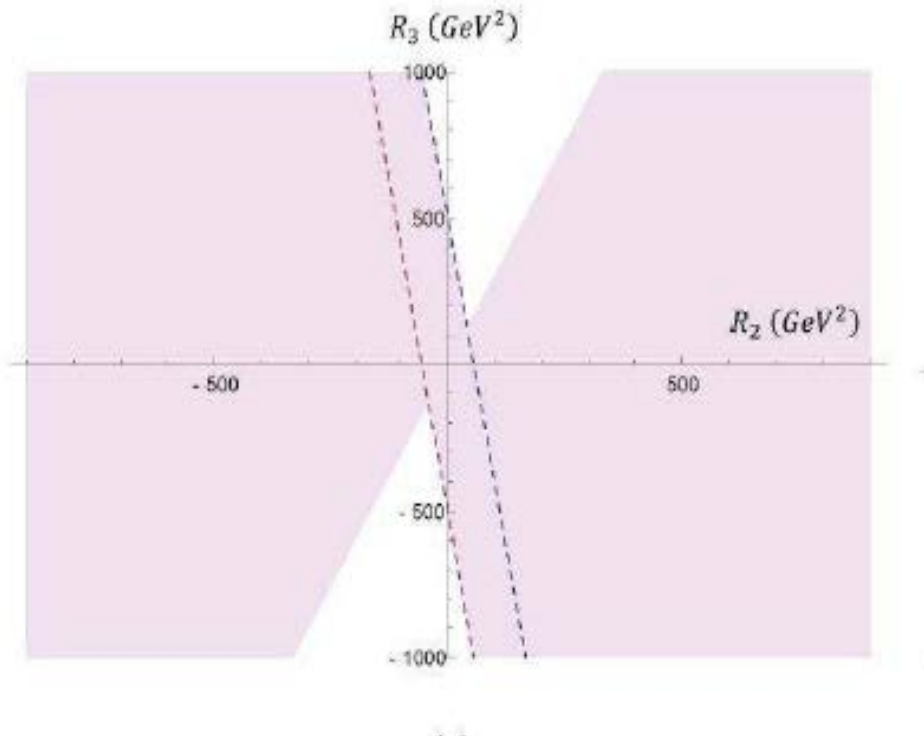
$$\lambda > 0, \quad \lambda_s > 0, \quad \Lambda + \sqrt{2\lambda\lambda_s} > 0.$$

Region of spont. CP violation

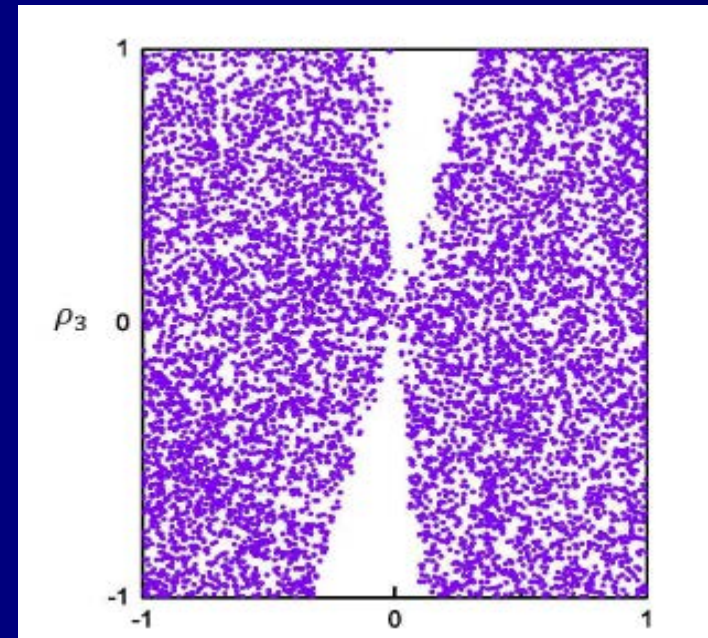
$$-4m_4^2 \cos \xi + 3R_2(1 + 2 \cos 2\xi) + R_3 + R_4 = 0,$$

$$R_2 = \sqrt{2}w\kappa_2, \quad R_3 = \sqrt{2}w\kappa_3, \quad R_4 = \frac{\sqrt{2}v^2\kappa_4}{w},$$

for $\kappa_4 = 0$



full scan



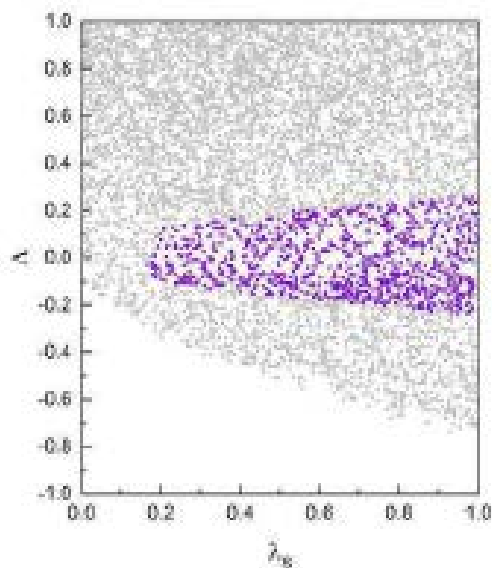
Scan – similar to IDMS but here low $\langle \chi \rangle \sim w$ possible

$$M_{h_1} \in [124.00, 127.00] \text{ GeV}, \quad M_{h_3} \gtrsim M_{h_2} > 150 \text{ GeV}$$

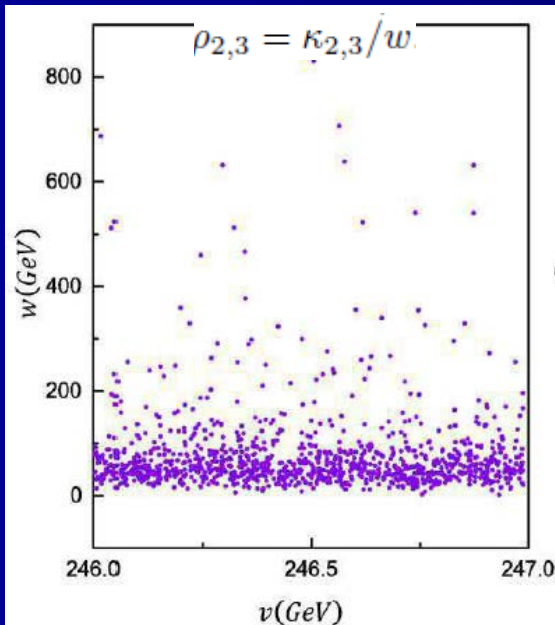
$$0.2 < \lambda_1 < 0.3.$$

$$-1 < \Lambda < 1, \quad 0 < \lambda_s < 1, \quad -1 < \rho_{2,3} < 1, \quad 0 < \xi < \pi,$$

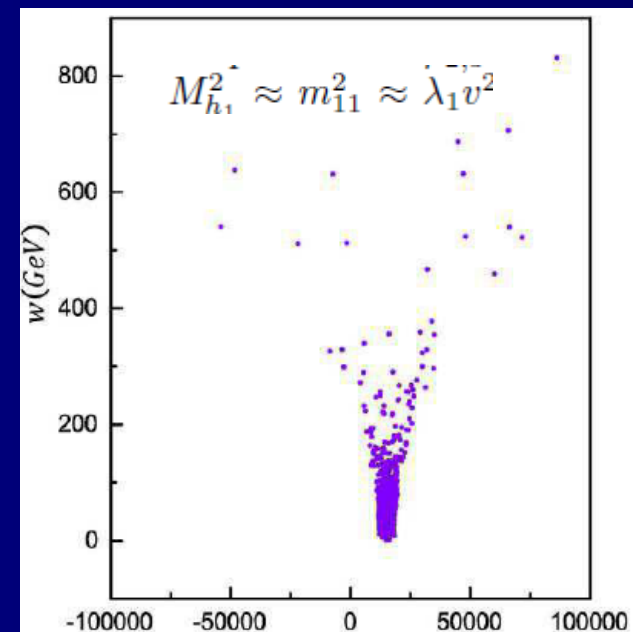
$$-90000 \text{ GeV}^2 < \mu_1^2, \mu_2^2, m_{11}^2 < 90000 \text{ GeV}^2. \quad \mu_1^2 = m_s^2 + 2m_4^2, \quad \mu_2^2 = m_s^2 - 2m_4^2.$$



(a)



(b)



J invariant

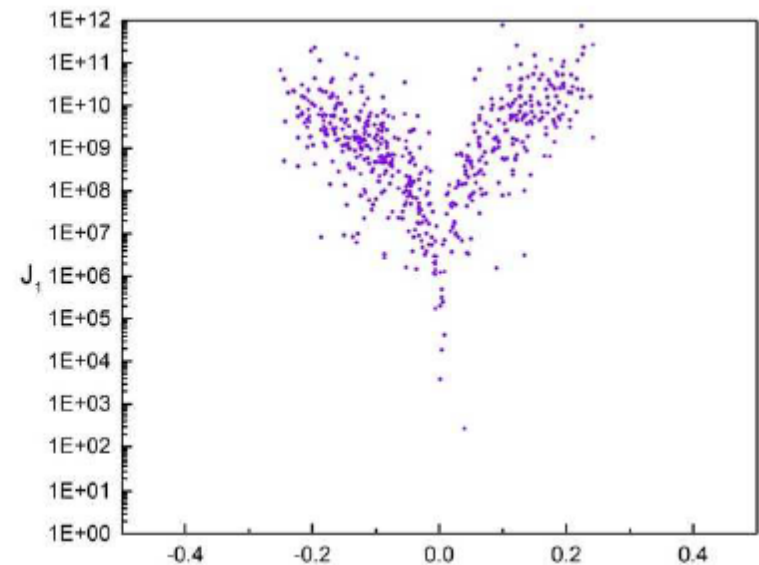
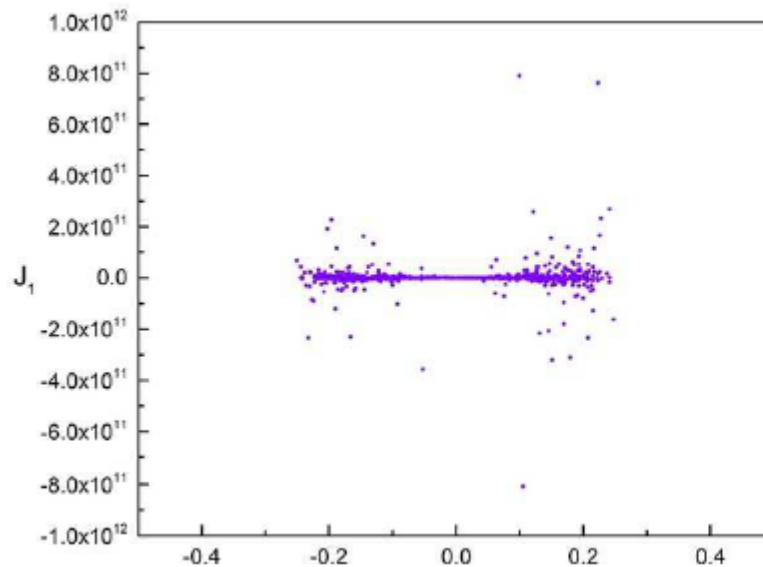
$$J_1 = M_{12}M_{13}(M_{22} - M_{33}) + M_{23}(M_{13}^2 - M_{12}^2),$$

$$J_1 = 2\Lambda^2 v^2 w^2 m_4^2 \sin \xi \cos \xi = 2\Lambda^2 v^2 w_1 w_2 m_4^2$$

where

$$-4m_4^2 \cos \xi + 3R_2(1 + 2 \cos 2\xi) + R_3 + R_4 = 0,$$

$$R_{2,3,4} \sim w \kappa_{2,3,4}$$

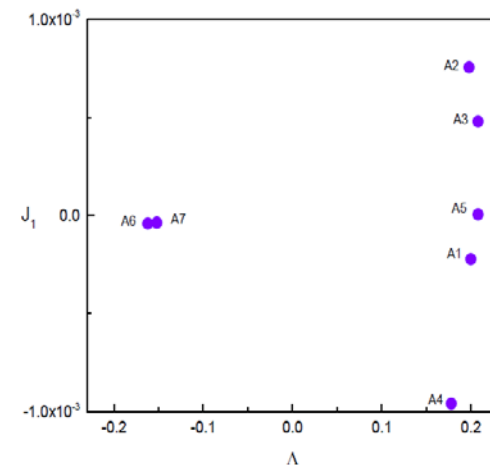


Benchmarks

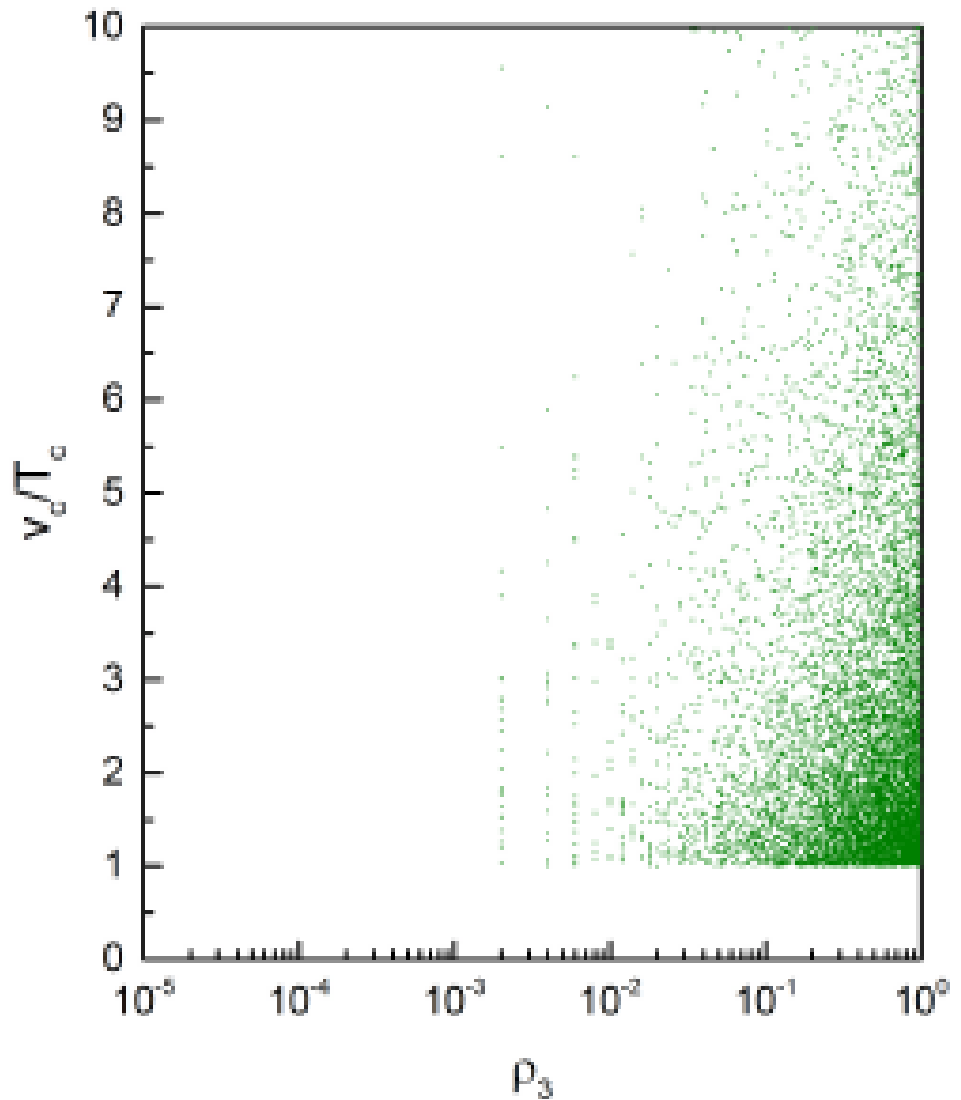
Benchmark	α_1	α_2	α_3	M_{h_1}	M_{h_2}	M_{h_3}	S	T	J_1/v^6
A1	-0.047	-0.053	1.294	124.64	652.375	759.984	-0.072	-0.094	-2.2×10^{-4}
A2	-0.048	0.084	0.084	124.26	512.511	712.407	-0.001	-0.039	7.2×10^{-4}
A3	0.078	0.297	0.364	124.27	582.895	650.531	0.003	-0.046	4.5×10^{-4}
A4	0.006	-0.276	0.188	125.86	466.439	568.059	-0.013	-0.169	-9.5×10^{-4}
A5	0.062	-0.436	0.808	125.21	303.545	582.496	0.002	-0.409	5.0×10^{-6}
A6	-0.210	0.358	0.056	124.92	181.032	188.82	0.003	-0.010	-4.0×10^{-5}
A7	-0.205	0.403	0.057	125.01	175.45	178.52	0.002	-0.020	-3.5×10^{-5}

Table I. Benchmark points A1 – A7, masses are given in GeV

Benchmark	$R_{\gamma\gamma}^{h_1}$	$R_{\gamma\gamma}^{h_2}$	$R_{\gamma\gamma}^{h_3}$	$\Gamma_{tot}^{h_1}$	$\Gamma_{tot}^{h_2}$	$\Gamma_{tot}^{h_3}$
A1	0.98	0.0021	0.0028	0.0042	0.304	0.781
A2	0.98	0.0021	0.0070	0.0042	0.145	1.31
A3	0.98	0.0055	0.085	0.0042	0.566	12.24
A4	0.92	3.3×10^{-5}	0.074	0.0043	0.001	7.08
A5	0.81	0.0029	0.17	0.0043	0.002	17.51
A6	0.82	0.19	0.11	0.0043	0.119	0.163
A7	0.81	0.18	0.15	0.0043	0.871	0.083



Strong 1st order PT



cubic term

Baryogenesis with iso-doublet vector-like quarks

Neda Darvishi

1608.02820

Iso-doublet vector quarks as in

McDonald, Phys. Rev. D 53 (1996) 645.

28 Nov-3 Dec 2016
Warsaw, Poland



TOPICS

T, C, P, CP and CPT symmetries

Emergence of symmetries from entanglement

Discrete symmetries in cosmology

Neutrino masses, mixing

and discrete symmetries Role of symmetries
beyond the SM

Discrete symmetries and models of flavour mixing

Strings,
branes,

New results from LHC,
new facilities

extra dimensions,
discrete symmetries

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Photo by Andrzej Jórczak (1944-81), astronomer and artist
specializing in experimental photography, from the series *Commutations* (1972).

Courtesy of Liliana Jórczak Sami


Evolution of Universe to the Inert Phase

Evolution of the Universe in 2HDM— through different vacua in the past

Ginzburg, Ivanov, Kanishev 2009

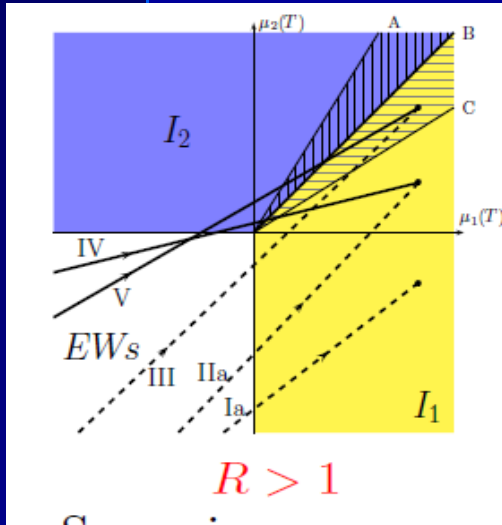
*Ginzburg, Kanishev, MK, Sokołowska PRD 2010,
Sokołowska 2011*

We consider 2HDM with an explicit D symmetry assuming that today the **Inert Doublet Model** describes reality. In the simplest approximation only *mass terms* in V vary with temperature like T^2 , while λ 's are fixed

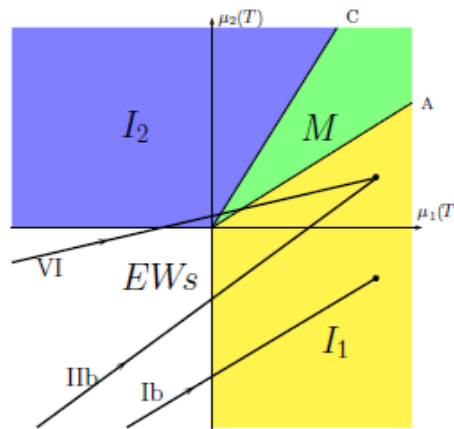
Various evolution from EWs to Inert phase possible in one, two or three steps, 
with 1st or 2nd order phase transitions...

Evolution of vacua

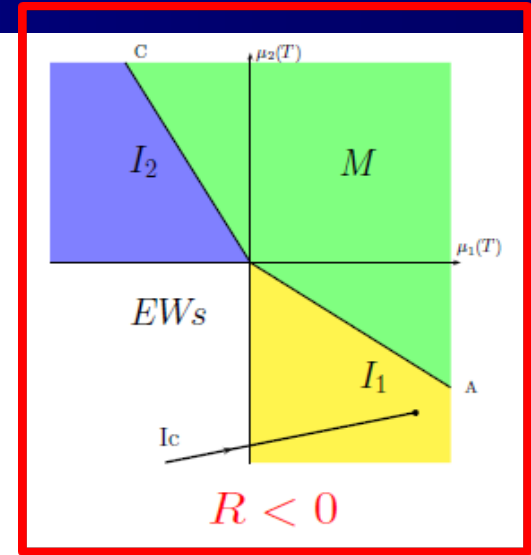
$EWs \rightarrow I_2 \rightarrow I_1$



$R > 1$



$1 > R > 0$



$R < 0$

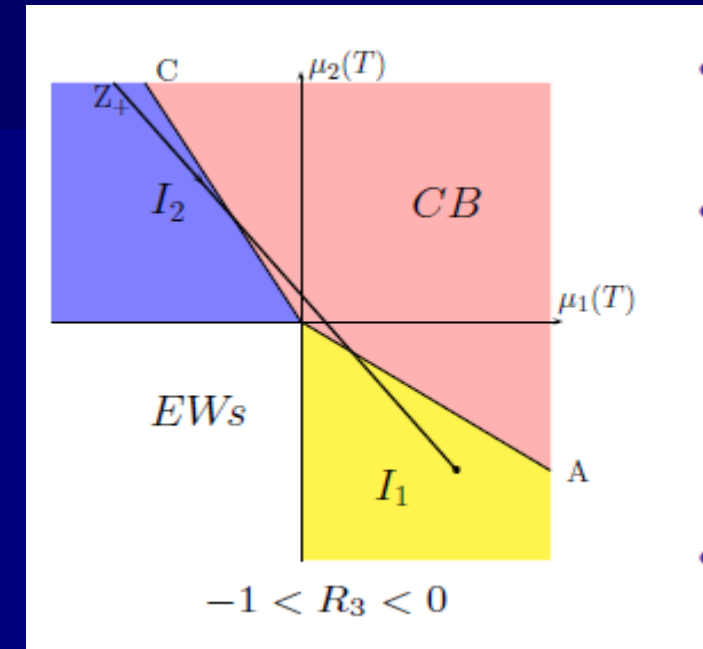
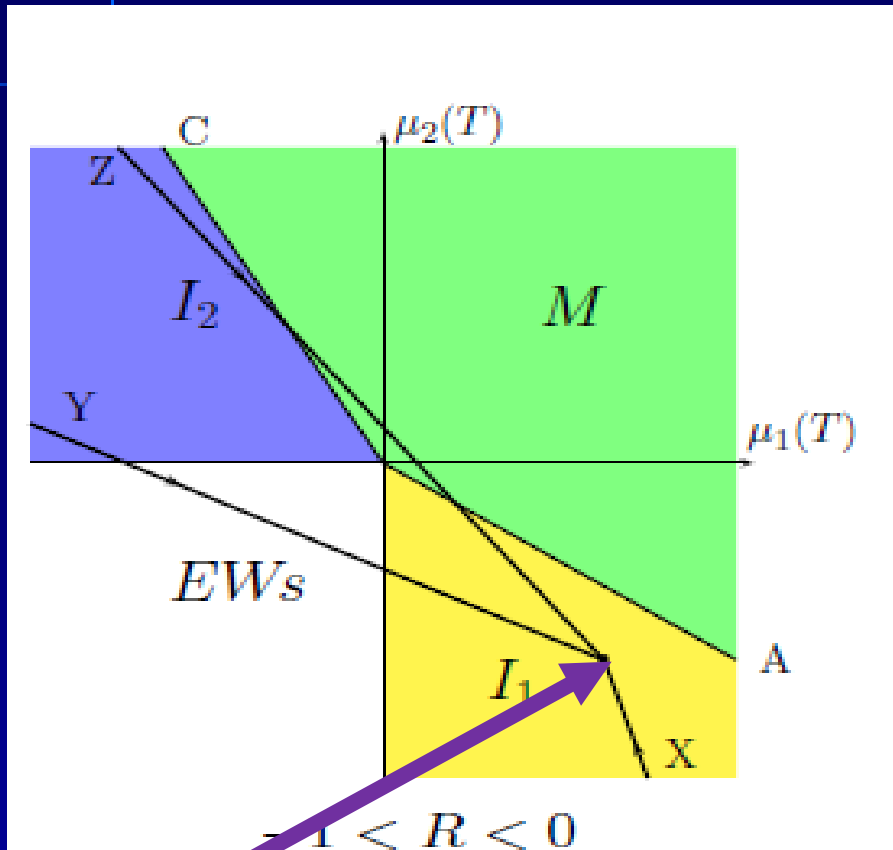
$EWs \rightarrow I_1$

T^2 corrections

→ rays from EWs phase to Inert phase
 one, two or three stages of Universe
 (II order phase transitions, one I order)

$$R = \frac{\lambda_{345}}{\sqrt{\lambda_1 \lambda_2}}$$

Nonrestoration of EW symmetry for $\lambda_{345} < 0$



Charged breaking
phase

Only one ray with EW restoration in the past
(in one step)

Beyond T^2 corrections \gg strong 1st order phase transition in IDM

EW baryogenesis?

G. Gil MsThesis'2011, G. Gil, P. Chankowski, MK 1207.0084 [hep-ph] PLB 2012

We applied one-loop effective potential at $T=0$ (Coleman-Wienberg term) and temperature dependent effective potential at $T \neq 0$ (with sum of ring diagrams)

$$V_T^{(1L)}(v_1, v_2) = V_{\text{eff}}^{(1L)}(v_1, v_2) + \Delta^{(1L)} V_{T \neq 0}(v_1, v_2).$$

The one-loop effective potential $V_{\text{eff}}(v_1, v_2)$ is given in the Landau gauge by standard formula

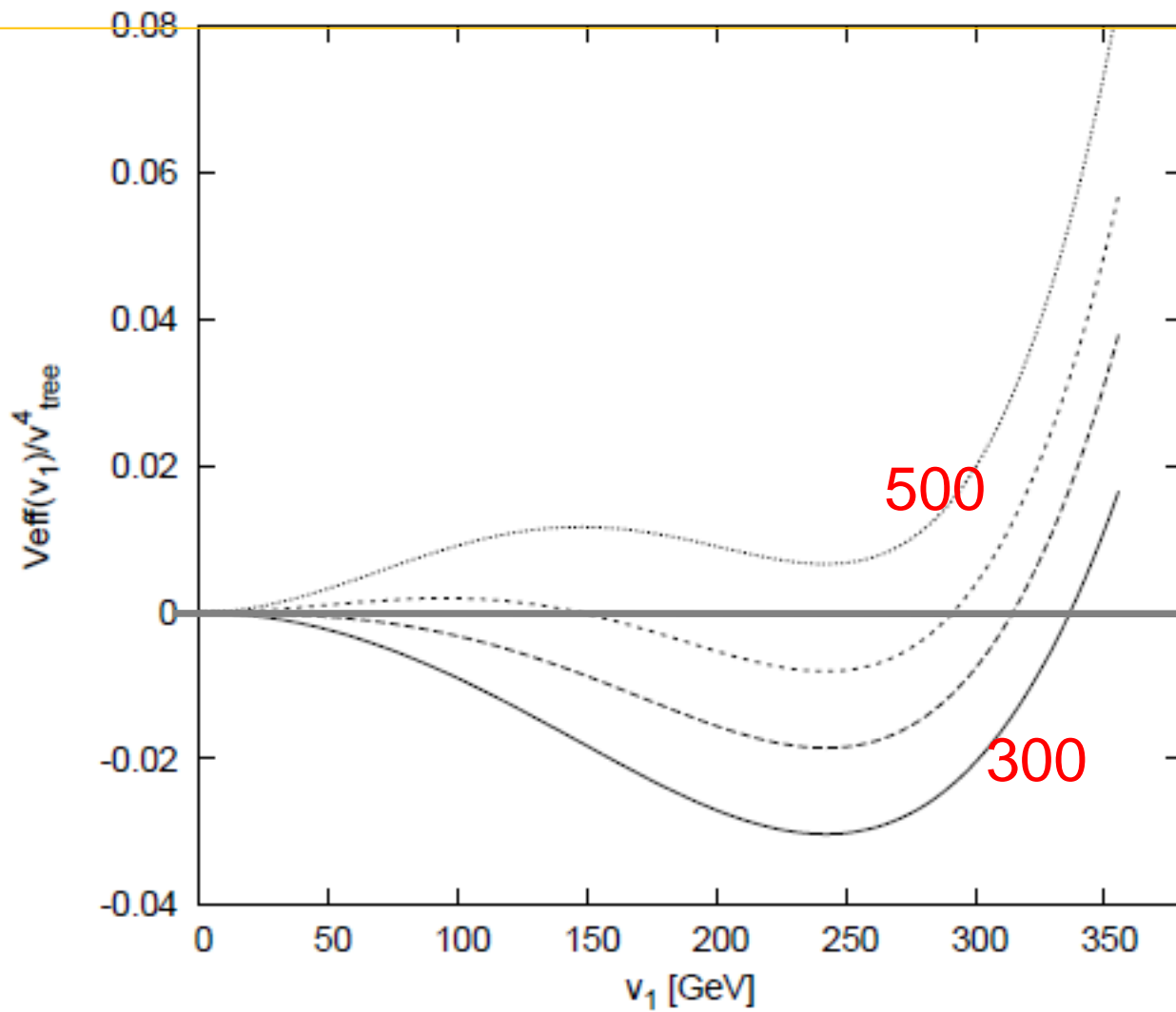
$$V_{\text{eff}}^{(1L)} = V_{\text{tree}} + \frac{1}{64\pi^2} \sum_{\text{fields}} C_s \left\{ \mathcal{M}_s^4 \left(\ln \frac{\mathcal{M}_s^2}{4\pi\mu^2} - \frac{3}{2} + \frac{2}{d-2} - \gamma_E \right) \right\} + \text{CT},$$

mass matrices

number of states

counter terms \rightarrow

Effective $T=0$ potential



$M_h = 125$ GeV

$M_H = 65$ GeV

$M_H + M_A =$
500, 450, 400, 300
GeV

$\lambda_{345} = 0.2,$
 $\lambda_2 = 0.2$

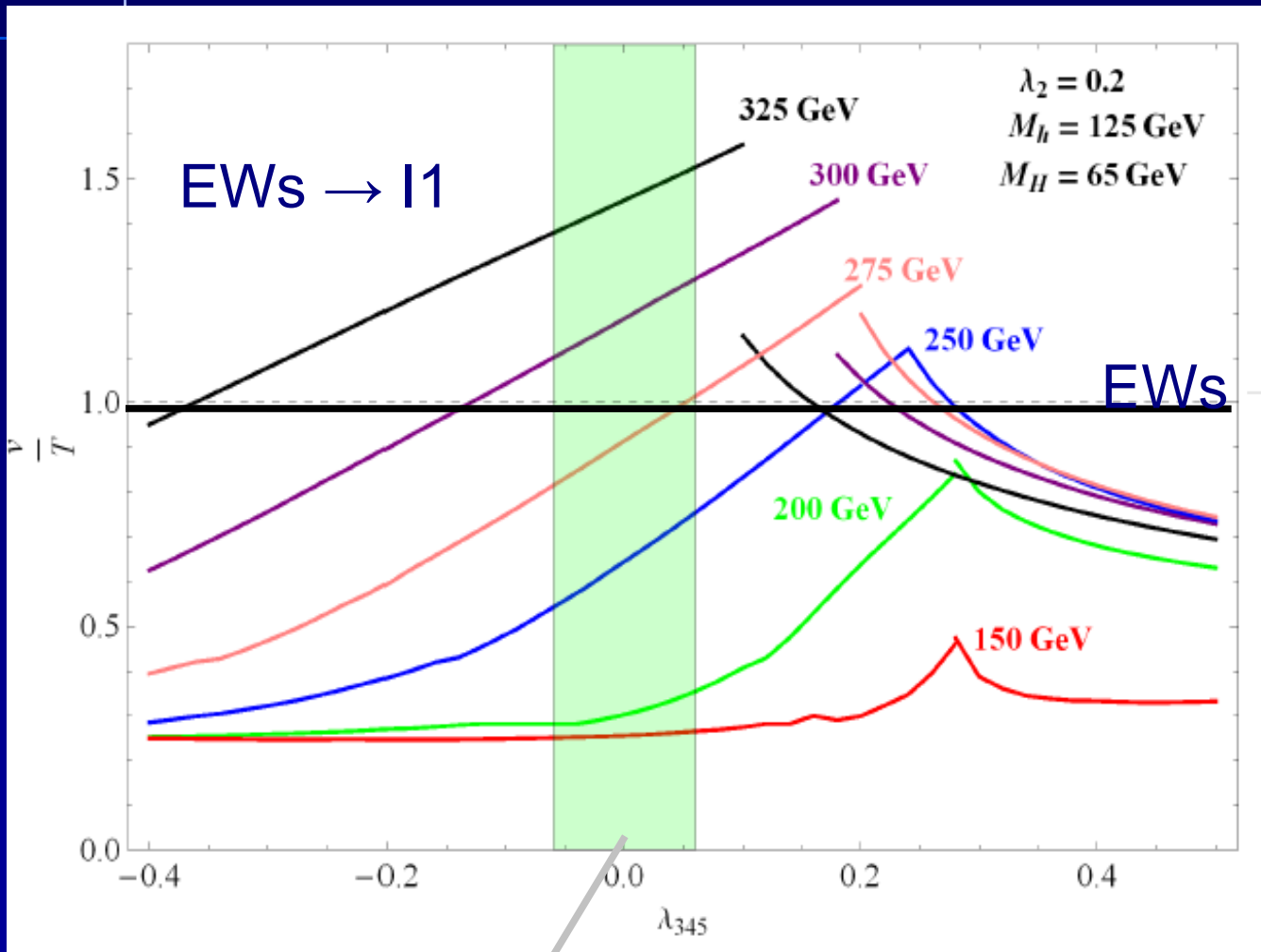
$v_{2(D)} = 0$

Critical temperature T_{EW} : V at new minimum = V at $(v_{1(s)} = v_{2(D)} = 0)$ ₅₁

Results for $v(T_{EW})/T_{EW} > 1$

$M_h=125$ GeV, $M_H=65$ GeV, $\lambda_2=0.2$

strong 1st order
phase transition
if ratio > 1



Allowed
MH+=MA
between 275
and 380 GeV
(one step)

R<0 Xenon100 bound R>0

λ_{345}