

Neutrino Mass Generation with Multi-Higgs

Hiroaki SUGIYAMA
(Toyama Prefectural Univ.)

Based on "S. Kanemura, HS, PLB753, 161"
"S. Kanemura, K. Sakurai, HS, PLB758, 465"
"M. Aoki, S. Kanemura, K. Sakurai, HS, PLB763, 352"

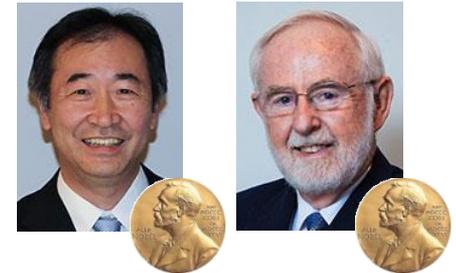
Introduction

Neutrino mass

- { The standard model (SM) \longrightarrow Massless
 { Neutrino oscillation \longrightarrow **Non-zero masses**



Two possible neutrino masses



Dirac mass : $m_D \bar{\nu}_L \nu_R$ (introduce ν_R)

$$y_\nu \bar{L} \epsilon \Phi^* \nu_R \implies m_D = \frac{y_\nu}{\sqrt{2}} v$$

$$(m_D \sim 0.1 \text{ eV} \implies y_\nu \sim 10^{-12})$$

Unnaturally small

Nontrivial way ?



Majorana mass : $m_L \bar{\nu}_L (\nu_L)^c$ (Lepton # violation)

$$Q_{\text{EM}} : 0 + 0 = 0$$

Specific to neutrinos \longrightarrow **How to generate ?**

Many models for m_ν \longrightarrow Which is the true one ?

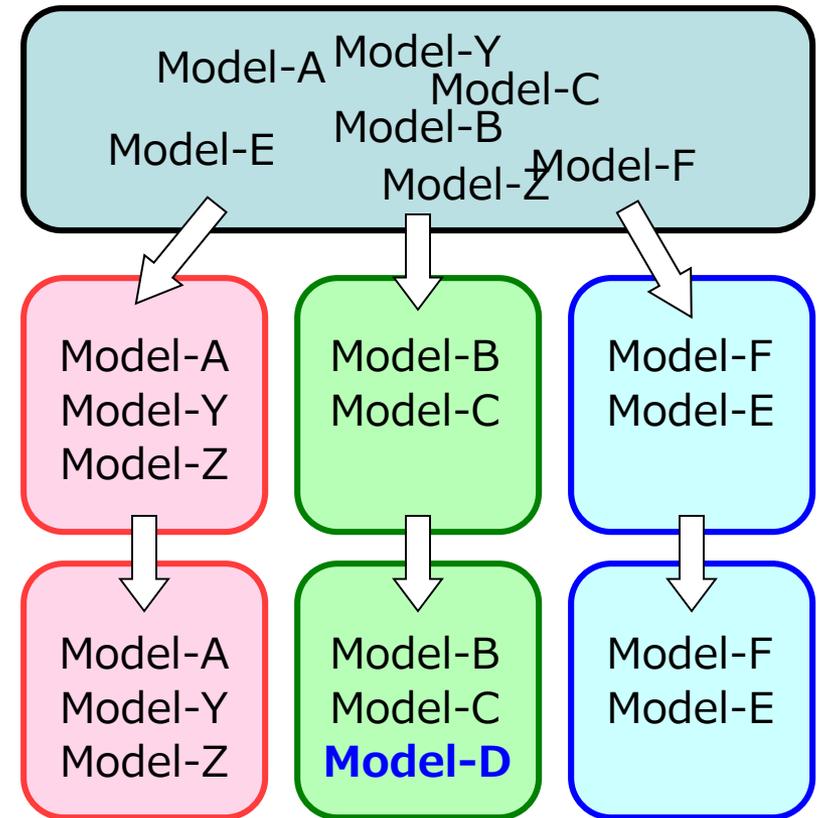
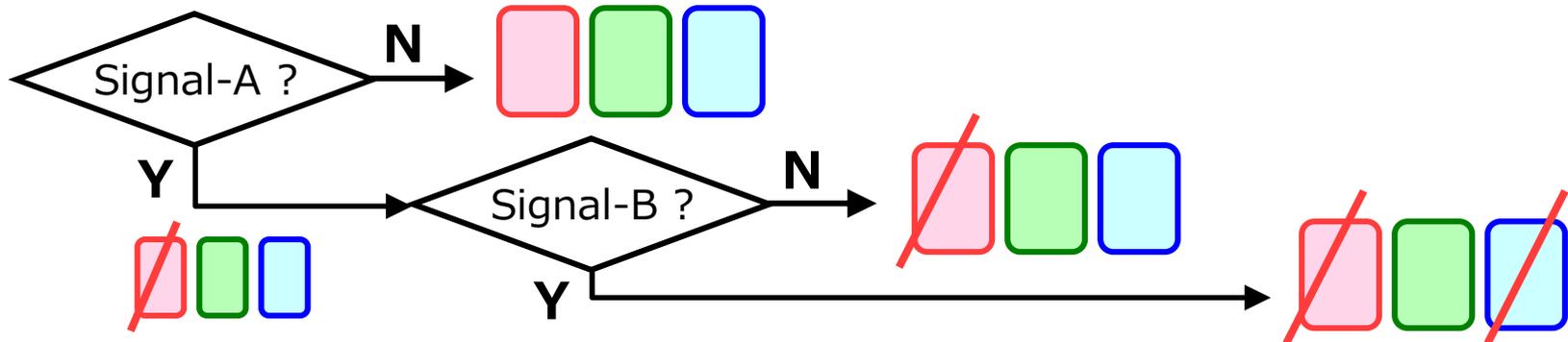
Motivation

There are (too) many models.

Classification is desired.

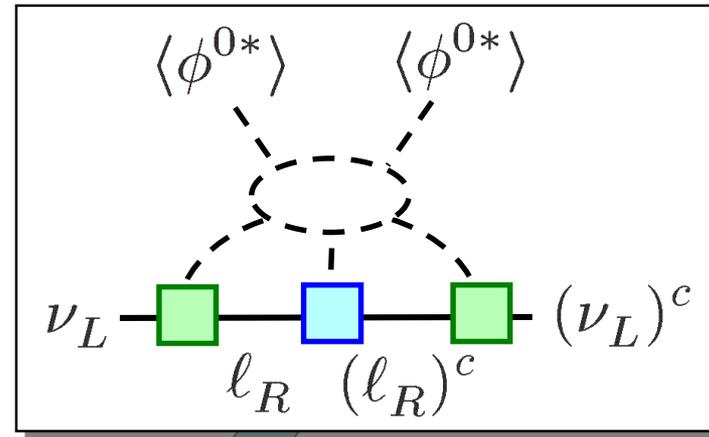
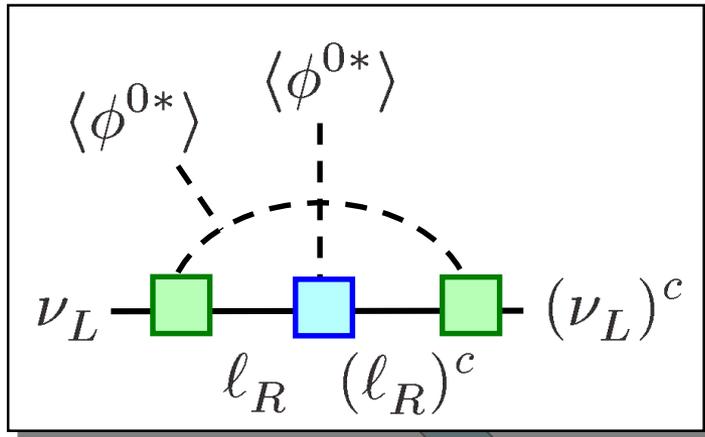
Missing models ?

Efficient tests ?

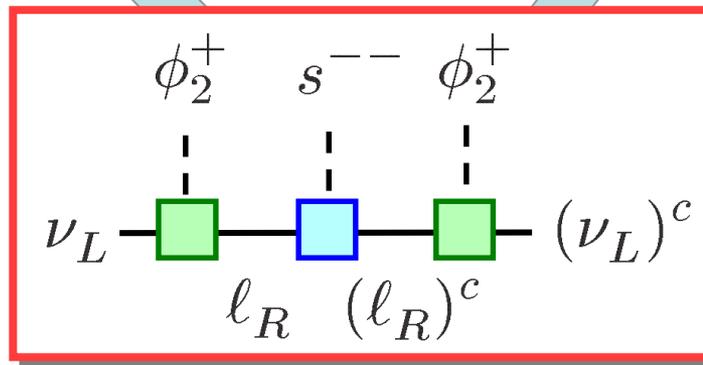


Classification according to Yukawa Int.

● Models (full Lagrangian)



● “Mechanisms” : Concentrating on **Yukawa int.** of leptons



Common structure of neutrino mass matrix
(cf. Overall scale depends on detail of models)

Contents



- Introduction
- **Classification of Models**
- Experimental Tests
- Summary

Classification of Models

Set up for "simple" models

Symmetries : $\left\{ \begin{array}{l} \text{SM gauge sym.} + \text{Unbroken } Z_2 \quad (\text{for dark matter}) \\ Z'_2 \text{ (Softly broken } Z_2 \text{) for Dirac neutrino} \\ \text{to forbid } y_\nu \bar{L} \epsilon \Phi^* \nu_R \text{ at the tree level} \\ \text{Another softly broken } Z_2 \text{ for 2HD cases} \\ \text{to forbid FCNC at the tree level} \end{array} \right.$

New fermions : $\left\{ \begin{array}{l} \nu_R : \text{Only for Dirac neutrino, } Z'_2\text{-odd} \\ \psi_R^0 : \text{Gauge-singlet, } Z_2\text{-odd} \end{array} \right.$

New scalars : $\left\{ \begin{array}{l} \text{Yukawa int. with leptons} \\ \left\{ \begin{array}{l} [\text{New scalar}] - [\text{Lepton}] - [\text{Lepton}] \\ [Z_2\text{-odd scalar}] - \psi_R^0 - [\text{Lepton}] \end{array} \right. \\ \text{We do not take cases without new scalars} \\ \text{(e.g. type-I seesaw)} \end{array} \right.$

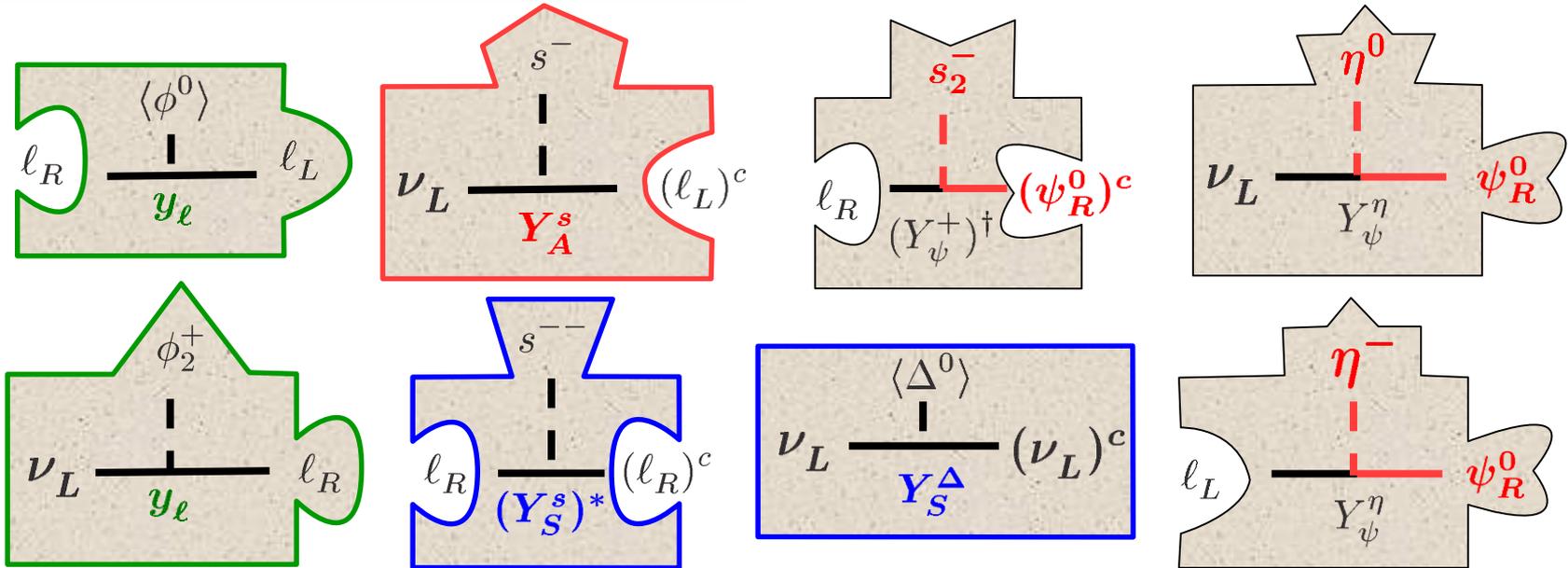
Scalars with Leptonic Yukawa Int.



Majorana

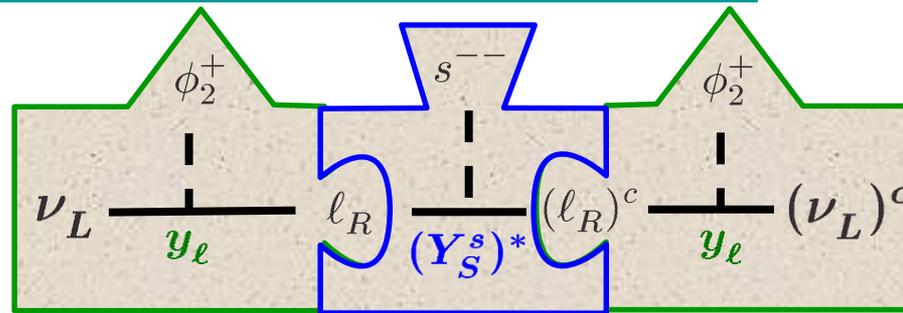
	SU(2) _L	U(1) _Y	Yukawa int.	Note
s_L^+	<u>1</u>	1	$(Y_A^s)_{\ell\ell'} [\overline{L}_\ell \epsilon L_{\ell'}^c s_L^-]$	Antisym.
s^{++}	<u>1</u>	2	$(Y_S^s)_{\ell\ell'} [(\ell_R)^c \ell'_R s^{++}]$	Sym.
$\Phi_2 = \begin{pmatrix} \phi_2^+ & \phi_2^0 \end{pmatrix}^T$	<u>2</u>	1/2	$y_\ell [\overline{L}_\ell \Phi_2 \ell_R]$	Diag.
$\Delta = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}$	<u>3</u>	1	$(Y_S^\Delta)_{\ell\ell'} [\overline{L}_\ell \Delta^\dagger \epsilon L_{\ell'}^c]$	Sym.
s_2^+ Z₂ odd	<u>1</u>	1	$(Y_\psi^+)_{\ell i} [(\ell_R)^c \psi_{iR}^0 s_2^+]$	ψ_R^0
$\eta = (\eta^+ \ \eta^0)^T$ Z₂ odd	<u>2</u>	1/2	$(Y_\psi^\eta)_{\ell i} [\overline{L}_\ell \epsilon \eta^* \psi_{iR}^0]$	ψ_R^0

Yukawa int. : Puzzle pieces



Neutrino mass : Solutions of the puzzle

For example,



Introduce s^{++} , $\Phi_2 \Rightarrow$ Majorana neutrino mass

$$m_L \propto \mathbf{y} (Y_S^s)^* \mathbf{y}$$

For Majorana ν masses

	Scalar with leptonic Yukawa int.					
					Z_2 -odd	
	s_L^+	s^{++}	Φ_2	Δ	s_2^+	η
$SU(2)_L$	<u>1</u>	<u>1</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>2</u>
$U(1)_Y$	1	2	1/2	1	1	1/2
Unbroken Z_2	+	+	+	+	-	-
	✓		✓			
M1	✓	✓				
M2		✓	✓			
M3		✓				
M4				✓		
M5	✓				✓	
M6			✓		✓	
M7					✓	
M8						✓

Zee-Wolfenstein model

Zee-Babu model

Cheng-Li model

Gustafsson-No
-Rivera model

Higgs Triplet model

Crauss-Nasri
-Trodden model

Aoki-Kanemura
-Seto model

New

Ma model

8 Mechanisms



Majorana

	SU(2) _L	U(1) _Y	Yukawa int.	Note
s_L^+	<u>1</u>	1	$(Y_A^s)_{\ell\ell'} [\overline{L}_\ell \epsilon L_{\ell'}^c s_L^-]$	Antisym.
s^{++}	<u>1</u>	2	$(Y_S^s)_{\ell\ell'} [(\overline{\ell}_R)^c \ell'_R s^{++}]$	Sym.
$\Phi_2 = \begin{pmatrix} \phi_2^+ & \phi_2^0 \end{pmatrix}^T$	<u>2</u>	1/2	$y_\ell [\overline{L}_\ell \Phi_2 \ell_R]$	Diag.
$\Delta = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}$	<u>3</u>	1	$(Y_S^\Delta)_{\ell\ell'} [\overline{L}_\ell \Delta^\dagger \epsilon L_{\ell'}^c]$	Sym.
$s_2^+ \quad Z_2 \text{ odd}$	<u>1</u>	1	$(Y_\psi^+)_{li} [(\overline{\ell}_R)^c \psi_{iR}^0 s_2^+]$	ψ_R^0
$\eta = (\eta^+ \quad \eta^0)^T \quad Z_2 \text{ odd}$	<u>2</u>	1/2	$(Y_\psi^\eta)_{li} [\overline{L}_\ell \epsilon \eta^* \psi_{iR}^0]$	ψ_R^0



Dirac

	SU(2) _L	U(1) _Y	Yukawa int.	Note
s^0	<u>1</u>	0	$(Y_S^0)_{ij} \left[\overline{(\nu_{iR})^c} \nu_{jR} s^0 \right]$	Sym. ν_R
s_L^+	<u>1</u>	1	$(Y_A^s)_{\ell\ell'} \left[\overline{L_\ell} \epsilon L_{\ell'}^c s_L^- \right]$	Antisym.
s_R^+	<u>1</u>	1	$(Y^s)_{li} \left[\overline{(\ell_R)^c} \nu_{iR} s_R^+ \right]$	ν_R
s^{++}	<u>1</u>	2	$(Y_S^s)_{\ell\ell'} \left[\overline{(\ell_R)^c} \ell'_R s^{++} \right]$	Sym.
$\Phi_\nu = \begin{pmatrix} \phi_\nu^+ & \phi_\nu^0 \end{pmatrix} T$	<u>2</u>	1/2	$(Y_\nu)_{li} \left[\overline{L_\ell} \epsilon \Phi_\nu^* \nu_{iR} \right]$	ν_R
$\Phi_2 = \begin{pmatrix} \phi_2^+ & \phi_2^0 \end{pmatrix} T$	<u>2</u>	1/2	$y_\ell \left[\overline{L_\ell} \Phi_2 \ell_R \right]$	Diag.
$\Delta = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}$	<u>3</u>	1	$(Y_S^\Delta)_{\ell\ell'} \left[\overline{L_\ell} \Delta^\dagger \epsilon L_{\ell'}^c \right]$	Sym.
s_2^0 Z₂ odd	<u>1</u>	0	$(Y_\psi^0)_{ij} \left[\overline{(\nu_{iR})^c} \psi_{jR}^0 s_2^0 \right]$	ν_R ψ_R^0
s_2^+ Z₂ odd	<u>1</u>	1	$(Y_\psi^+)_{li} \left[\overline{(\ell_R)^c} \psi_{iR}^0 s_2^+ \right]$	ψ_R^0
$\eta = (\eta^+ \quad \eta^0)^T$ Z₂ odd	<u>2</u>	1/2	$(Y_\psi^\eta)_{li} \left[\overline{L_\ell} \epsilon \eta^* \psi_{iR}^0 \right]$	ψ_R^0



Dirac

Lepton number
(conserved)forbid $y_\nu \bar{L} \in \Phi^* \nu_R$ at tree level
($\nu_R: -$)

	SU(2) _L	U(1) _Y	L#	Z' ₂	Yukawa int.	Note
s^0	<u>1</u>	0	-2	+	$(Y_S^0)_{ij} \left[\overline{(\nu_{iR})^c} \nu_{jR} s^0 \right]$	Sym. ν_R
s_L^+	<u>1</u>	1	-2	+	$(Y_A^s)_{\ell\ell'} \left[\bar{L}_\ell \in L_{\ell'}^c, s_L^- \right]$	Antisym.
s_R^+	<u>1</u>	1	-2	-	$(Y^s)_{li} \left[\overline{(\ell_R)^c} \nu_{iR} s_R^+ \right]$	ν_R
s^{++}	<u>1</u>	2	-2	+	$(Y_S^s)_{\ell\ell'} \left[\overline{(\ell_R)^c} \ell'_R s^{++} \right]$	Sym.
$\Phi_\nu = \begin{pmatrix} \phi_\nu^+ & \phi_\nu^0 \end{pmatrix} T$	<u>2</u>	1/2	0	-	$(Y_\nu)_{li} \left[\bar{L}_\ell \in \Phi_\nu^* \nu_{iR} \right]$	ν_R
$\Phi_2 = \begin{pmatrix} \phi_2^+ & \phi_2^0 \end{pmatrix} T$	<u>2</u>	1/2	0	+	$y_\ell \left[\bar{L}_\ell \Phi_2 \ell_R \right]$	Diag.
$\Delta = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}$	<u>3</u>	1	-2	+	$(Y_S^\Delta)_{\ell\ell'} \left[\bar{L}_\ell \Delta^\dagger \in L_{\ell'}^c \right]$	Sym.
s_2^0 Z₂ odd	<u>1</u>	0	-1	-	$(Y_\psi^0)_{ij} \left[\overline{(\nu_{iR})^c} \psi_{jR}^0 s_2^0 \right]$	ν_R ψ_R^0
s_2^+ Z₂ odd	<u>1</u>	1	-1	+	$(Y_\psi^+)_{li} \left[\overline{(\ell_R)^c} \psi_{iR}^0 s_2^+ \right]$	ψ_R^0
$\eta = (\eta^+ \ \eta^0)^T$ Z₂ odd	<u>2</u>	1/2	-1	+	$(Y_\psi^\eta)_{li} \left[\bar{L}_\ell \in \eta^* \psi_{iR}^0 \right]$	ψ_R^0

For Dirac ν masses (without DM)

Scalar with leptonic Yukawa int.

								Z_2 -odd		
	s^0	s_L^+	s_R^+	s^{++}	Φ_ν	Φ_2	Δ	s_2^0	s_2^+	η
$SU(2)_L$	<u>1</u>	<u>1</u>	<u>1</u>	<u>1</u>	<u>2</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>1</u>	<u>2</u>
$U(1)_Y$	0	1	1	2	1/2	1/2	1	0	1	1/2
L#	-2	-2	-2	-2	0	0	-2	-1	-1	-1
Z'_2	+	+	-	+	-	+	+	-	+	+
D1		✓	✓							
D2			✓				✓			
D3			✓	✓			✓			
D4			✓	✓						
D5	✓		✓				✓			
D6	✓		✓							
D7					✓					

Nasri-Moussa
model

New

ν -philic 2HDM

For Dirac ν masses (with DM)

Scalar with leptonic Yukawa int.

								Z_2 -odd		
	s^0	s_L^+	s_R^+	s^{++}	Φ_ν	Φ_2	Δ	s_2^0	s_2^+	η
$SU(2)_L$	<u>1</u>	<u>1</u>	<u>1</u>	<u>1</u>	<u>2</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>1</u>	<u>2</u>
$U(1)_Y$	0	1	1	2	1/2	1/2	1	0	1	1/2
L#	-2	-2	-2	-2	0	0	-2	-1	-1	-1
Z'_2	+	+	-	+	-	+	+	-	+	+
D8		✓						✓	✓	
D9							✓	✓	✓	
D10			✓							✓
D11			✓			✓			✓	
D12			✓						✓	
D13			✓			✓		✓		
D14			✓					✓		
D15						✓		✓	✓	
D16								✓	✓	
D17			✓						✓	✓
D18								✓		✓

New

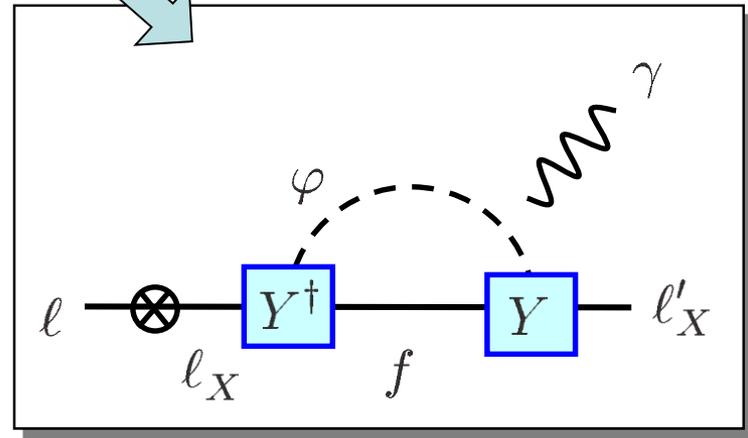
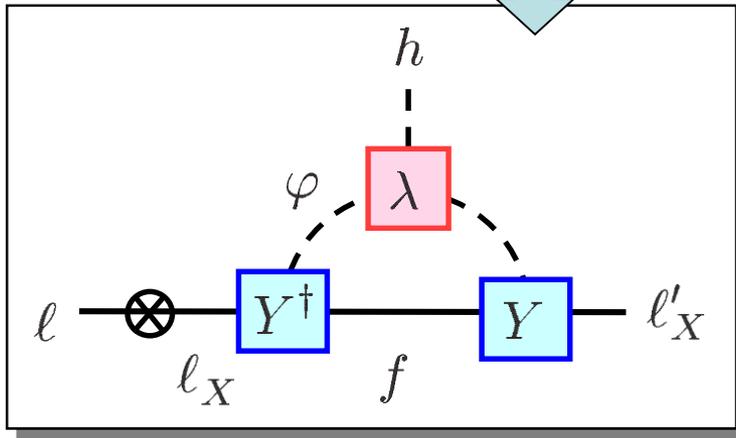
Gu-Sarkar
model

18 Mechanisms

- 
- Contents
- Introduction
 - Classification of Models
 - **Experimental Tests**
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$$h \rightarrow ll' \text{ and } l \rightarrow l'\gamma$$

Toy model : $\mathcal{L} = Y_{ae} \left[\overline{f_a} l_X \varphi \right] - \lambda |\Phi|^2 |\varphi|^2 + \dots$
 $X = L, R$



$$\text{BR}(h \rightarrow ll'_X) \sim 0.1 \frac{\lambda^2}{(2 - 3Q_\varphi)^2} \text{BR}(l \rightarrow l'_X \gamma)$$

Too small $\text{BR}(h \rightarrow ll')$
to be observed.
If observed,
the toy model is excluded.



$$\left\{ \begin{array}{l} \text{BR}(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13} \\ \text{MEG collab., EPJC76, no.8, 434} \\ \text{BR}(\tau \rightarrow e\gamma) < 3.3 \times 10^{-8} \\ \text{BR}(\tau \rightarrow \mu\gamma) < 4.4 \times 10^{-8} \\ \text{Babar collab., PRL104, 021802} \end{array} \right.$$

If $h \rightarrow \ell\ell'$ is Observed

M. Aoki, S. Kanemura, K. Sakurai, HS,
PLB763, 352 (2016)

"Mechanisms" for Majorana ν mass

	Scalar with leptonic Yukawa int.						$\ell \rightarrow \ell'\gamma$	
	s_L^+	s^{++}	Φ_2	Δ	s_2^+	η	ℓ'_L	ℓ'_R
SU(2) _L	<u>1</u>	<u>1</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>2</u>		
U(1) _Y	1	2	1/2	1	1	1/2		
Unbroken Z ₂	+	+	+	+	-	-		
M1	✓	✓					✓	✓
M2		✓	✓					✓
M3		✓						✓
M4				✓			✓	
M5	✓				✓		✓	✓
M6			✓		✓			✓
M7					✓			✓
M8						✓	✓	✓

$h \rightarrow \ell\ell'$ signal \Rightarrow BR($\ell \rightarrow \ell'\gamma$) is too large \Rightarrow **Excluded**

If $h \rightarrow \ell\ell'$ is Observed

M. Aoki, S. Kanemura, K. Sakurai, HS,
PLB763, 352 (2016)

"Mechanisms" for Majorana ν mass

	Scalar with leptonic Yukawa int.						$\ell \rightarrow \ell'\gamma$	
	s_L^+	s^{++}	Φ_2	Δ	Z_2 -odd		ℓ'_L	ℓ'_R
					s_2^+	η		
$SU(2)_L$	<u>1</u>	<u>1</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>2</u>		
$U(1)_Y$	1	2	1/2	1	1	1/2		
Unbroken Z_2	+	+	+	+	-	-		
M1	✓	✓					✓	✓
M2		✓	✓					✓
M3		✓						✓
M4				✓			✓	
M5	✓				✓		✓	✓
M6			✓		✓			✓
M7					✓			✓
M8						✓	✓	

All excluded for Majorana neutrinos ! (also type-I seesaw)

“Mechanisms” for Dirac ν mass (without DM)

	Scalar with leptonic Yukawa int.										$l \rightarrow l'\gamma$	
								Z ₂ -odd				
	s^0	s_L^+	s_R^+	s^{++}	Φ_ν	Φ_2	Δ	s_2^0	s_2^+	η		
SU(2) _L	<u>1</u>	<u>1</u>	<u>1</u>	<u>1</u>	<u>2</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>1</u>	<u>2</u>		
U(1) _Y	0	1	1	2	1/2	1/2	1	0	1	1/2		
L#	-2	-2	-2	-2	0	0	-2	-1	-1	-1		
Z' ₂	+	+	-	+	-	+	+	-	+	+		
D1		✓	✓								✓	✓
D2			✓			✓					✓	✓
D3			✓	✓		✓						✓✓
D4			✓	✓								✓✓
D5	✓		✓									✓
D6	✓		✓									✓
D7					✓						✓	

D3 and D4 survive

$$\text{BR}(l \rightarrow l'\gamma) \propto \left| \frac{(Y^{s\dagger}Y^s)_{\ell\ell'}}{m_{s_R^+}^2} + \frac{(Y_S^{s\dagger}Y_S^s)_{\ell\ell'}}{m_{s^{++}}^2} \right|^2 \ll \left| \frac{(Y^{s\dagger}Y^s)_{\ell\ell'}}{m_{s_R^+}^2} \right|^2$$

Cancellation is possible

“Mechanisms” for Dirac ν mass (with DM)

Scalar with leptonic Yukawa int.

	Scalar with leptonic Yukawa int.									$l \rightarrow l' \gamma$		
								Z_2 -odd			l'_L	l'_R
	s^0	s_L^+	s_R^+	s^{++}	Φ_ν	Φ_2	Δ	s_2^0	s_2^+	η		
$SU(2)_L$	<u>1</u>	<u>1</u>	<u>1</u>	<u>1</u>	<u>2</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>1</u>	<u>2</u>		
$U(1)_Y$	0	1	1	2	1/2	1/2	1	0	1	1/2		
LN	-2	-2	-2	-2	0	0	-2	-1	-1	-1		
Z'_2	+	+	-	+	-	+	+	-	+	+		
D8		✓						✓	✓		✓	✓
D9							✓	✓	✓		✓	✓
D10			✓							✓	✓	✓
D11			✓			✓			✓			✓✓
D12			✓						✓			✓✓
D13			✓			✓		✓				✓
D14			✓					✓				✓
D15										✓		✓
D16								✓	✓			✓
D17			✓						✓	✓	✓	✓✓
D18								✓		✓	✓	✓

D11, D12, and D17 survive

“Mechanisms” for Dirac ν mass

	Scalar with leptonic Yukawa int.										$l \rightarrow l' \gamma$	
								Z ₂ -odd				
	s^0	s_L^+	s_R^+	s^{++}	Φ_ν	Φ_2	Δ	s_2^0	s_2^+	η		
SU(2) _L	<u>1</u>	<u>1</u>	<u>1</u>	<u>1</u>	<u>2</u>	<u>2</u>	<u>3</u>	<u>1</u>	<u>1</u>	<u>2</u>		
U(1) _Y	0	1	1	2	1/2	1/2	1	0	1	1/2		
LN	-2	-2	-2	-2	0	0	-2	-1	-1	-1		
Z' ₂	+	+	-	+	-	+	+	-	+	+		
D3			✓	✓		✓						✓✓
D4			✓	✓								✓✓
D11			✓			✓			✓			✓✓
D12			✓						✓			✓✓
D17			✓						✓	✓	✓	✓✓

These mechanisms for generating masses of **Dirac neutrinos**

can **survive** after discovery of $h \rightarrow ll'$

Summary



Simple models to generate **Majorana** ν masses can be classified into **8 Mechanisms**

Simple models to generate **Dirac** ν masses can be classified into **18 Mechanisms**



$$h \rightarrow ll'$$


Simple models for **Majorana** ν masses are **excluded**

Some simple models for **Dirac** ν masses can **survive**

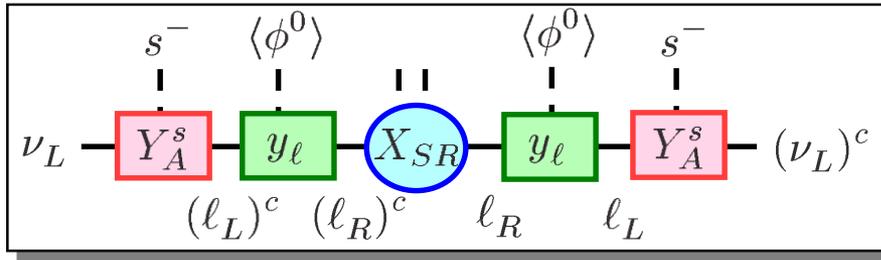
See you again in Osaka

Backup

3 Groups for Majorana ν mass

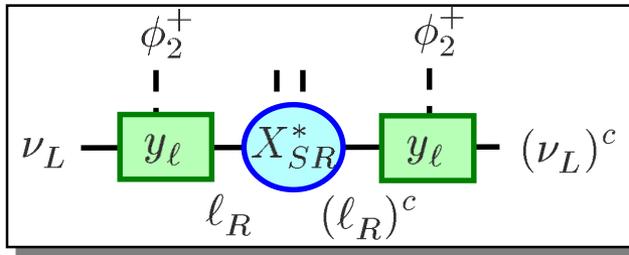
$$\text{I) } m_L \propto Y_A^s y_l X_{SR} y_l Y_A^{sT}$$

Y_A : **Antisym.** Yukawa for s^-



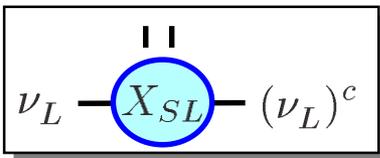
$$\text{II) } m_L \propto y_l X_{SR}^* y_l$$

y_l : **Diagonal** Yukawa ($\propto m_\ell$) for Φ



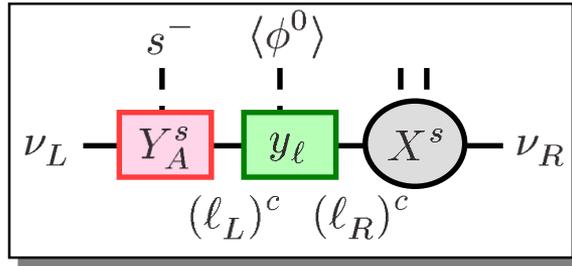
$$\text{III) } m_L \propto X_{SL}$$

$X_{SL}(X_{SR})$: **Sym.** matrix



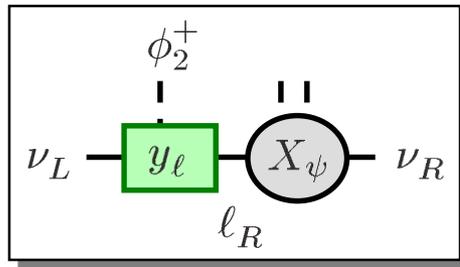
7 Groups for Dirac ν mass

$$\mathbf{I}') m_D \propto Y_A^s y_\ell X^s$$



Y_A : **Antisym.** Yukawa for s^-

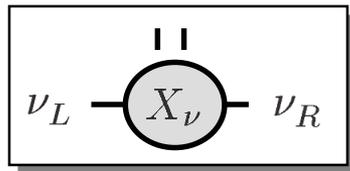
$$\mathbf{II}') m_D \propto y_\ell X_\psi$$



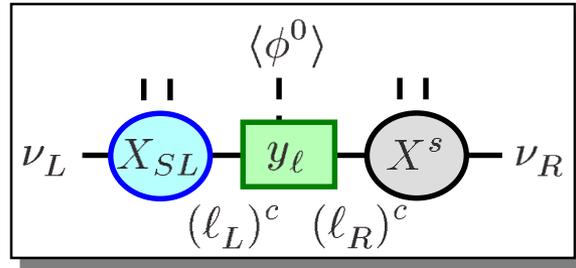
y_ℓ : **Diagonal** Yukawa ($\propto m_\ell$) for Φ

$X_\nu(X^s, X_\psi)$: Arbitrary

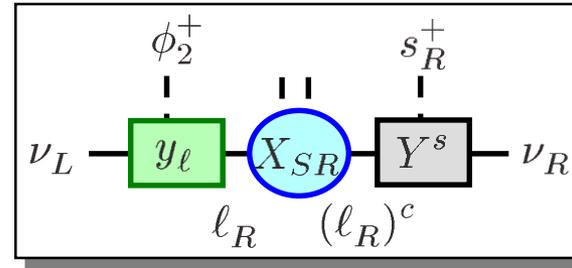
$$\mathbf{III}') m_D \propto X_\nu$$



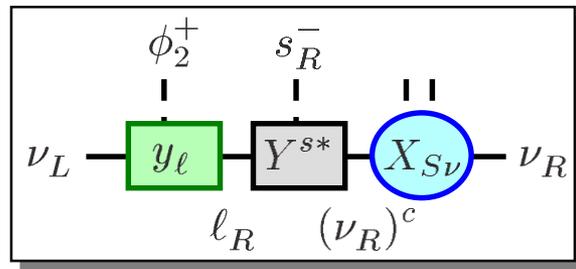
$$\text{IV) } m_D \propto X_{SL} y_l X^s$$



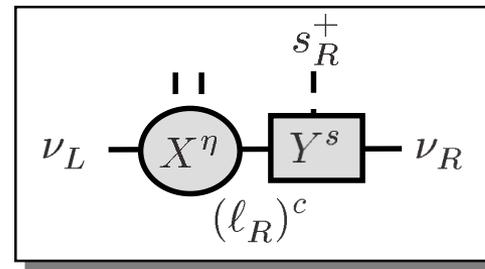
$$\text{V) } m_D \propto y_l X_{SR} Y^s$$



$$\text{VI) } m_D \propto y_l (Y^s)^* X_{S\nu}$$



$$\text{VII) } m_D \propto X^\eta Y^s$$



How can we test them ?



Concentrating on Yukawa

⇒ Flavor experiments

$$l \rightarrow \bar{l}_1 l_2 l_3$$

$$l \rightarrow l' \nu \bar{\nu}$$

$$l \rightarrow l' \gamma$$

$$h \rightarrow ll'$$

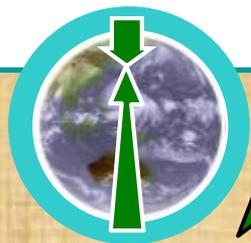
$$\theta_{23} \simeq 45^\circ$$

$$\theta_{13} \simeq 9^\circ$$

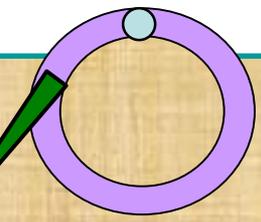
$$\delta ?$$

$$\theta_{12} \simeq 33^\circ$$

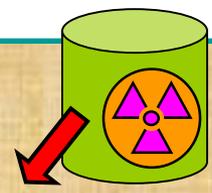
$$U_{MNS} \simeq \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0.71 & 0.71 \\ 0 & -0.71 & 0.71 \end{pmatrix} \begin{pmatrix} 0.99 & 0 & 0.15 \\ 0 & 1 & 0 \\ 0.15 & 0 & 0.99 \end{pmatrix} \begin{pmatrix} 0.83 & 0.55 & 0 \\ -0.55 & 0.83 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



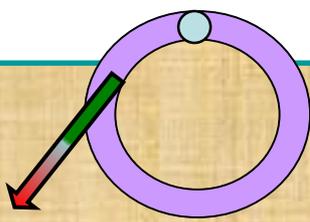
Atmospheric



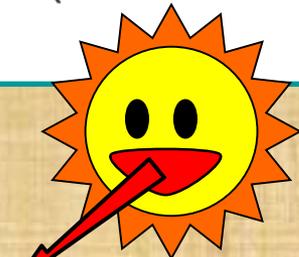
Accelerator



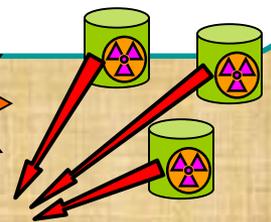
Reactor



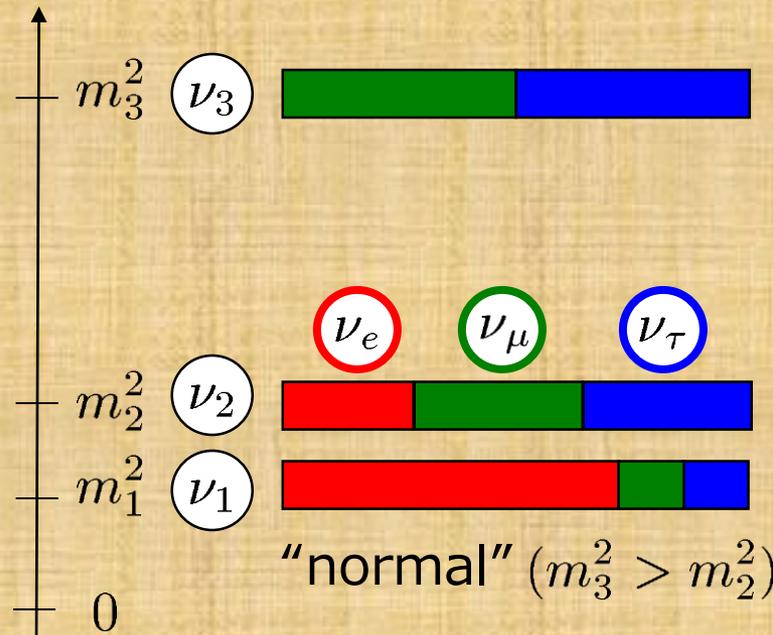
T2K



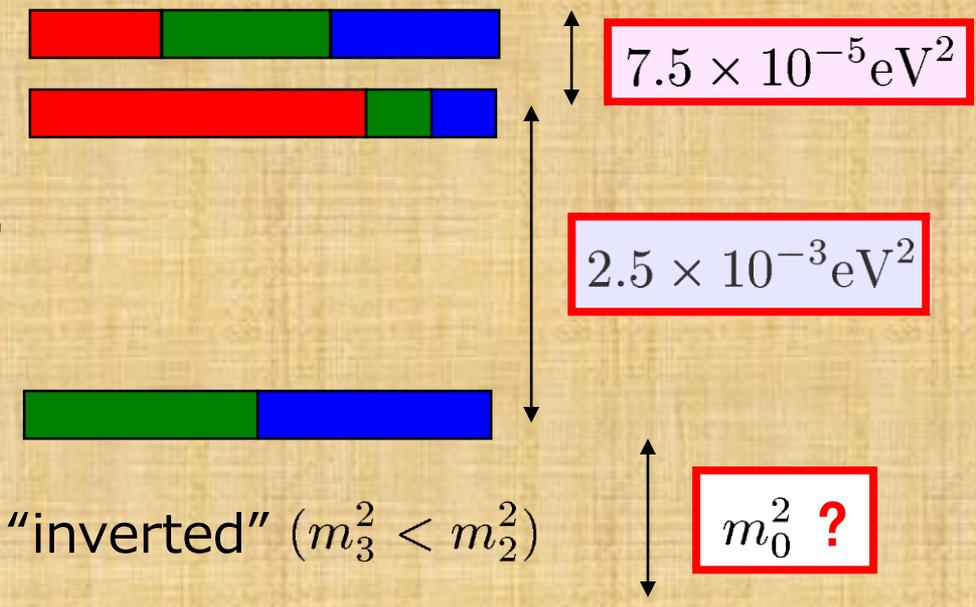
Solar



KamLAND



or



Oscillation Data

$$m_\nu = U_{\text{MNS}}^* \text{diag}(m_1 e^{i\alpha_{12}}, m_2, m_3 e^{i\alpha_{32}}) U_{\text{MNS}}^\dagger$$

$$U_{\text{MNS}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$c_{ij} \equiv \cos \theta_{ij} \quad s_{ij} \equiv \sin \theta_{ij}$$

$$\text{T2K : } \begin{cases} \sin^2 \theta_{23} = 0.514_{-0.056}^{+0.055}, & \Delta m_{32}^2 = (2.51 \pm 0.10) \times 10^{-3} \text{ eV}^2 \\ \sin^2 \theta_{23} = 0.511_{-0.055}^{+0.055}, & \Delta m_{23}^2 = (2.48 \pm 0.10) \times 10^{-3} \text{ eV}^2 \end{cases}$$

K. Abe *et al.*, PRL**112**, 181801 (2014)

$$\text{Daya Bay : } \sin^2 2\theta_{13} = 0.090_{-0.009}^{+0.008} \quad \text{R.P. An *et al.*, PRL**112**, 061801 (2014)}$$

$$\sin^2 2\theta_{13} = 0.084 \pm 0.005 \quad \text{R.P. An *et al.*, PRL**115**, 111802 (2015)}$$

$$\text{SNO : } \tan^2 \theta_{12} = 0.427_{-0.024}^{+0.027}, \quad \Delta m_{21}^2 = 7.46_{-0.19}^{+0.20} \times 10^{-5} \text{ eV}^2$$

B. Aharmim *et al.*, PRC**88**, 025501 (2013)

Absolute Mass Scale

$$\mathbf{I}) \quad m_\nu \propto Y_A y_\ell X_{SR} y_\ell Y_A^T$$

$$\mathbf{I}') \quad m_D \propto Y_A^s y_\ell X^s$$

$$(Y_A)_{\ell\ell'} \left[\overline{L}_\ell \in L_{\ell'}^c s^- \right]$$

$$\text{Det}(m_\nu) \propto \text{Det}(Y_A) = 0$$

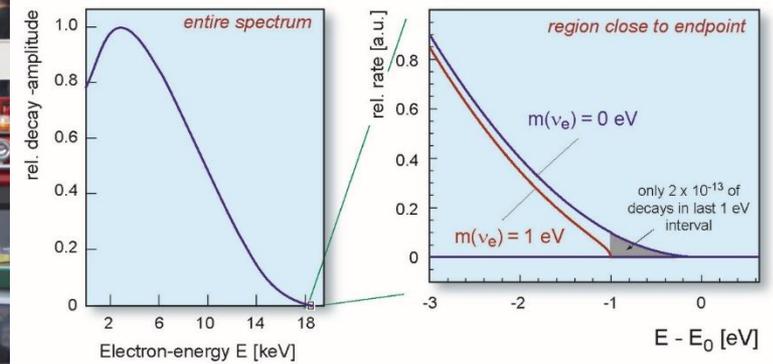
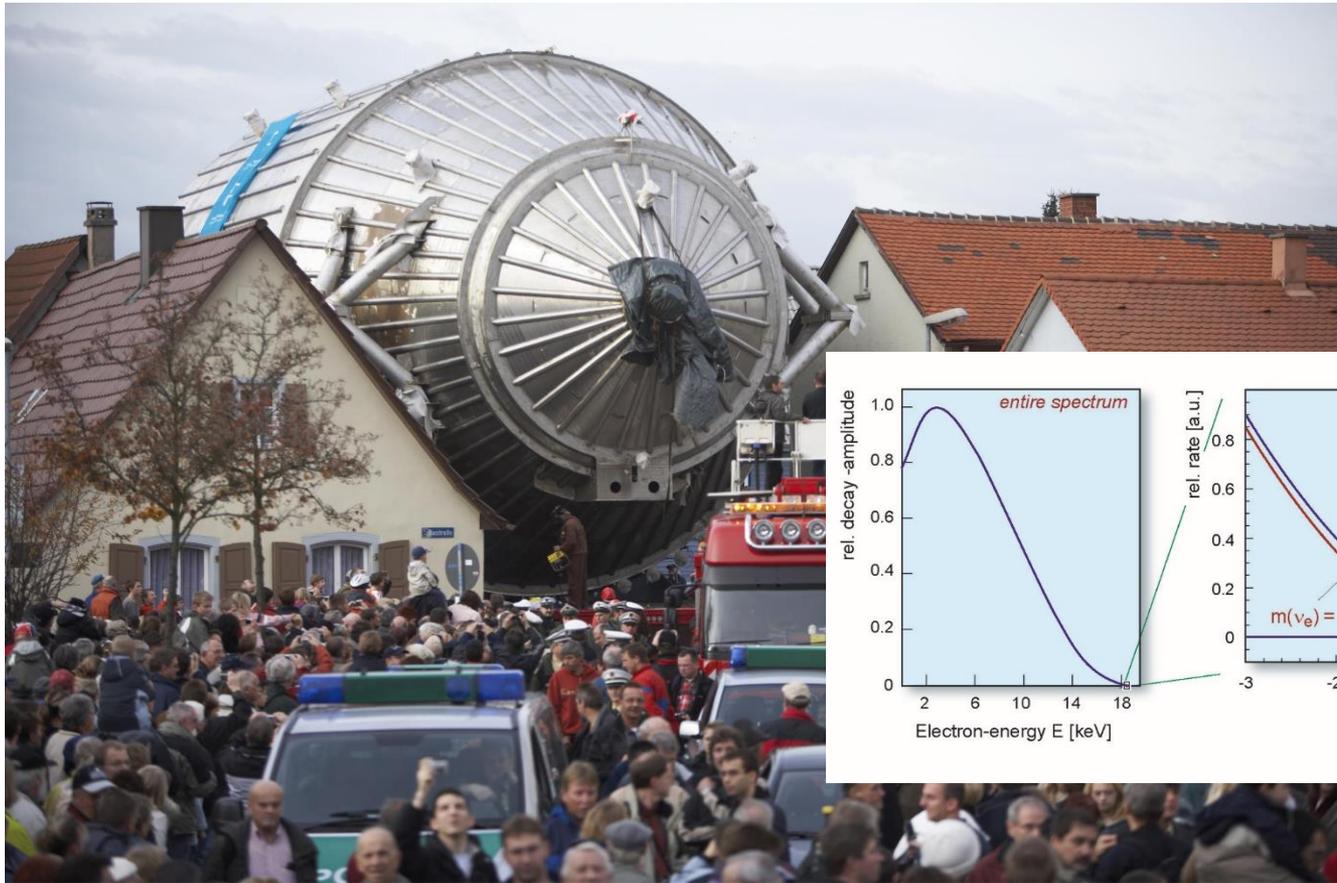
\longrightarrow $m_1 = 0$ or $m_3 = 0$ m_i : Neutrino mass eigenvalues
 ($m_2 \neq 0$ due to solar ν osc.)

\longrightarrow If $\min(m_i) \neq 0$ experimentally \Rightarrow **Excluded**

Direct (${}^3\text{H}$ β -decay): $m_\nu \simeq 0.35$ eV (5σ sensitivity)

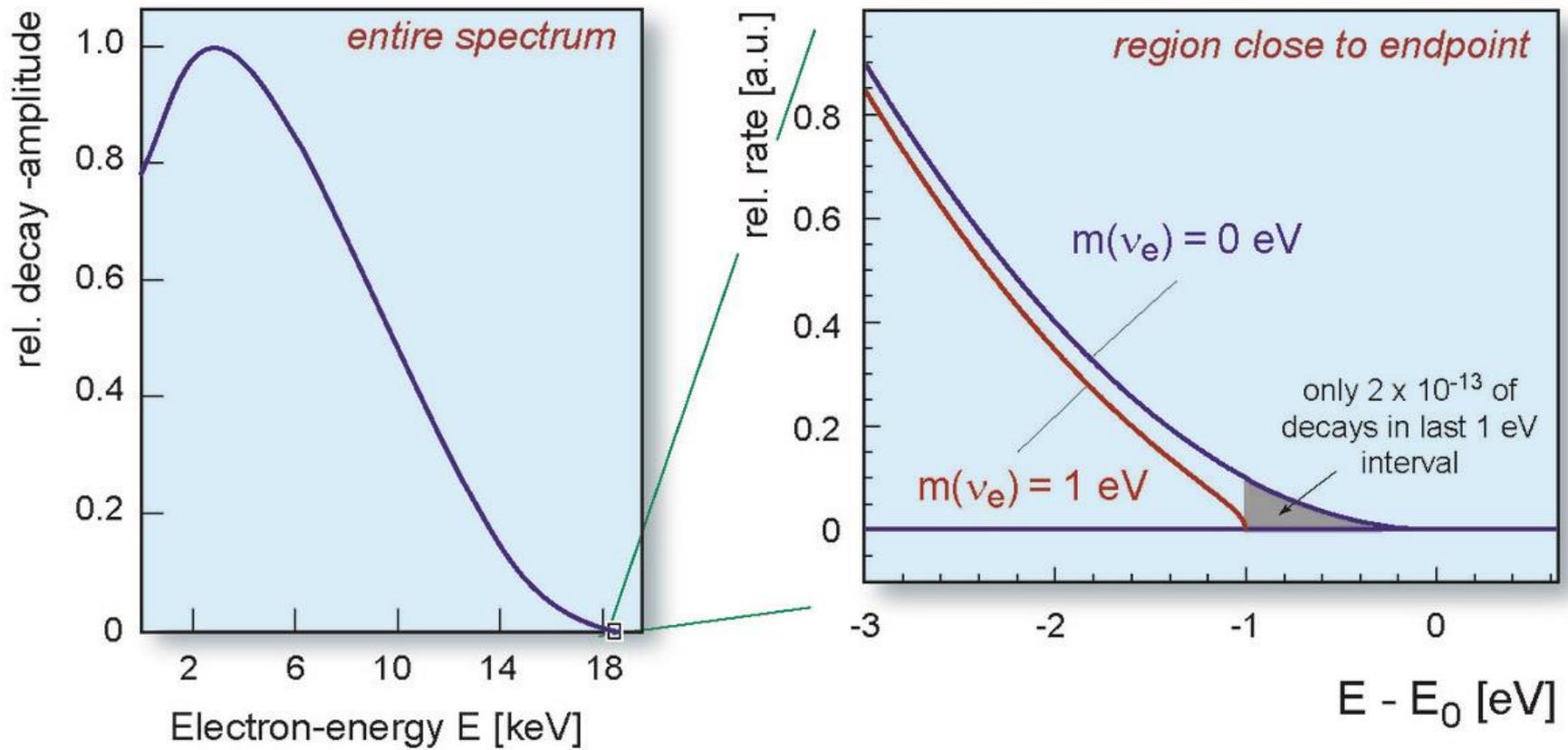
Indirect (Cosmology): $\Sigma m_i = \mathcal{O}(0.01)$ eV (95% CL sens.)

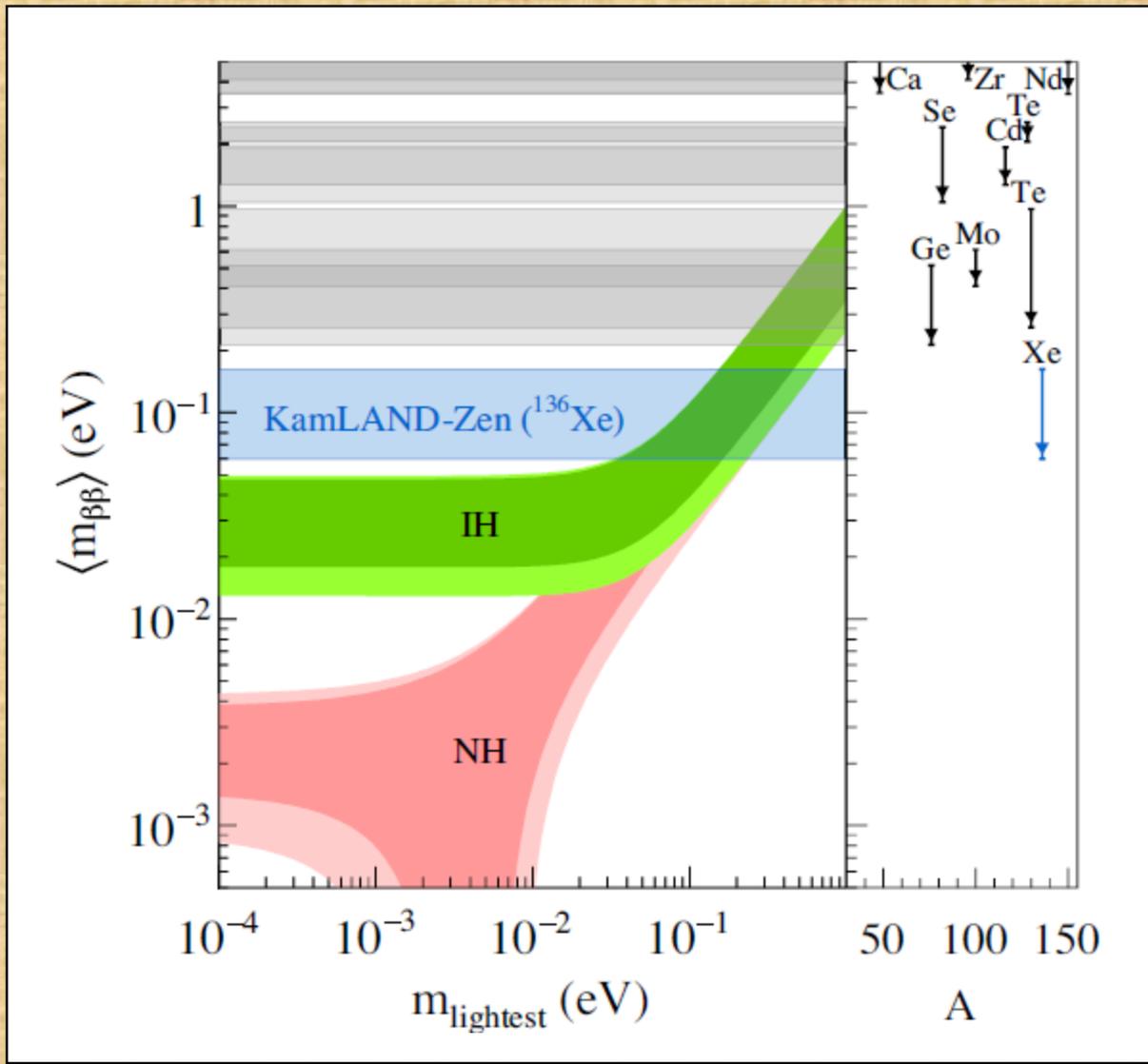
Not for a model but **for all models belong to Group-I and I'**
 (e.g. ZB model, KNT model, NM model)



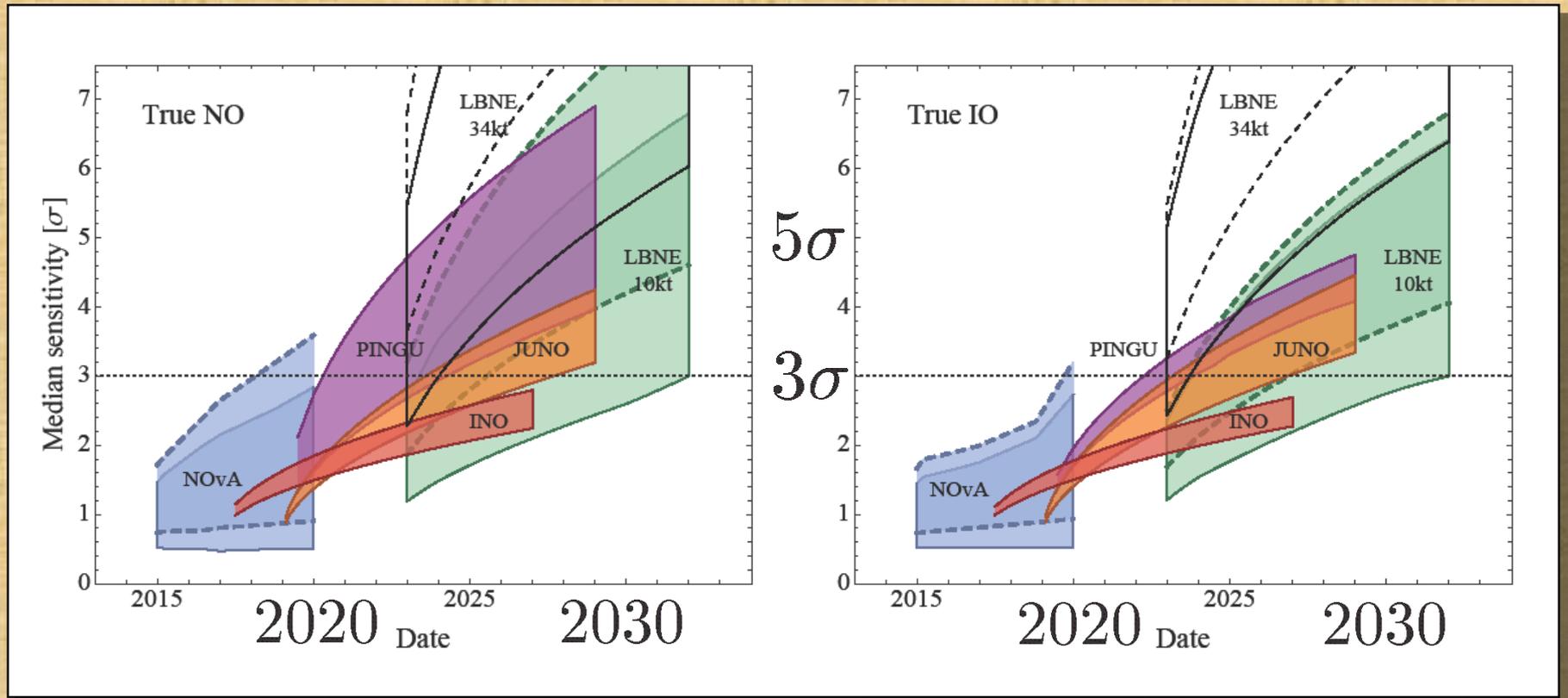
<https://www.katrin.kit.edu/213.php>

Main spectrometer of KATRIN experiment
 Transport through Leopoldshafen in Germany





KamLAND-ZEN collab., PRL**117**, 082503 (2016)



M. Blennow *et al.*, JHEP1403, 028 (2014)

LFV Decay of the Higgs Boson

$$\text{BR}(h \rightarrow \ell\ell') \equiv \text{BR}(h \rightarrow \bar{\ell}\ell') + \text{BR}(h \rightarrow \ell\bar{\ell}')$$

		ATLAS EPJC 77 , 70 (2017) (8 TeV, 20.3 fb ⁻¹)	CMS [1] PLB 763 , 472 (2016) [2] PLB 749 , 337 (2015) [3] JHEP 1806 , 001 (2018)	(8 TeV, 19.7 fb ⁻¹) (13 TeV, 35.9 fb ⁻¹)
$h \rightarrow \mu e$	Limit		$< 0.035 \times 10^{-2}$ [1]	
	Best fit			
$h \rightarrow \tau e$	Limit	$< 1.04 \times 10^{-2}$	$< 0.61 \times 10^{-2}$ [3]	
	Best fit	$-0.34^{+0.64}_{-0.66} \times 10^{-2}$	$0.30^{+0.18}_{-0.18} \times 10^{-2}$ [3]	
$h \rightarrow \tau\mu$	Limit	$< 1.43 \times 10^{-2}$	$< 1.51 \times 10^{-2}$ [2] $< 0.25 \times 10^{-2}$ [3]	
	Best fit	$0.53^{+0.51}_{-0.51} \times 10^{-2}$	$\left\{ \begin{array}{l} 0.84^{+0.39}_{-0.37} \times 10^{-2} (2.4\sigma) [2] \\ 0.00^{+0.12}_{-0.12} \times 10^{-2} [3] \end{array} \right.$	

My naïve calc.

$$(0.8 \pm 1.1) \times 10^{-3}$$

$h \rightarrow \ell\ell'$ at the loop level

Dim.-4 operator

$$\mathcal{L} = \mathbf{Y}_4 \left[\bar{L} \Phi \ell'_R \right]$$

$$\text{Mass : } \frac{v}{\sqrt{2}} \mathbf{Y}_4 \left[\bar{\ell}_L \ell'_R \right] \xrightarrow{\text{Diagonalize}} m_\ell \left[\bar{\ell}_L \ell_R \right]$$

$$\text{Int. : } \frac{1}{\sqrt{2}} \mathbf{Y}_4 \left[\bar{\ell}_L \ell'_R h \right] \dashrightarrow \frac{m_\ell}{v} \left[\bar{\ell}_L \ell_R h \right] \quad \text{no LFV}$$

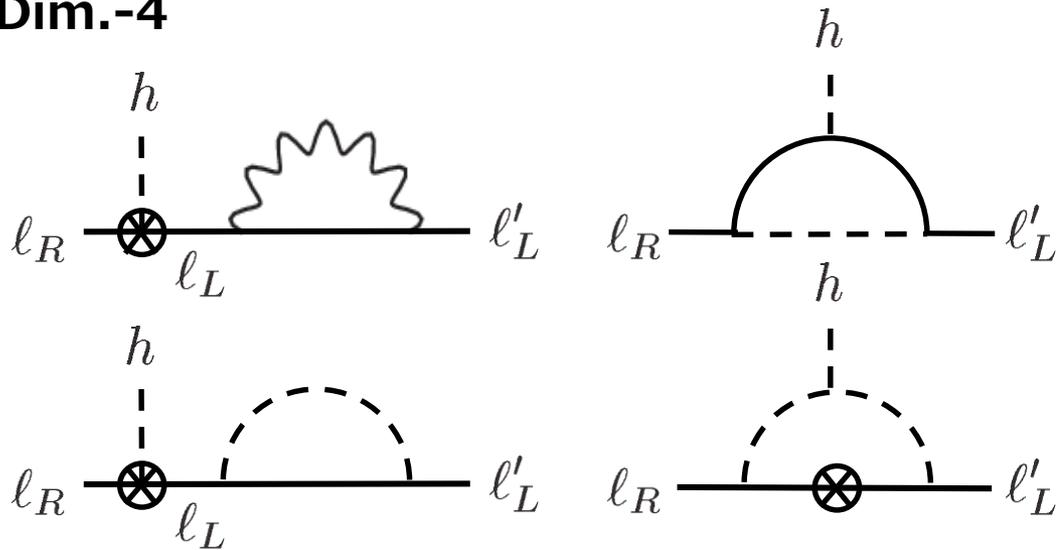
Dim.-6 operator

$$\mathcal{L} = \mathbf{Y}_4 \left[\bar{L} \Phi \ell'_R \right] + \frac{\mathbf{Y}_6}{\Lambda^2} \left[\bar{L} \Phi \ell'_R (\Phi^\dagger \Phi) \right]$$

$$\text{Mass : } v \left(\frac{1}{\sqrt{2}} \mathbf{Y}_4 + \frac{v^2}{2\Lambda^2} \mathbf{Y}_6 \right) \left[\bar{\ell}_L \ell'_R \right] \xrightarrow{\text{Diag.}} m_\ell \left[\bar{\ell}_L \ell_R \right] \quad \text{LFV}$$

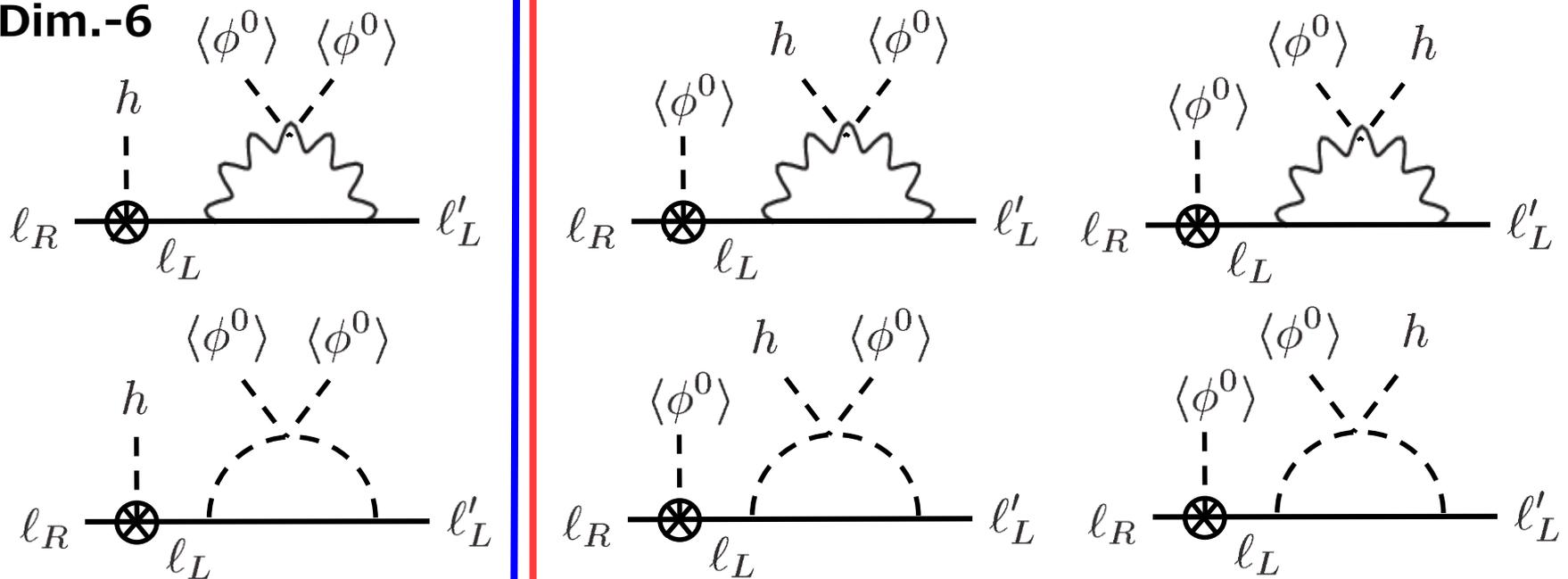
$$\text{Int. : } \left(\frac{1}{\sqrt{2}} \mathbf{Y}_4 + \frac{3v^2}{2\Lambda^2} \mathbf{Y}_6 \right) \left[\bar{\ell}_L \ell'_R h \right] \dashrightarrow \left(\frac{m_\ell}{v} \delta_{\ell\ell'} + \frac{v^2}{\Lambda^2} \mathbf{Y}_6 \right) \left[\bar{\ell}_L \ell'_R h \right]$$

Dim.-4

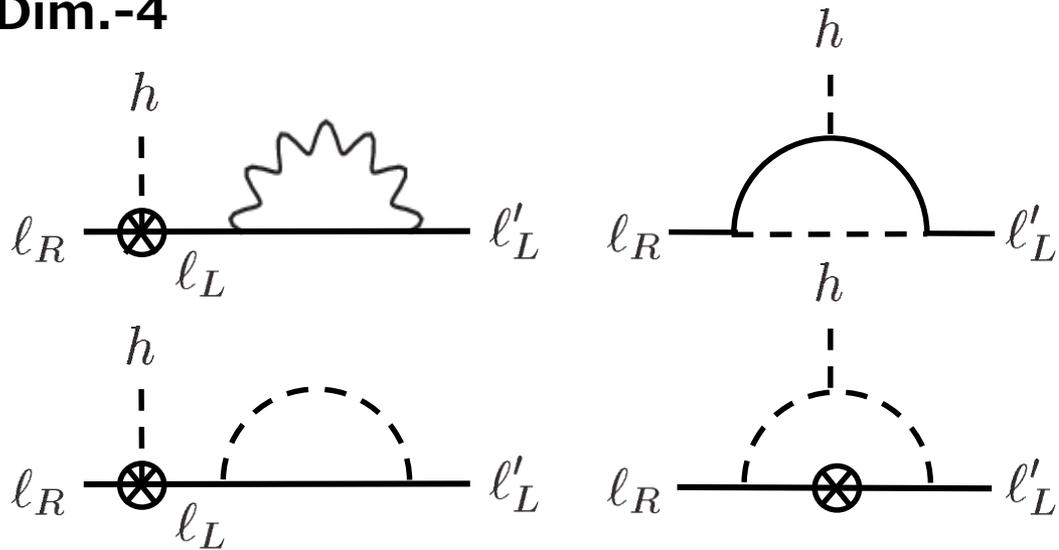


$$\left(\frac{m_l}{v} \delta_{ll'} + \frac{v^2}{\Lambda^2} \mathbf{Y}_6 \right) [\overline{l'_L} l_R h]$$

Dim.-6

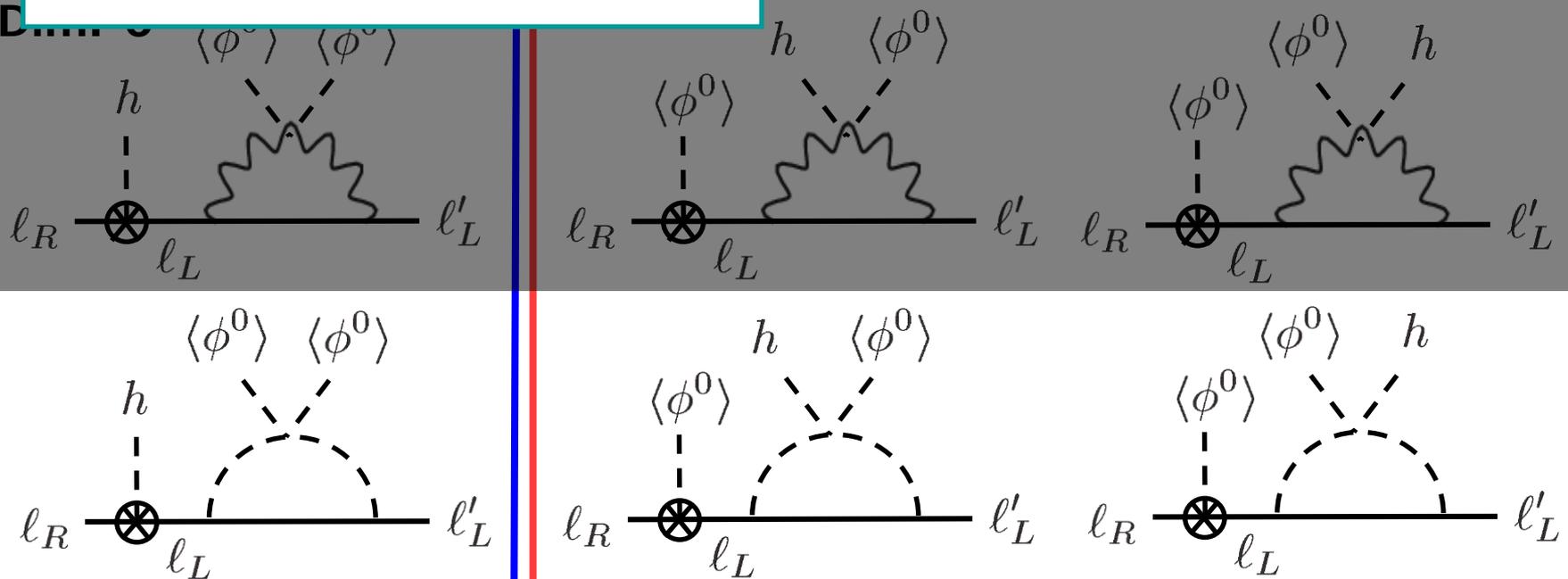


Dim.-4

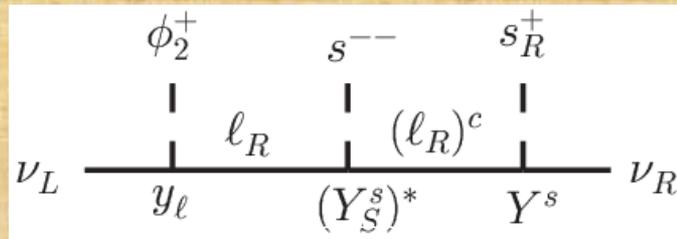


$$\left(\frac{m_l}{v} \delta_{ll'} + \frac{v^2}{\Lambda^2} \mathbf{Y}_6 \right) [\overline{l'_L} l_R h]$$

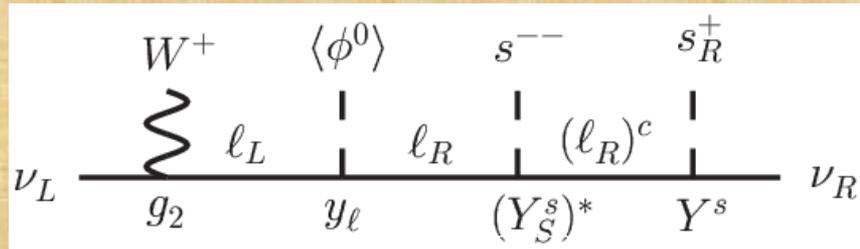
Suppressed by GIM mechanism



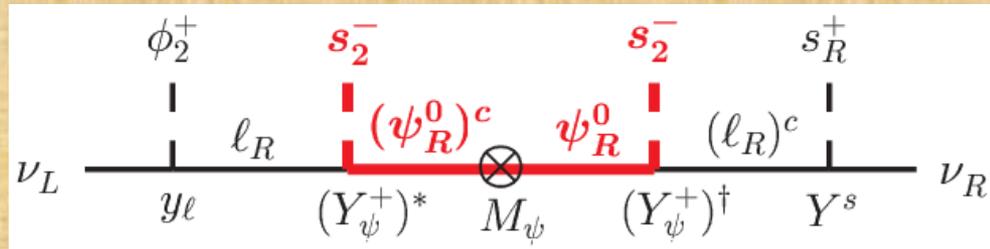
D3



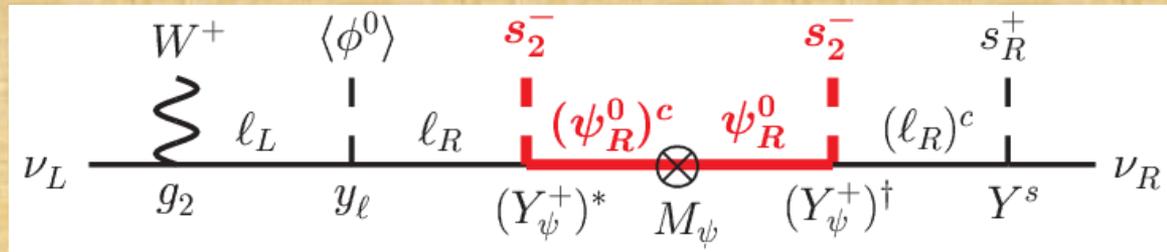
D4



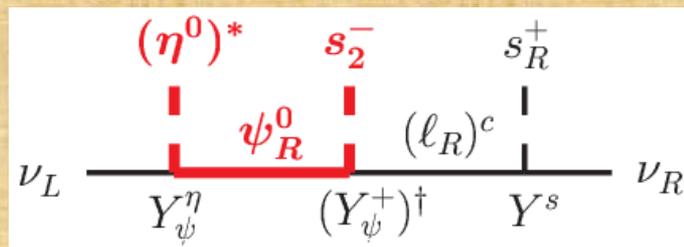
D11

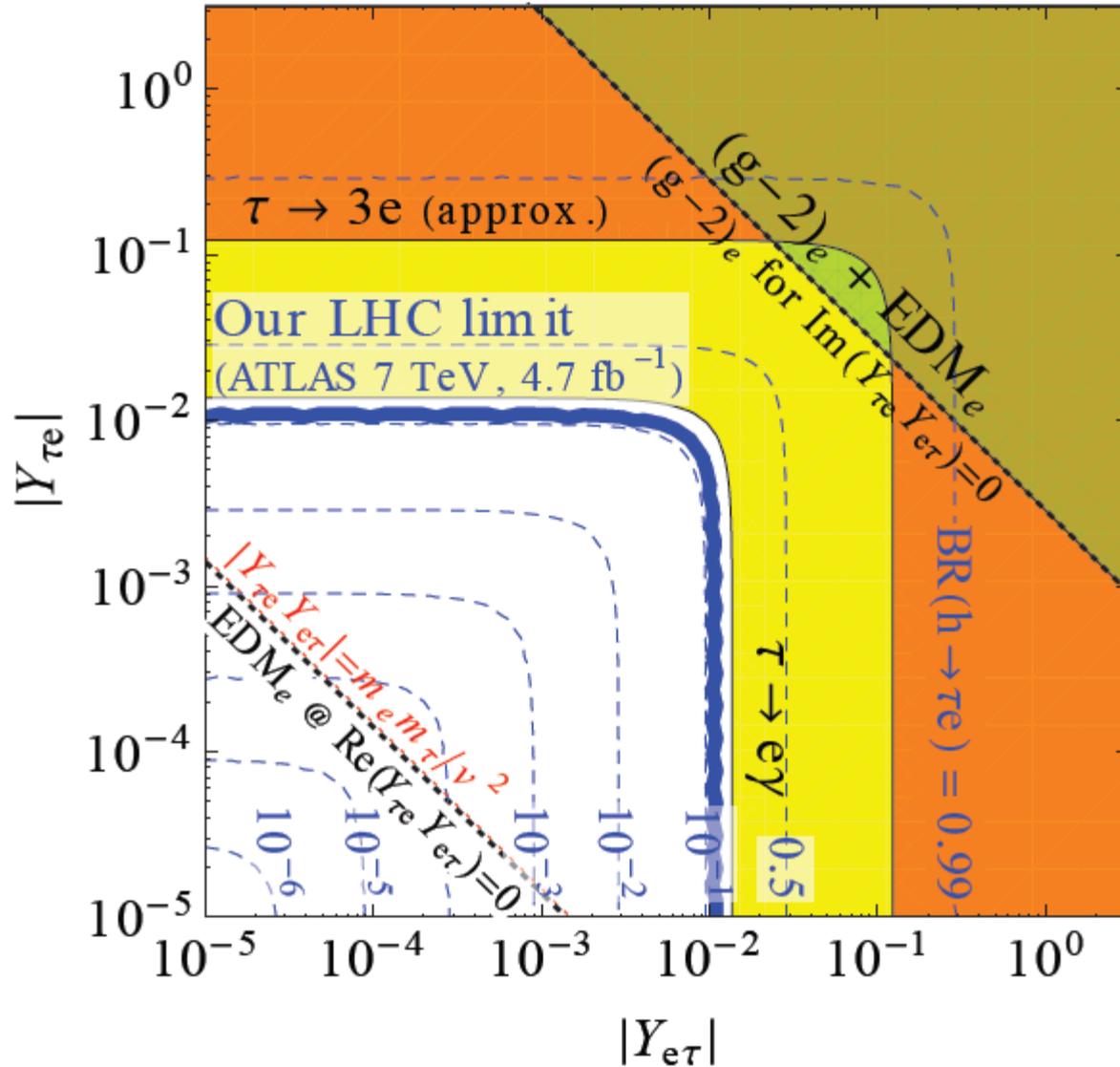


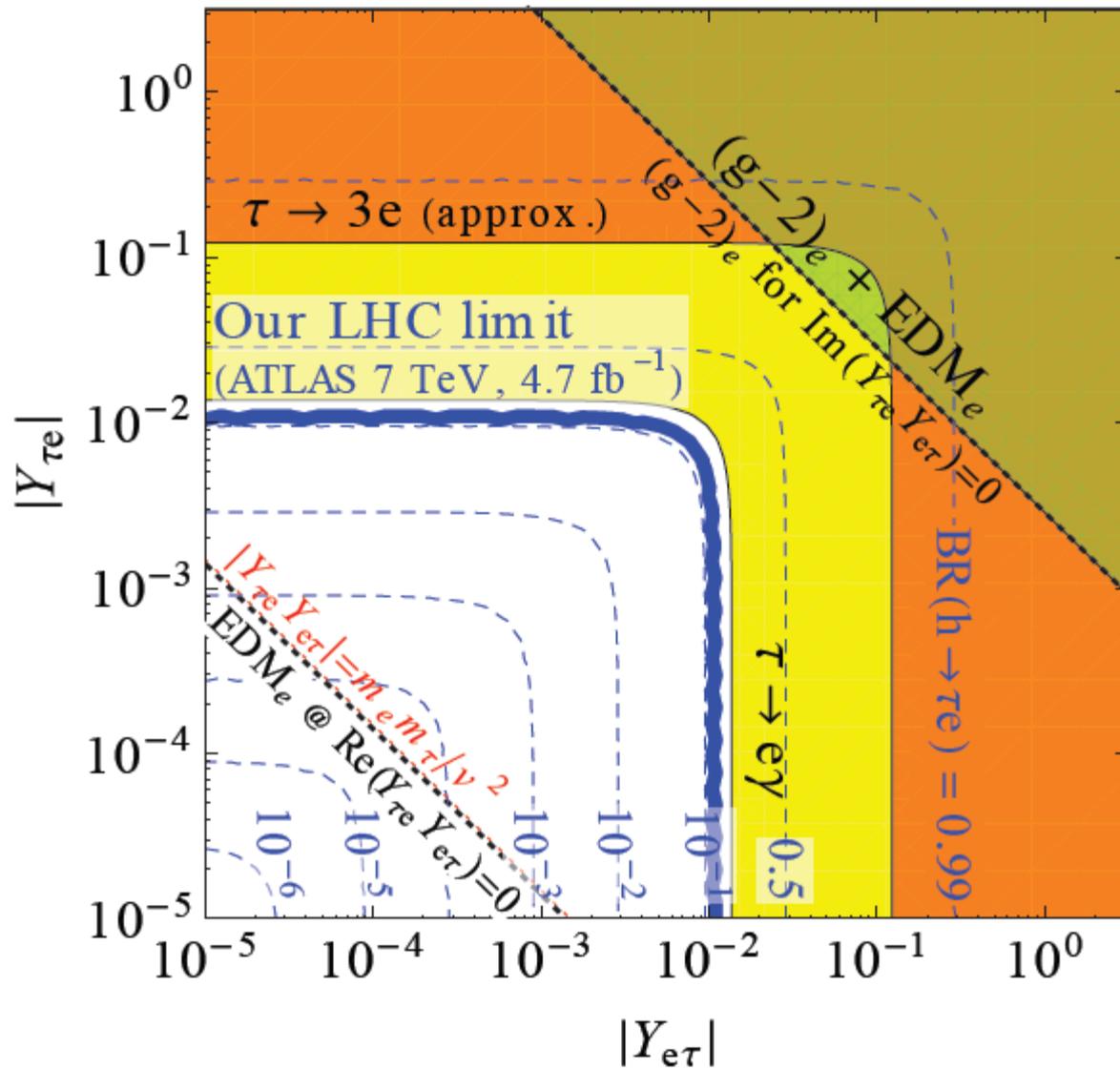
D12

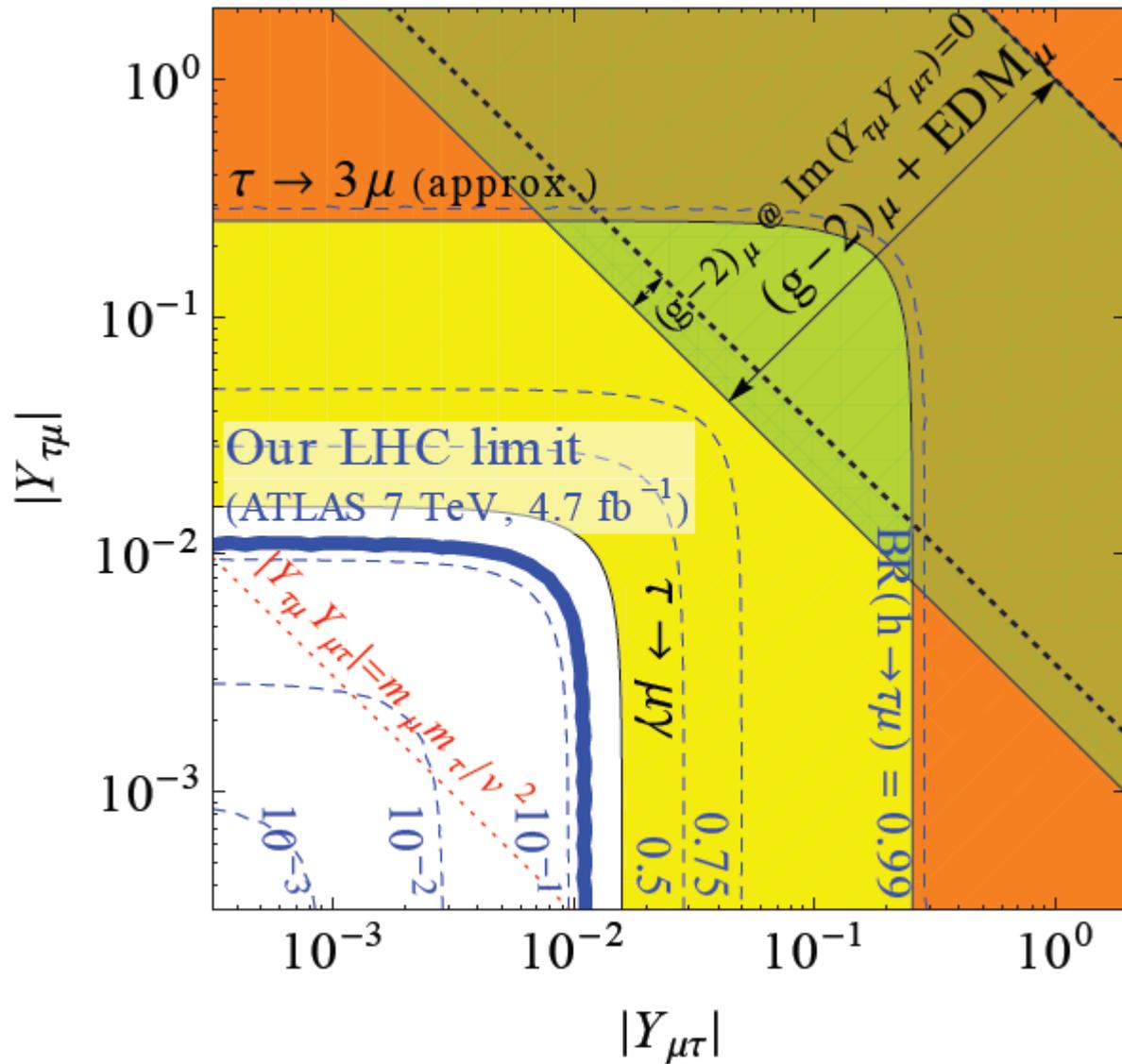


D17









$$\Phi_{LQ} = (\phi_{LQ}^{2/3}, \phi_{LQ}^{-1/3})^T \quad Y = 1/6$$

$$\begin{cases} \bar{Q} \Phi_{LQ} \nu_R \longrightarrow \bar{u}_L \nu_R \phi_{LQ}^{2/3} \\ \bar{d}_R \Phi_{LQ}^T \epsilon L \longrightarrow \bar{d}_R \ell_L \phi_{LQ}^{2/3} \end{cases}$$

$$\Phi'_{LQ} = (\phi'_{LQ}{}^{5/3}, \phi'_{LQ}{}^{2/3})^T \quad Y = 7/6$$

$$\begin{cases} \bar{Q} \Phi'_{LQ} \ell_R \longrightarrow \bar{d}_L \ell_R \phi'_{LQ}{}^{2/3} \\ \bar{u}_R \Phi'_{LQ}{}^T \epsilon L \longrightarrow \bar{u}_R \nu_L \phi'_{LQ}{}^{2/3} \end{cases}$$