

An $SU(2)_{\mu\tau}$ Model for Muon $g-2$ Anomaly

Koji Tsumura (Kyoto U.)

Workshop on Multi-Higgs Models,

Lisbon, Portugal, Sep 4-7, 2018

Based JHEP 05 (2018) 069 with C.-W. Chiang (NTU)

“Model with a gauged lepton flavor $SU(2)$ symmetry”

Contents

- Introduction
 - Muon $g-2$ anomaly
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 - Michel Decay
 - Collider Signature
 - Neutrino Mass
- Conclusion

Magnetic Moment of Muon

- g-factor : Int. btw Spin-B (Magnetic field)

- ✓ Classical (Dirac Eq.) $\rightarrow g=2$

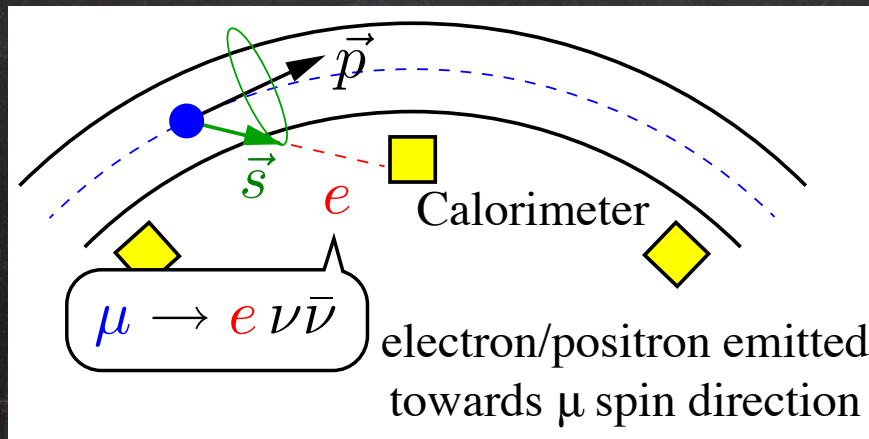
- ✓ Quantum $\rightarrow g=2(1+a_\mu)$ [$a_\mu=(g-2)/2$]

$$\mathcal{H} = -\vec{\mu} \cdot \vec{B} \quad \left(\vec{\mu} = g \frac{e}{2m} \vec{S} \right)$$

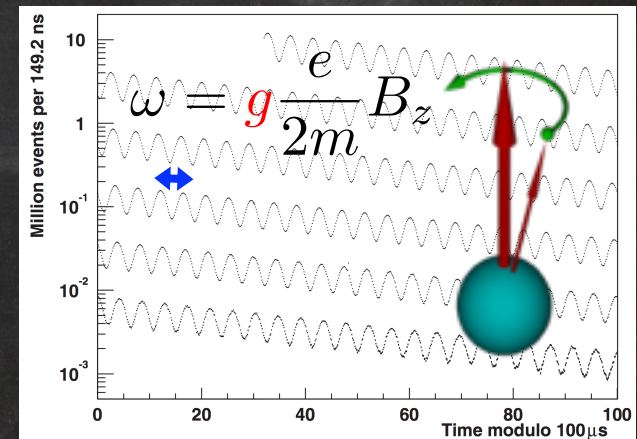
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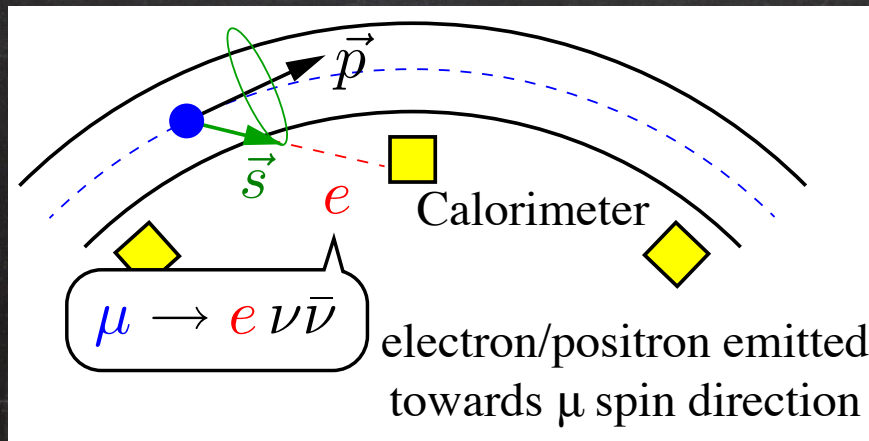
Larmor Precession



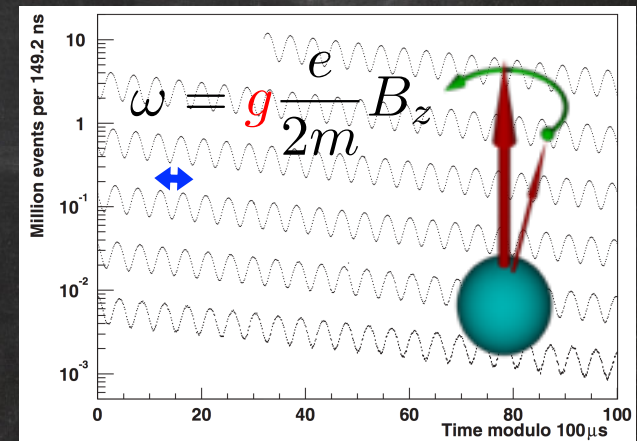
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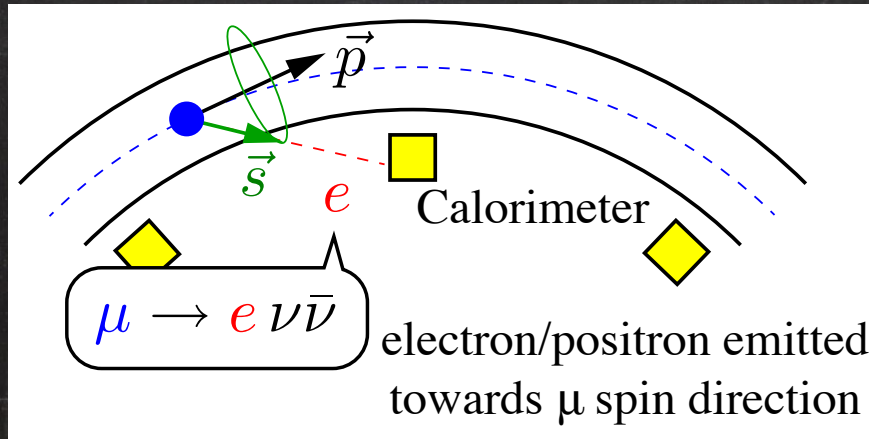
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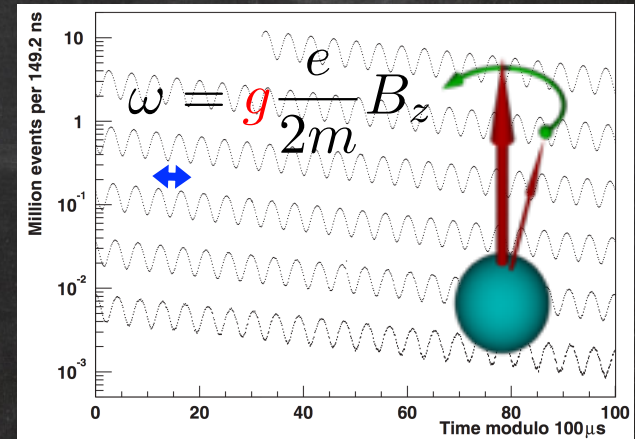
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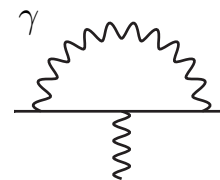
- **New Expt : FNAL E989 has been started**

KNT18 a_μ^{SM} update [KNT18: arXiv:1802.02995, PRD (in press)]

	<u>2011</u>	→	<u>2017</u>	
QED	11658471.81 (0.02)	→	11658471.90 (0.01)	[arXiv:1712.06060]
EW	15.40 (0.20)	→	15.36 (0.10)	[Phys. Rev. D 88 (2013) 053005]
LO HLbL	10.50 (2.60)	→	9.80 (2.60)	[EPJ Web Conf. 118 (2016) 01016]
NLO HLbL			0.30 (0.20)	[Phys. Lett. B 735 (2014) 90]
	<u>HLMNT11</u>		<u>KNT18</u>	
LO HVP	694.91 (4.27)	→	693.27 (2.46)	this work
NLO HVP	-9.84 (0.07)	→	-9.82 (0.04)	this work
NNLO HVP			1.24 (0.01)	[Phys. Lett. B 734 (2014) 144]
Theory total	11659182.80 (4.94)	→	11659182.05 (3.56)	this work
Experiment			11659209.10 (6.33)	world avg
Exp - Theory	26.1 (8.0)	→	27.1 (7.3)	this work
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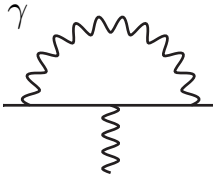
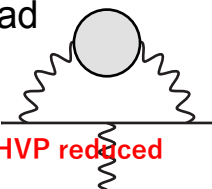
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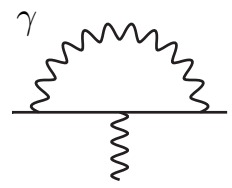
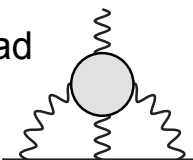
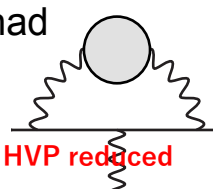
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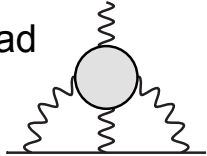
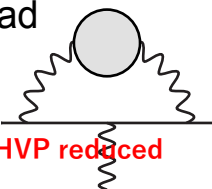
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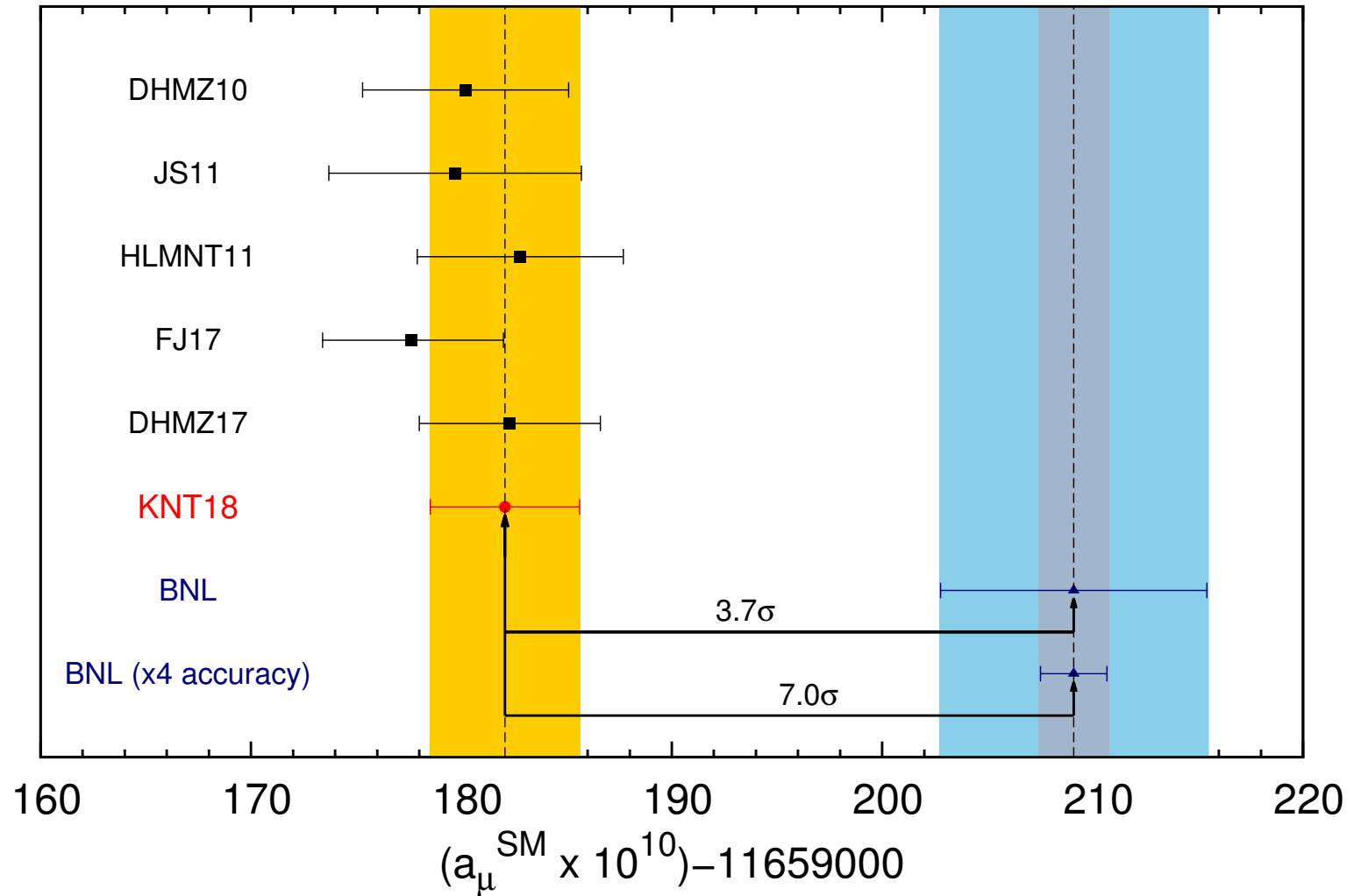
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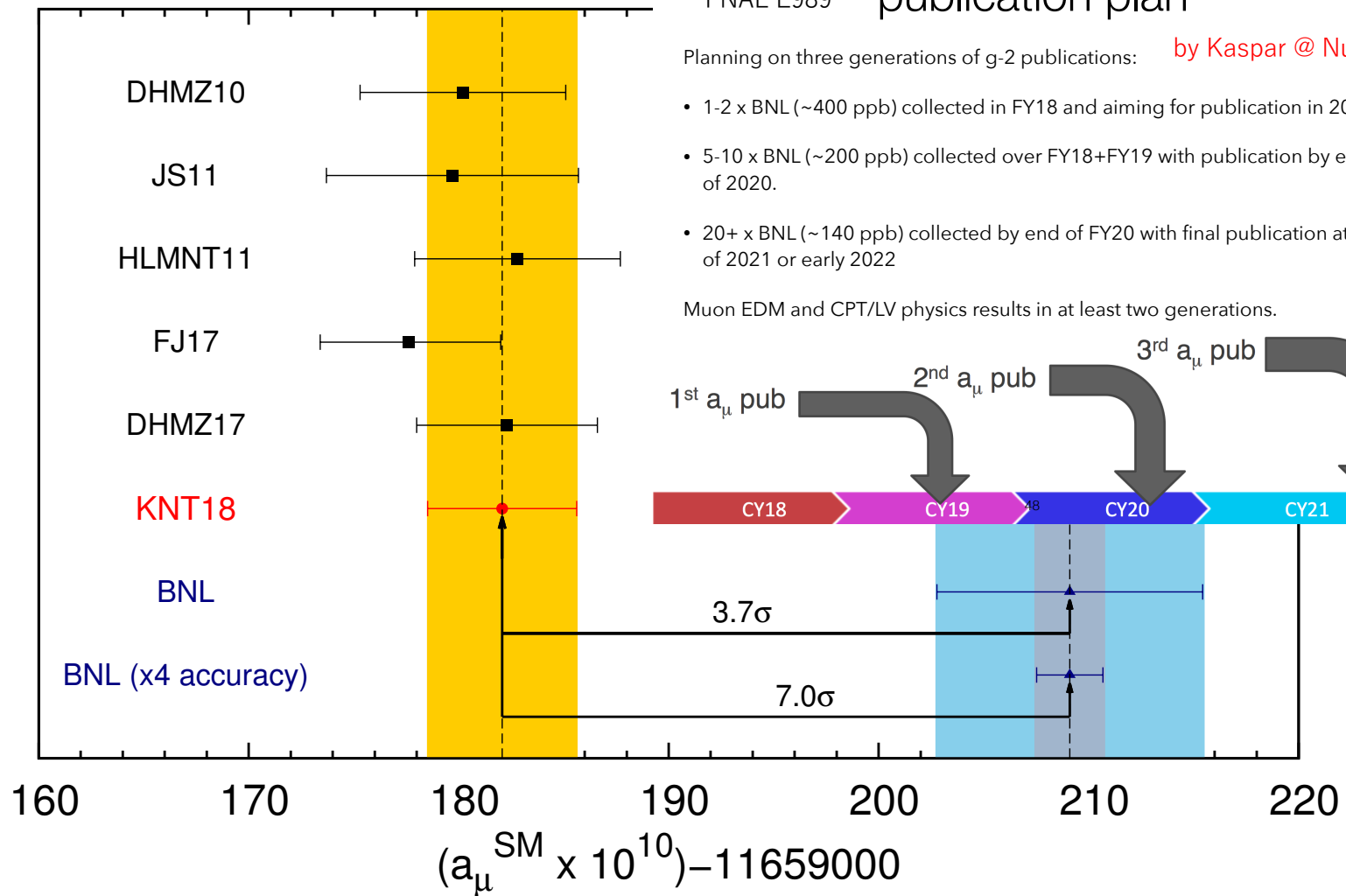
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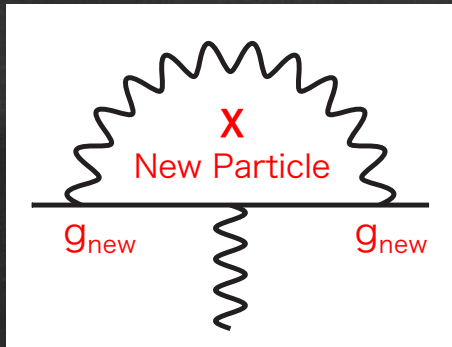
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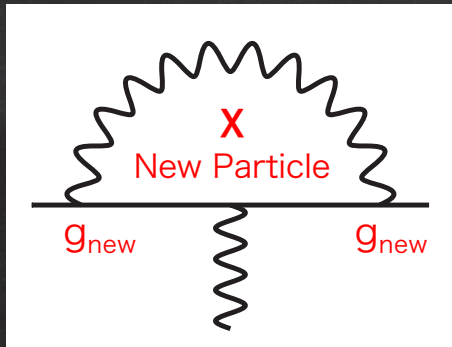


New Physics



$$\Delta a_{\mu}^{\text{new}} \sim \frac{g_{\text{new}}^2}{(4\pi)^2} \frac{M_{\mu}^2}{M_X^2}$$

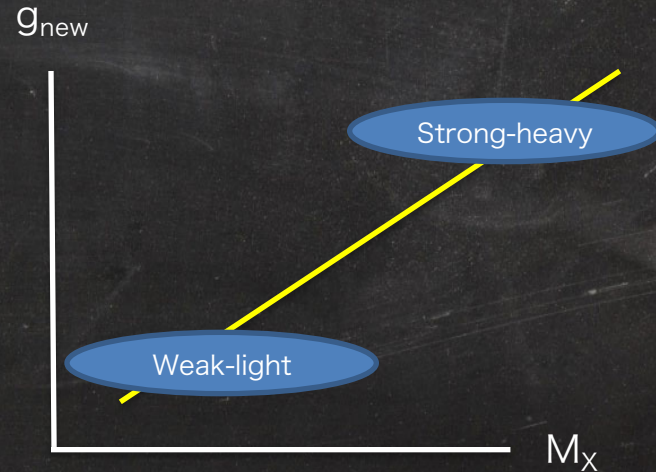
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To fit the observed discrepancy

$$g_{\text{new}} \approx g_{\text{weak}} \text{ and } M_X \approx M_W$$
$$(g_{\text{new}} \approx 10^{-3} \text{ and } M_X \approx M_{\mu})$$



U(1)_{μ-τ} Model

From Shimomura's slide @ NHWG20

Gauged U(1)_{L_μ-L_τ} model

He, Joshi, Lew, Volkas, PRD (1991) ,
R. Foot, Mod.Phys.Lett. (1991)

- A minimal extension of the SM
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- Large neutrino mixing (approx.)

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Asai, Hamaguchi, Nagata, 1705.00419

	l_e	e_R	l_μ	μ_R	l_τ	τ_R
L_μ	0	0	1	1	0	0
L_τ	0	0	0	0	1	1

The Lagrangian

Kinetic mixing with SM gauge boson

$$\mathcal{L}_{\text{new}} = -\frac{1}{4} Z'_{\mu\nu} Z'^{\mu\nu} + \frac{m_{Z'}^2}{2} Z'_\mu Z'^\mu + \frac{\epsilon}{2} B_{\mu\nu} Z'^{\mu\nu}$$

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new interactions for μ , τ and ν

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SU(2)_{μτ} Model

- Gauged Lepton Flavor SU(2)
 $e \sim$ singlet, $(\mu, \tau) \sim$ doublet
- Flavor Higgs is also needed
 (discuss later)
- Flavor Gauge bosons X

	SM leptons			
	$L = L_e$	$L^\alpha = \begin{bmatrix} L_\mu \\ L_\tau \end{bmatrix}$	e_R	$e_R^\alpha = \begin{bmatrix} \mu_R \\ \tau_R \end{bmatrix}$
$SU(2)_{\mu\tau}$	1	2	1	2
$SU(2)_L$	2	2	1	1
$U(1)_Y$	-1/2	-1/2	-1	-1

$$(\mathcal{D}^\lambda)^\alpha_\beta = \delta^\alpha_\beta D^\lambda + i g_X (X^\lambda)^\alpha_\beta$$

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$$(X^\lambda)^\alpha_\beta = \frac{(\sigma_a)^\alpha_\beta}{2} X_a^\lambda = \frac{1}{2} \begin{bmatrix} X_3^\lambda & \sqrt{2} X_+^\lambda \\ \sqrt{2} X_-^\lambda & -X_3^\lambda \end{bmatrix}, \quad X_\pm^\lambda = \frac{1}{\sqrt{2}} (X_1^\lambda \mp i X_2^\lambda)$$

X_3 is nothing but Z' in $U(1)_{L_\mu - L_\tau}$

The Lagrangian

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③

Flavor-charged X boson (New)

The Lagrangian

$$\mathcal{L}_{\text{new}} = -\frac{1}{4} Z'_{\mu\nu} Z'^{\mu\nu} + \frac{m_{Z'}^2}{2} Z'_\mu Z'^\mu + \frac{\epsilon}{2} B_{\mu\nu} Z'^{\mu\nu}$$

~~$\frac{\epsilon}{2} B_{\mu\nu} Z'^{\mu\nu}$~~

$$+ g' Z'_\mu (\bar{\mu} \gamma^\mu \mu + \bar{\nu}_\mu \gamma^\mu \nu_\mu - \bar{\tau} \gamma^\mu \tau - \bar{\nu}_\tau \gamma^\mu \nu_\tau)$$

Flavor Changing X

- Flavor Changing Interactions

$$-\frac{g_X}{\sqrt{2}} [X_+^\lambda (\bar{\mu} \gamma_\lambda \tau + \bar{\nu}_{\mu L} \gamma_\lambda \nu_{\tau L}) + X_-^\lambda (\bar{\tau} \gamma_\lambda \mu + \bar{\nu}_{\tau L} \gamma_\lambda \nu_{\mu L})]$$

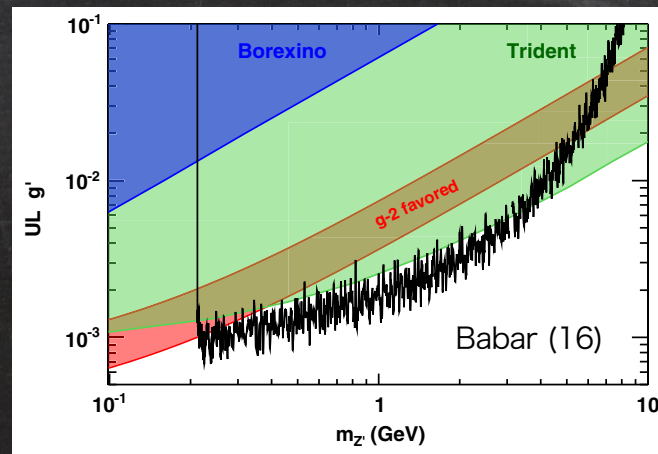
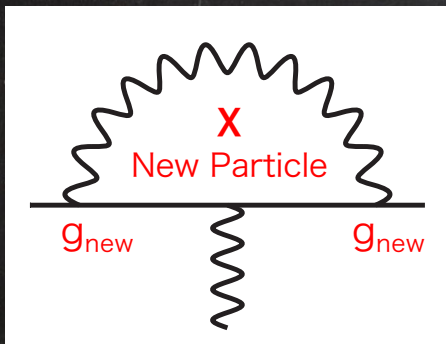
Flavor Changing X

- Flavor Changing Interactions

$$-\frac{g_X}{\sqrt{2}} [X_+^\lambda (\bar{\mu}\gamma_\lambda\tau + \overline{\nu_{\mu L}}\gamma_\lambda\nu_{\tau L}) + X_-^\lambda (\bar{\tau}\gamma_\lambda\mu + \overline{\nu_{\tau L}}\gamma_\lambda\nu_{\mu L})]$$

- Muon g-2

$$\Delta a_\mu(X_3) \approx \frac{g_X^2}{48\pi^2} \frac{M_\mu^2}{M_X^2}$$



Flavor Changing X

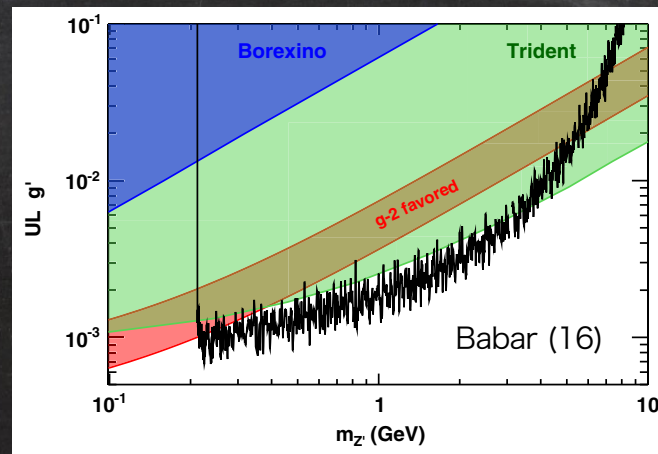
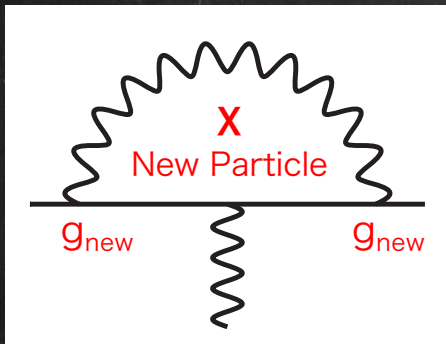
- Flavor Changing Interactions

$$-\frac{g_X}{\sqrt{2}} [X_+^\lambda (\bar{\mu}\gamma_\lambda\tau + \bar{\nu}_{\mu L}\gamma_\lambda\nu_{\tau L}) + X_-^\lambda (\bar{\tau}\gamma_\lambda\mu + \bar{\nu}_{\tau L}\gamma_\lambda\nu_{\mu L})]$$

- Muon g-2

Chirality Enhancement!!

$$\Delta a_\mu(X_3) \approx \frac{g_X^2}{48\pi^2} \frac{M_\mu^2}{M_X^2} \quad \Delta a_\mu(X_\pm) \approx \frac{g_X^2}{48\pi^2} \frac{M_\mu^2}{M_X^2} \left(6 \frac{M_\tau}{M_\mu} - 4\right) \sim 96.9 \times \Delta a_\mu(X_3)$$



Flavor Changing X

- Flavor Changing Interactions

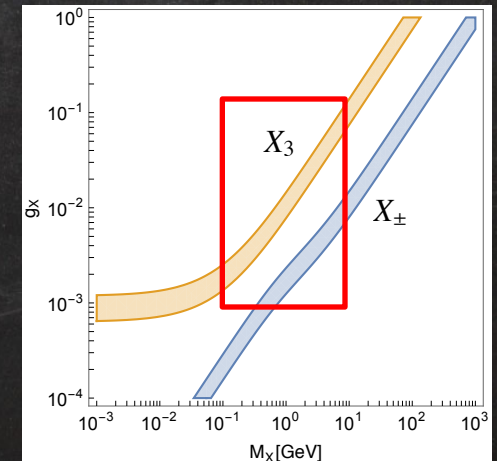
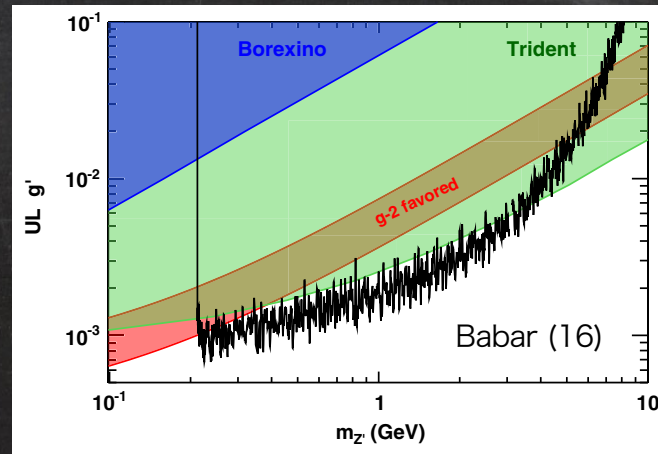
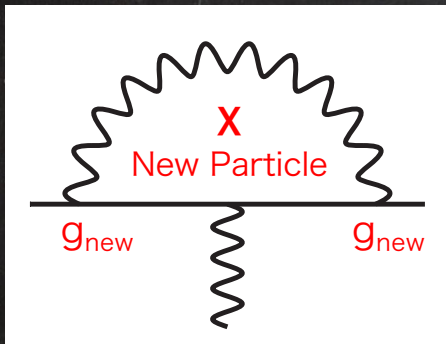
$$-\frac{g_X}{\sqrt{2}} [X_+^\lambda (\bar{\mu}\gamma_\lambda\tau + \bar{\nu}_{\mu L}\gamma_\lambda\nu_{\tau L}) + X_-^\lambda (\bar{\tau}\gamma_\lambda\mu + \bar{\nu}_{\tau L}\gamma_\lambda\nu_{\mu L})]$$

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Chirality Enhancement!!

$$\Delta a_\mu(X_3) \approx \frac{g_X^2}{48\pi^2} \frac{M_\mu^2}{M_X^2}$$

$$\Delta a_\mu(X_\pm) \approx \frac{g_X^2}{48\pi^2} \frac{M_\mu^2}{M_X^2} \left(6 \frac{M_\tau}{M_\mu} - 4\right) \sim 96.9 \times \Delta a_\mu(X_3)$$



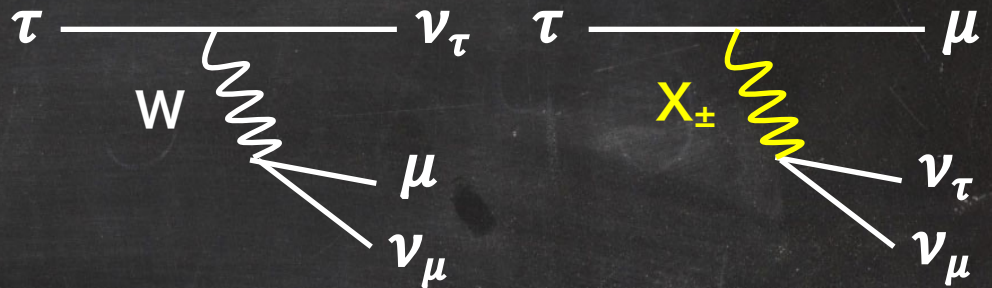
τ Michel Decay

- Flavor Changing Interactions

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- Lepton Universality

$$R_{\tau \rightarrow \mu / \tau \rightarrow e} = \frac{\Gamma(\tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau)}{\Gamma(\tau \rightarrow e \bar{\nu}_e \nu_\tau)}$$



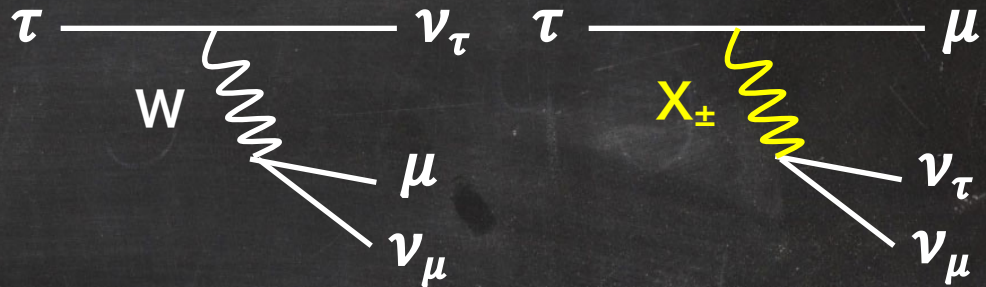
τ Michel Decay

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- Lepton Universality

$$R_{\tau \rightarrow \mu / \tau \rightarrow e} = \frac{\Gamma(\tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau)}{\Gamma(\tau \rightarrow e \bar{\nu}_e \nu_\tau)}$$



$$\eta^X \equiv \frac{R_{\tau \rightarrow \mu / \tau \rightarrow e}^X}{R_{\tau \rightarrow \mu / \tau \rightarrow e}^{\text{SM}}} = \left(1 + \alpha_X + \frac{1}{2} \alpha_X^2 \right) - 2 \alpha_X \frac{M_\mu}{M_\tau} \frac{g(\rho^{-2})}{f(\rho^{-2})}$$

$$\frac{\alpha_X}{2} \equiv \frac{g_X^2}{4\sqrt{2}G_F M_X^2}$$

Interference

$$f(z) \equiv 1 - 8z + 8z^3 - z^4 - 12z^2 \log z \approx 1$$

$$g(z) \equiv 1 + 9z - 9z^2 - z^3 + 6z \log z + 6z^2 \log z \approx 1$$

τ Michel Decay

- Flavor Changing Interactions

$$-\frac{g_X}{\sqrt{2}} [X_+^\lambda (\bar{\mu} \gamma_\lambda \tau + \bar{\nu}_{\mu L} \gamma_\lambda \nu_{\tau L}) + X_-^\lambda (\bar{\tau} \gamma_\lambda \mu + \bar{\nu}_{\tau L} \gamma_\lambda \nu_{\mu L})]$$

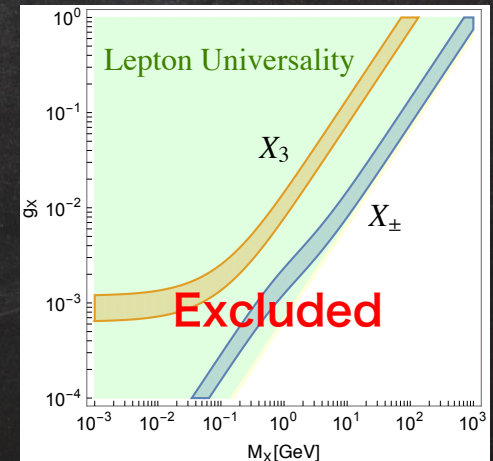
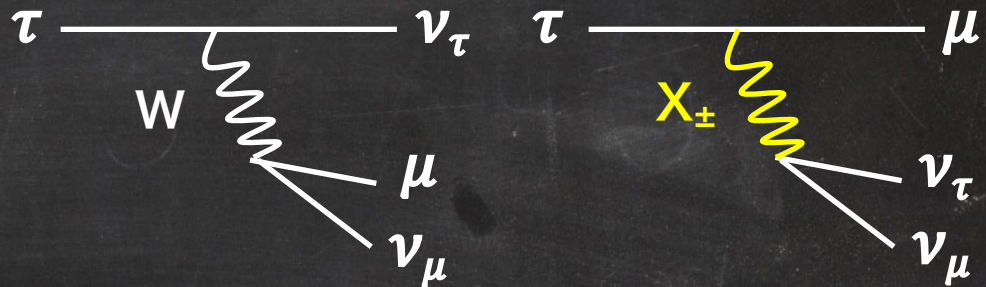
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So far, No Higgs ...

μ - τ Mass Degeneracy

- Yukawa sector

$$y_0 \bar{L}_\alpha \Phi_0 e_R^\alpha \rightarrow \frac{y_0 v_0}{\sqrt{2}} (\bar{\mu}_L \mu_R + \bar{\tau}_L \tau_R)$$

	SM leptons			
	$L = L_e$	$L^\alpha = \begin{bmatrix} L_\mu \\ L_\tau \end{bmatrix}$	e_R	$e_R^\alpha = \begin{bmatrix} \mu_R \\ \tau_R \end{bmatrix}$
$SU(2)_{\mu\tau}$	1	2	1	2
$SU(2)_L$	2	2	1	1
$U(1)_Y$	-1/2	-1/2	-1	-1

μ - τ Mass Degeneracy

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- Flavor Adjoint Higgs

$$\bar{L}_\alpha (y_0 \Phi_0 \delta_\beta^\alpha + 2y \Phi_\beta^\alpha) e_R^\beta$$

$$\Phi_\beta^\alpha = \frac{(\sigma_a)^\alpha_\beta}{2} \Phi_a = \frac{1}{2} \begin{bmatrix} \Phi_3 & \sqrt{2}\Phi_+ \\ \sqrt{2}\Phi_- & -\Phi_3 \end{bmatrix}, \quad \Phi_\pm = (\Phi_1 \mp i\Phi_2)/\sqrt{2}$$

Scalars		
Φ_0	Φ_β^α	S^α
1	3	2
2	2	1
+1/2	+1/2	0

μ - τ Mass Degeneracy

- Yukawa sector

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Scalars		
Φ_0	Φ^α_β	S^α
1	3	2
2	2	1
+1/2	+1/2	0

$$y_0 = \frac{\sqrt{2}}{v_0} \frac{M_\mu + M_\tau}{2}$$

$$y = \frac{\sqrt{2}}{v_3} \frac{M_\mu - M_\tau}{2}$$

μ - τ Mass Degeneracy

- Yukawa sector

$$y_0 \bar{L}_\alpha \Phi_0 e_R^\alpha \rightarrow \frac{y_0 v_0}{\sqrt{2}} (\bar{\mu}_L \mu_R + \bar{\tau}_L \tau_R)$$

	SM leptons			
	$L = L_e$	$L^\alpha = \begin{bmatrix} L_\mu \\ L_\tau \end{bmatrix}$	e_R	$e_R^\alpha = \begin{bmatrix} \mu_R \\ \tau_R \end{bmatrix}$
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$$\bar{L}_\alpha (y_0 \Phi_0 \delta^\alpha_\beta + 2y \Phi^\alpha_\beta) e_R^\beta$$

Scalars		
Φ_0	Φ^α_β	S^α
1	3	2
2	2	1
+1/2	+1/2	0

$$\Phi^\alpha_\beta = \frac{(\sigma_a)^\alpha_\beta}{2} \Phi_a = \frac{1}{2} \begin{bmatrix} \Phi_3 & \sqrt{2}\Phi_+ \\ \sqrt{2}\Phi_- & -\Phi_3 \end{bmatrix}, \quad \Phi_\pm = (\Phi_1 \mp i\Phi_2)/\sqrt{2}$$

- Flavor Fundamental SM singlet Higgs S^α

$$\langle \Phi \rangle \propto \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} : SU(2)_{\mu\tau} \rightarrow U(1)_{L_\mu - L_\tau}$$

$$\langle S \rangle \propto \begin{bmatrix} 0 \\ 1 \end{bmatrix} : SU(2)_{\mu\tau} \rightarrow \text{None}$$

$$y_0 = \frac{\sqrt{2}}{v_0} \frac{M_\mu + M_\tau}{2}$$

$$y = \frac{\sqrt{2}}{v_3} \frac{M_\mu - M_\tau}{2}$$

Symmetry Breaking Sector

- Scalar Potential

$$\begin{aligned} V(\Phi_0, \Phi, S) = & -\mu_S^2 |S|^2 - \mu_0^2 \Phi_{0i}^\dagger \Phi_0^i + 2m_\Phi^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + \lambda_S |S|^4 + \lambda_{0S} |S|^2 \Phi_{0i}^\dagger \Phi_0^i \\ & + 2\lambda_{\Phi S1} |S|^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + 2\lambda_{\Phi S2} S^\dagger [\Phi_i^\dagger, \Phi^i] S \\ & + 2\lambda'_1 (\text{Tr}(\Phi_i^\dagger \Phi^i))^2 + 2\lambda''_1 \text{Tr}(\Phi_i^\dagger \Phi^j) \text{Tr}(\Phi_j^\dagger \Phi^i) + 2\lambda'''_1 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \text{Tr}(\Phi^i \Phi^j) \\ & + \frac{\lambda_2}{2} (\Phi_{0i}^\dagger \Phi_0^i)^2 + 2\lambda_3 \text{Tr}(\Phi_j^\dagger \Phi^j) \Phi_{0i}^\dagger \Phi_0^i + 2\lambda_4 \text{Tr}(\Phi^i \Phi_j^\dagger) \Phi_{0i}^\dagger \Phi_0^j \\ & + \left\{ \lambda'_{\Phi S} S i \sigma_2 [\Phi_i^\dagger, \Phi^i] S + \lambda_5 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \Phi_0^i \Phi_0^j + 4\kappa \Phi_{0i}^\dagger \text{Tr}(\Phi_j^\dagger \Phi^j \Phi^i) \right. \\ & \left. + \lambda_\mu \Phi_{0i}^\dagger S^\dagger \Phi^i S + \lambda'_{\mu 1} \Phi_{0i}^\dagger S i \sigma_2 \Phi^i S + \lambda'_{\mu 2} S i \sigma_2 \Phi_i^\dagger S \Phi_0^i + \text{H.c.} \right\} \end{aligned}$$

Symmetry Breaking Sector

- Scalar Potential

Assume Global Sym. for $S \rightarrow$ a simpler vacuum structure $\rightarrow \Phi$ gets Diagonal VEV

$$\begin{aligned}
 V(\Phi_0, \Phi, S) = & -\mu_S^2 |S|^2 - \mu_0^2 \Phi_{0i}^\dagger \Phi_0^i + 2m_\Phi^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + \lambda_S |S|^4 + \lambda_{0S} |S|^2 \Phi_{0i}^\dagger \Phi_0^i \\
 & + 2\lambda_{\Phi S1} |S|^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + 2\lambda_{\Phi S2} S^\dagger [\Phi_i^\dagger, \Phi^i] S \\
 & + 2\lambda'_1 (\text{Tr}(\Phi_i^\dagger \Phi^i))^2 + 2\lambda''_1 \text{Tr}(\Phi_i^\dagger \Phi^j) \text{Tr}(\Phi_j^\dagger \Phi^i) + 2\lambda'''_1 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \text{Tr}(\Phi^i \Phi^j) \\
 & + \frac{\lambda_2}{2} (\Phi_{0i}^\dagger \Phi_0^i)^2 + 2\lambda_3 \text{Tr}(\Phi_j^\dagger \Phi^j) \Phi_{0i}^\dagger \Phi_0^i + 2\lambda_4 \text{Tr}(\Phi^i \Phi_j^\dagger) \Phi_{0i}^\dagger \Phi_0^j \\
 & + \left\{ \cancel{\lambda'_{\Phi S} S i \sigma_2 [\Phi_i^\dagger, \Phi^i] S} + \lambda_5 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \Phi_0^i \Phi_0^j + 4\kappa \Phi_{0i}^\dagger \text{Tr}(\Phi_j^\dagger \Phi^j \Phi^i) \right. \\
 & \left. + \lambda_\mu \Phi_{0i}^\dagger S^\dagger \Phi^i S + \cancel{\lambda'_{\mu 1} \Phi_{0i}^\dagger S i \sigma_2 \Phi^i S} + \cancel{\lambda'_{\mu 2} S i \sigma_2 \Phi_i^\dagger S \Phi_0^i} + \text{H.c.} \right\}
 \end{aligned}$$

Symmetry Breaking Sector

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Assume Global Sym. for $S \rightarrow$ a simpler vacuum structure $\rightarrow \Phi$ gets Diagonal VEV

$$\begin{aligned}
 V(\Phi_0, \Phi, S) = & -\mu_S^2 |S|^2 - \mu_0^2 \Phi_{0i}^\dagger \Phi_0^i + 2m_\Phi^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + \lambda_S |S|^4 + \lambda_{0S} |S|^2 \Phi_{0i}^\dagger \Phi_0^i \\
 & + 2\lambda_{\Phi S 1} |S|^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + 2\lambda_{\Phi S 2} S^\dagger [\Phi_i^\dagger, \Phi^i] S \\
 & + 2\lambda'_1 (\text{Tr}(\Phi_i^\dagger \Phi^i))^2 + 2\lambda''_1 \text{Tr}(\Phi_i^\dagger \Phi^j) \text{Tr}(\Phi_j^\dagger \Phi^i) + 2\lambda'''_1 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \text{Tr}(\Phi^i \Phi^j) \\
 & + \frac{\lambda_2}{2} (\Phi_{0i}^\dagger \Phi_0^i)^2 + 2\lambda_3 \text{Tr}(\Phi_j^\dagger \Phi^j) \Phi_{0i}^\dagger \Phi_0^i + 2\lambda_4 \text{Tr}(\Phi^i \Phi_j^\dagger) \Phi_{0i}^\dagger \Phi_0^j \\
 & + \left\{ \cancel{\lambda'_{\Phi S} S i \sigma_2 [\Phi_i^\dagger, \Phi^i] S} + \lambda_5 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \Phi_0^i \Phi_0^j + 4\kappa \Phi_{0i}^\dagger \text{Tr}(\Phi_j^\dagger \Phi^j \Phi^i) \right. \\
 & \left. + \lambda_\mu \Phi_{0i}^\dagger S^\dagger \Phi^i S + \cancel{\lambda'_{\mu 1} \Phi_{0i}^\dagger S i \sigma_2 \Phi^i S} + \cancel{\lambda'_{\mu 2} S i \sigma_2 \Phi_i^\dagger S \Phi_0^i} + \text{H.c.} \right\}
 \end{aligned}$$

$$\langle \Phi \rangle = \begin{pmatrix} v_3 & 0 \\ 0 & v_3 \end{pmatrix} : SU(2)_{\mu\tau} \xrightarrow{\langle \Phi \rangle} U(1)_{\mu-\tau}$$

$$\begin{cases} (\Phi_0, \Phi_\pm) & \text{ZHDM} \\ (\Phi_+, \Phi_-) & \text{ZPDM} \dots \text{"flavor charged"} \end{cases}$$

\uparrow scalar (w/o VEV)

Symmetry Breaking Sector

- **Scalar Potential** Assume Global Sym. for S → a simpler vacuum structure → Φ gets Diagonal VEV

$$\begin{aligned}
 V(\Phi_0, \Phi, S) = & -\mu_S^2 |S|^2 - \mu_0^2 \Phi_{0i}^\dagger \Phi_0^i + 2m_\Phi^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + \lambda_S |S|^4 + \lambda_{0S} |S|^2 \Phi_{0i}^\dagger \Phi_0^i \\
 & + 2\lambda_{\Phi S1} |S|^2 \text{Tr}(\Phi_i^\dagger \Phi^i) + 2\lambda_{\Phi S2} S^\dagger [\Phi_i^\dagger, \Phi^i] S \\
 & + 2\lambda'_1 (\text{Tr}(\Phi_i^\dagger \Phi^i))^2 + 2\lambda''_1 \text{Tr}(\Phi_i^\dagger \Phi^j) \text{Tr}(\Phi_j^\dagger \Phi^i) + 2\lambda'''_1 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \text{Tr}(\Phi^i \Phi^j) \\
 & + \frac{\lambda_2}{2} (\Phi_{0i}^\dagger \Phi_0^i)^2 + 2\lambda_3 \text{Tr}(\Phi_j^\dagger \Phi^j) \Phi_{0i}^\dagger \Phi_0^i + 2\lambda_4 \text{Tr}(\Phi^i \Phi_j^\dagger) \Phi_{0i}^\dagger \Phi_0^j \\
 & + \left\{ \cancel{\lambda'_{\Phi S} S i \sigma_2 [\Phi_i^\dagger, \Phi^i] S} + \lambda_5 \text{Tr}(\Phi_i^\dagger \Phi_j^\dagger) \Phi_0^i \Phi_0^j + 4\kappa \Phi_{0i}^\dagger \text{Tr}(\Phi_j^\dagger \Phi^j \Phi^i) \right. \\
 & \left. + \lambda_\mu \Phi_{0i}^\dagger S^\dagger \Phi^i S + \cancel{\lambda'_{\mu 1} \Phi_{0i}^\dagger S i \sigma_2 \Phi^i S} + \cancel{\lambda'_{\mu 2} S i \sigma_2 \Phi_i^\dagger S \Phi_0^i} + \text{H.c.} \right\}
 \end{aligned}$$

- **Classification of Scalar fields** (Assume CP inv. Scalar sector for simplicity)

Original field	Mass eigenstates	Electric charge	CP	Lepton Flavor
(S_0^h, h_3, h_0)	(s, H, h)	0	even	conserving
(S_0^z, z_3, z_0)	$((X_3)_L, Z_L, A)$	0	odd	conserving
(ω_3^+, ω_0^+)	$((W^+)_L, H^+)$	+1	—	conserving
$(S_+, \phi_+^0, \phi_-^{0*})$	$((X_+)_L, H_+, h_+)$	0	—	changing
(ϕ_+^+, ϕ_-^+)	(ϕ_+^+, ϕ_-^+)	+1	—	changing

Flavor Changing Scalar

- New interactions

$$\tan \beta \equiv \frac{v_0}{v_3} \quad \tan \beta' \equiv \frac{2v_3}{v_S} \ll 1$$

$$\frac{M_\mu - M_\tau}{\sqrt{2} v c_\beta} \left\{ s_{\beta'} \bar{\mu} \tau S_+ + \underbrace{\bar{\mu} (c_{\alpha'} c_{\beta'} - s_{\alpha'} \gamma_5) \tau H_+}_{\text{NG boson}} - \bar{\mu} (s_{\alpha'} c_{\beta'} + c_{\alpha'} \gamma_5) \tau h_+ \right\}$$

Original field	Mass eigenstates	Electric charge	CP	Lepton Flavor
(S_0^h, h_3, h_0)	(s, H, h)	0	even	conserving
(S_0^z, z_3, z_0)	$((X_3)_L, Z_L, A)$	0	odd	conserving
(ω_3^+, ω_0^+)	$((W^+)_L, H^+)$	+1	—	conserving
$(S_+, \phi_+^0, \phi_-^{0*})$	$((\text{NG})_L, H_+, h_+)$	0	—	changing
(ϕ_+^+, ϕ_-^+)	(ϕ_+^+, ϕ_-^+)	+1	—	changing

Flavor Changing Scalar

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- Muon g-2

$$R = M_{h_\pm}^2 / M_{H_\pm}^2$$

$$\Delta a_{\mu\pm} \simeq \frac{1}{48\pi^2} \frac{M_\mu^2}{M_{h_\pm}^2} \left(\frac{M_\mu - M_\tau}{v} \sqrt{1 + t_\beta^2} \right)^2 \left\{ 1 + R + 3 \frac{M_\tau}{M_\mu} \underbrace{c_{2\alpha'}}_{\text{Chirality enhancement}} \left[(1 - R) \left(\ln \frac{M_\mu^2}{M_{h_\pm}^2} + \frac{3}{2} \right) - R \ln R \right] \right\}$$

Yukawa w/ $\tan\beta$ enhancement

Mixing angle dependence

Original field	Mass eigenstates	Electric charge	CP	Lepton Flavor
(S_0^h, h_3, h_0)	(s, H, h)	0	even	conserving
(S_0^z, z_3, z_0)	$((X_3)_L, Z_L, A)$	0	odd	conserving
(ω_3^+, ω_0^+)	$((W^+)_L, H^+)$	+1	—	conserving
$(S_+, \phi_+^0, \phi_-^{0*})$	$((\text{NG})_L, H_+, h_+)$	0	—	changing
(ϕ_+^+, ϕ_-^+)	(ϕ_+^+, ϕ_-^+)	+1	—	changing

Flavor Changing Scalar

- New interactions

$$\tan \beta \equiv \frac{v_0}{v_3} \quad \tan \beta' \equiv \frac{2v_3}{v_S} \ll 1$$

$$\frac{M_\mu - M_\tau}{\sqrt{2} v c_\beta} \left\{ s_{\beta'} \bar{\mu} \tau S_+ + \underbrace{\bar{\mu} (c_{\alpha'} c_{\beta'} - s_{\alpha'} \gamma_5) \tau H_+}_{\text{NG boson}} - \bar{\mu} (s_{\alpha'} c_{\beta'} + c_{\alpha'} \gamma_5) \tau h_+ \right\}$$

- Muon g-2

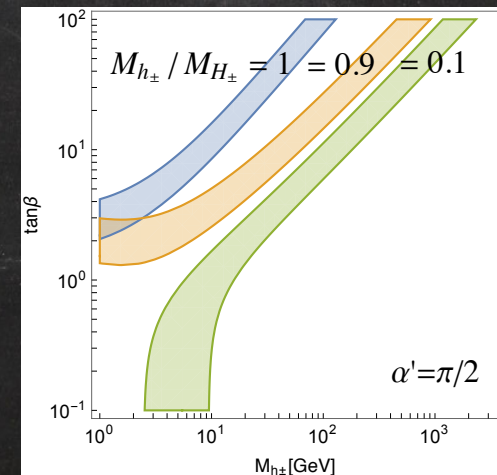
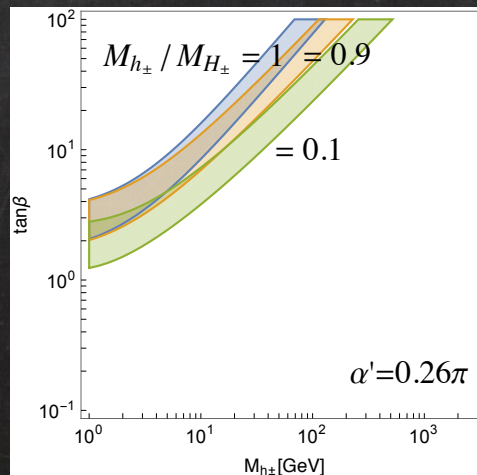
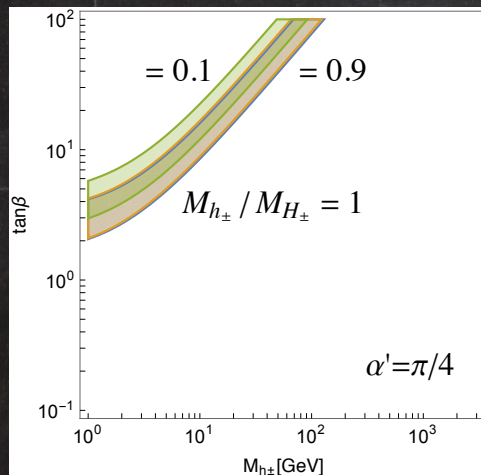
Chirality enhancement

$$R = M_{h_\pm}^2 / M_{H_\pm}^2$$

$$\Delta a_{\mu\pm} \simeq \frac{1}{48\pi^2} \frac{M_\mu^2}{M_{h_\pm}^2} \left(\frac{M_\mu - M_\tau}{v} \sqrt{1 + t_\beta^2} \right)^2 \left\{ 1 + R + 3 \frac{M_\tau}{M_\mu} c_{2\alpha'} \left[(1 - R) \left(\ln \frac{M_\mu^2}{M_{h_\pm}^2} + \frac{3}{2} \right) - R \ln R \right] \right\}$$

Yukawa w/ $\tan\beta$ enhancement

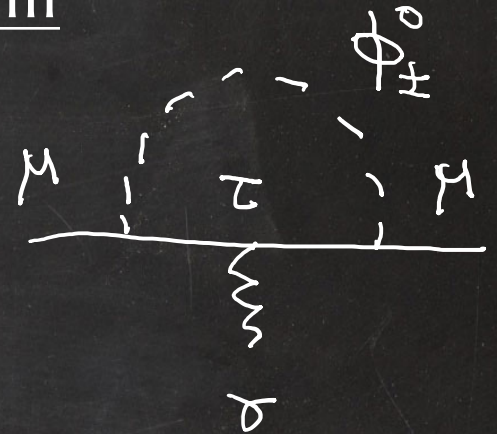
Mixing angle dependence



Comparison w/ 2HDM-III

2S Φ M ... Flavor charged

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \underbrace{\left(\frac{M_\tau}{v} t_\beta \right)^2}_{\text{coupling}} \times \underbrace{\left(\frac{M_\tau}{M_\mu} \right)}_{\text{Chirality Enhancement}}$$



2HDM-III (Omura, Senaha, Tobe)

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \underbrace{(P_{\mu\tau} P_{\tau\mu}^*)}_{\text{coupling}} \times \underbrace{\left(\frac{M_\tau}{M_\mu} \right)}_{\text{Chirality Enhancement}}$$



$P_{\mu\tau} \neq P_{\tau\mu}^*$ (τ independent complex parameters)

$$\propto \left(\frac{S_{R-d}}{M_H^2} - \frac{1}{M_A^2} \right) \text{ relate to higgs phys.}$$

\Downarrow
EDM constraint

$$\tan \beta \equiv \frac{v_0}{v_3}$$

Flavor Conserving Scalar

- New interactions

$$- \left(\overline{\mu}_L \left\{ \frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right\} \mu_R + \overline{\tau}_L \left\{ \frac{M_\tau}{v} \frac{1}{t_{2\beta}} + \frac{M_\mu}{v} \frac{1}{s_{2\beta}} \right\} \tau_R \right) (H + i A)$$

Set to 0 (simplicity)

SM-like limit $\sin(\beta-\alpha)=1$ (simplicity)

Original field	Mass eigenstates	Electric charge	CP	Lepton Flavor
(S_0^h, h_3, h_0)	$(\cancel{s}, H, \cancel{h})$	0	even	conserving
(S_0^z, z_3, z_0)	$((\cancel{X}_3)_L, \cancel{X}_L, A)$	0	odd	conserving
(ω_3^+, ω_0^+)	$((W^+)_L, H^+)$	+1	—	conserving
$(S_+, \phi_+^0, \phi_-^{0*})$	$((X_+)_L, H_+, h_+)$	0	—	changing
(ϕ_+^+, ϕ_-^+)	(ϕ_+^+, ϕ_-^+)	+1	—	changing

$$\tan \beta \equiv \frac{v_0}{v_3}$$

Flavor Conserving Scalar

- New interactions

$$- \left(\overline{\mu}_L \left\{ \frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right\} \mu_R + \overline{\tau}_L \left\{ \frac{M_\tau}{v} \frac{1}{t_{2\beta}} + \frac{M_\mu}{v} \frac{1}{s_{2\beta}} \right\} \tau_R \right) (H + i A)$$

- Muon g-2

$$\Delta a_{\mu 0} \simeq \frac{1}{12\pi^2} \frac{M_\mu^2}{M_\varphi^2} \left(\frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right)^2, \quad (M_\varphi \equiv M_H = M_A)$$

No Chirality enhancement

$$1/t_{2\beta} = (1 - t_\beta^2)/2t_\beta$$

$$1/s_{2\beta} = (1 + t_\beta^2)/2t_\beta$$

Yukawa w/ $\tan\beta$ enhancement

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- New interactions

$$- \left(\overline{\mu}_L \left\{ \frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right\} \mu_R + \overline{\tau}_L \left\{ \frac{M_\tau}{v} \frac{1}{t_{2\beta}} + \frac{M_\mu}{v} \frac{1}{s_{2\beta}} \right\} \tau_R \right) (H + iA)$$

- Muon $g-2$

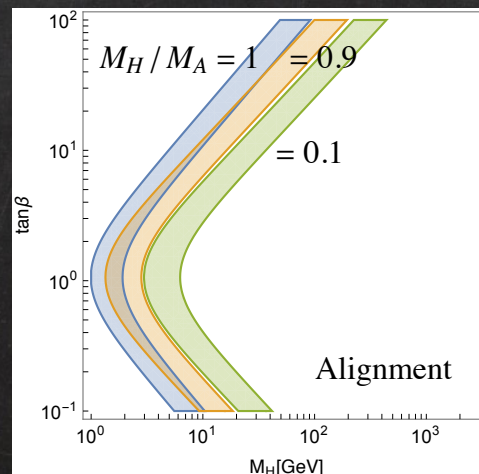
$$\Delta a_{\mu 0} \simeq \frac{1}{12\pi^2} \frac{M_\mu^2}{M_\varphi^2} \left(\frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right)^2, \quad (M_\varphi \equiv M_H = M_A)$$

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Yukawa w/ $\tan\beta$ enhancement



Flavor Conserving Scalar

$$\tan \beta \equiv \frac{v_0}{v_3}$$

- New interactions

$$- \left(\overline{\mu}_L \left\{ \frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right\} \mu_R + \overline{\tau}_L \left\{ \frac{M_\tau}{v} \frac{1}{t_{2\beta}} + \frac{M_\mu}{v} \frac{1}{s_{2\beta}} \right\} \tau_R \right) (H + i A)$$

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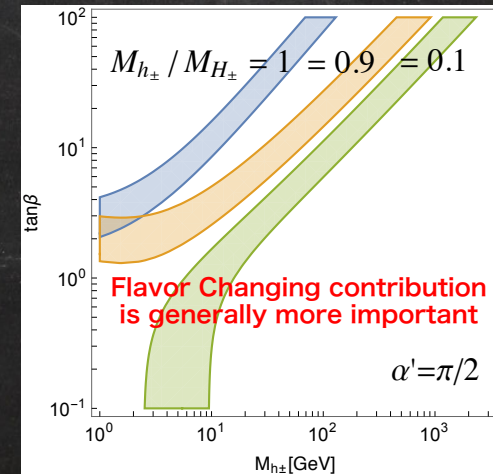
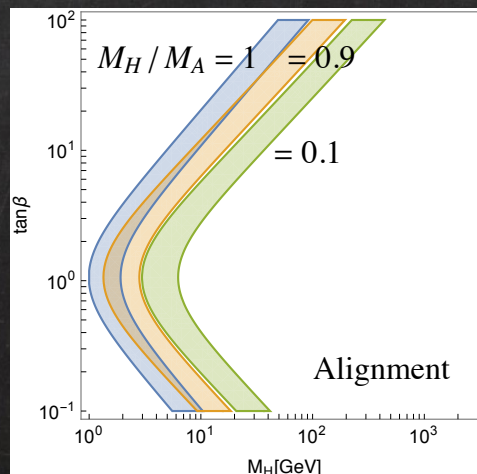
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Yukawa w/ $\tan\beta$ enhancement

$$M \sim M_\varphi \sim M_{h_\pm}$$



Comparison w/ 2HDM-X

2S \mathcal{P} M ... Flavor charged

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \left(\frac{M_\tau}{v} t_\beta \right)^2 \times \left(\frac{M_\tau}{M_\mu} \right)$$

coupling

Chirality
Enhancement



2HDM-X (Abe, Sato, Yagyu)

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \left(\frac{M_M}{v} t_\beta \right)^2 \times 1$$

light $\phi (= A)$
 $\sim 50 \text{ GeV}$

moderate t_β
 $\sim 10^2$



$$\left[\frac{M_\tau}{v} t_\beta \lesssim 4\pi \right]$$

Comparison w/ 2HDM-X

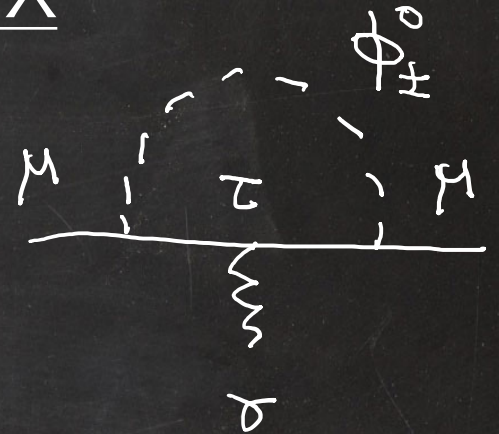
2S Φ M ... Flavor charged

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \left(\frac{M_\tau}{v} t_\beta \right)^2 \times \left(\frac{M_\tau}{M_\mu} \right)$$

$\times \left(\frac{M_\tau}{M_\mu} \right)^3$
 $\sim 5,000$

coupling

Chirality
Enhancement



2HDM-X

(Abe, Sato, Yagyu)

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \left(\frac{M_M}{v} t_\beta \right)^2 \times 1$$

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 $\sim 50 \text{ GeV}$

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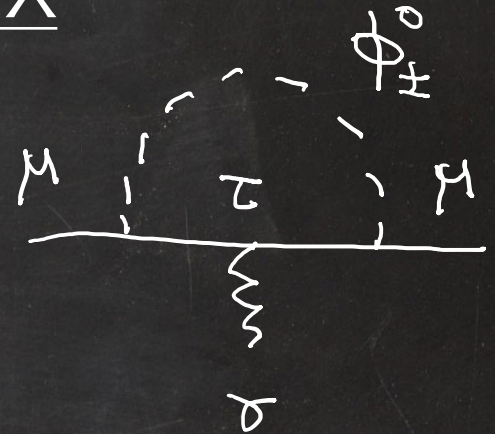
Comparison w/ 2HDM-X

2S \mathcal{P} M ... Flavor charged

$$\Delta a_\mu \sim \frac{1}{(4\pi)^2} \frac{M_\mu^2}{M_\phi^2} \times \left(\frac{M_\tau}{v} t_\beta \right)^2 \times \left(\frac{M_\tau}{M_\mu} \right)$$

coupling

Chirality
Enhancement



\mu-2H \mathcal{P} M (Abe, Sato, Yagyu)

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heavy ϕ
 $\sim 500 \text{ GeV}$

large t_β
 $\sim 10^3$

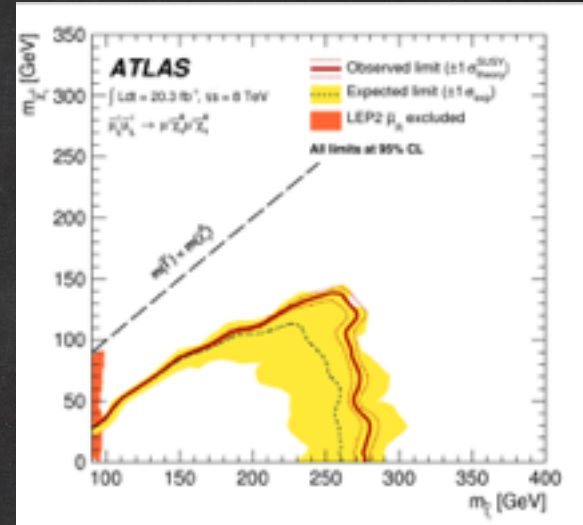
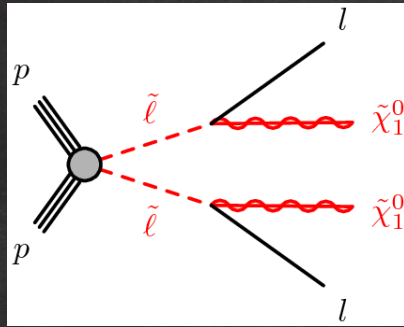
$$\left[\frac{M_M}{v} t_\beta \lesssim 4\pi \right]$$



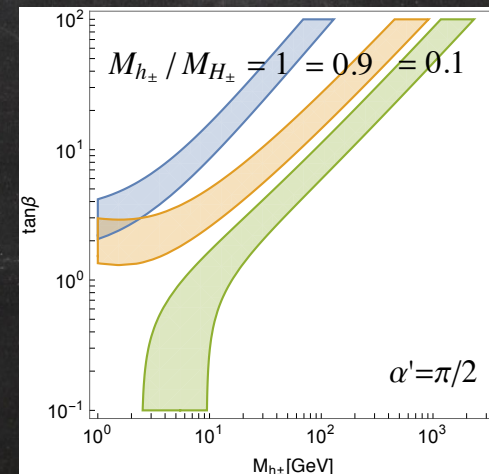
Slepton Search

- Pair Production of Charged Higgs

H^\pm has the same quantum charge with left-handed slepton



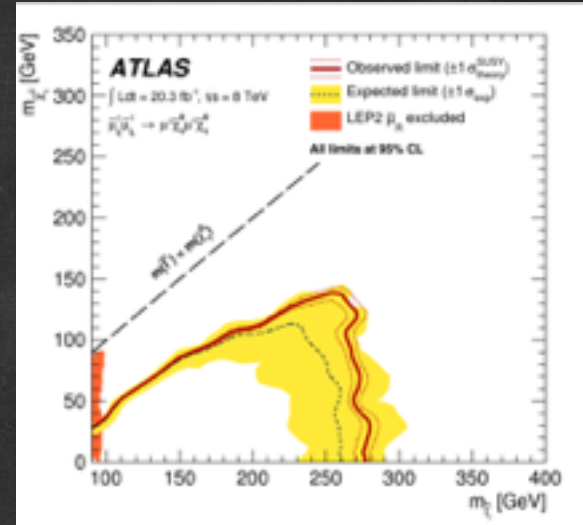
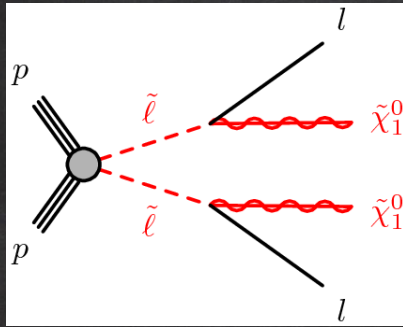
$$M \sim M_\varphi \sim M_{h_\pm}$$



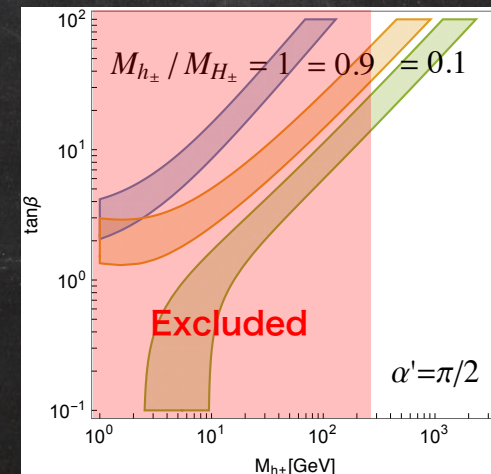
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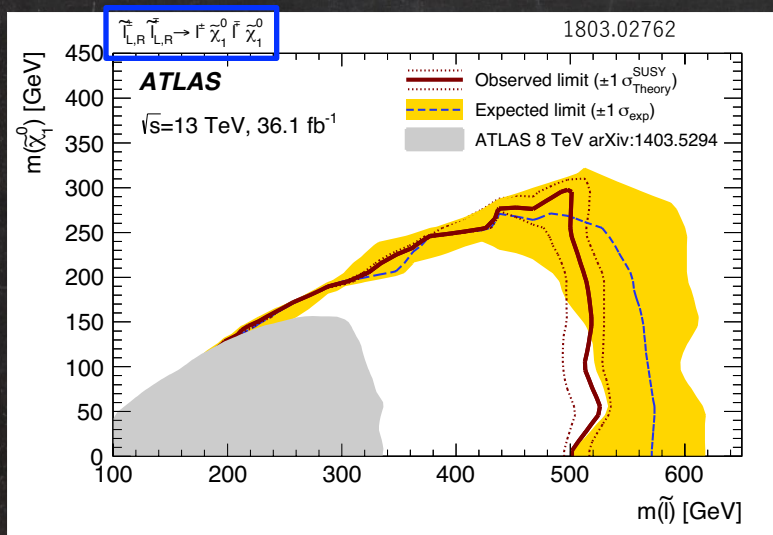
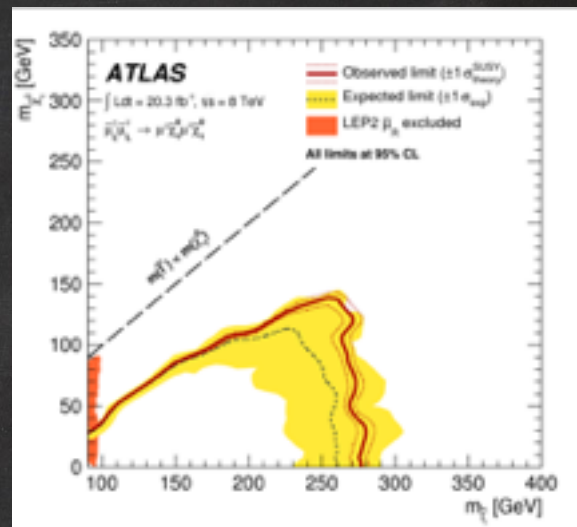
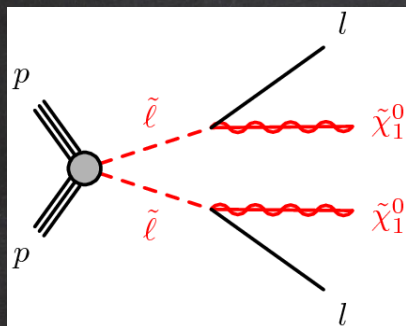
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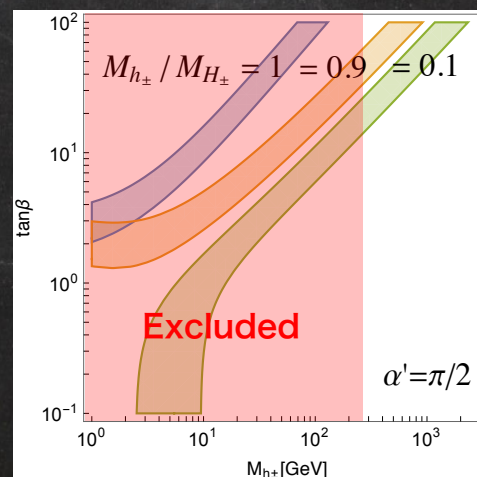
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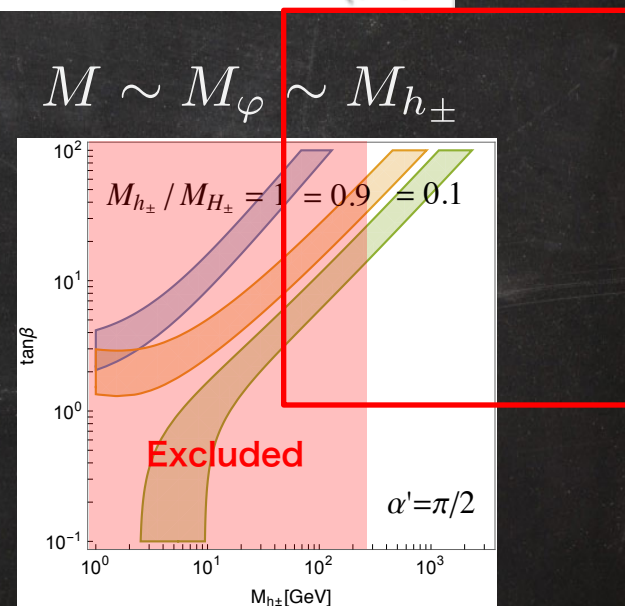
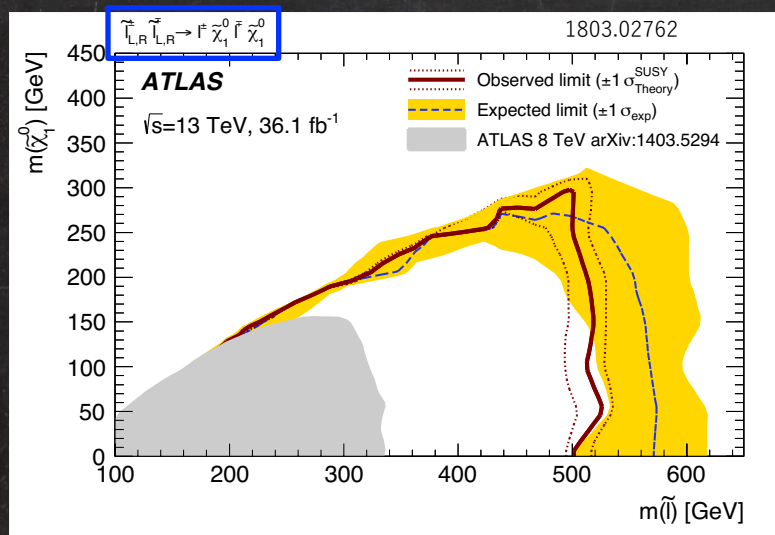
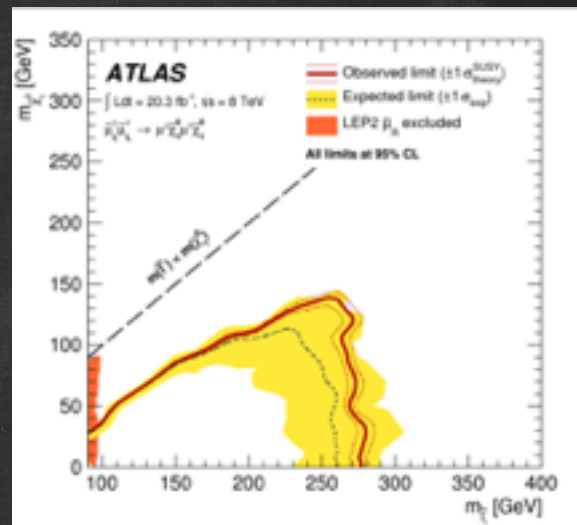
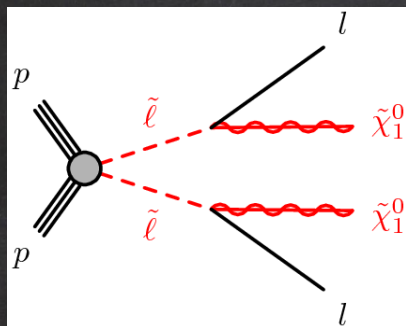
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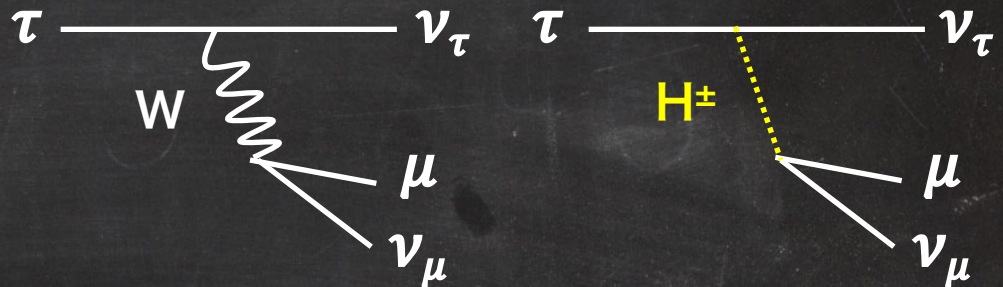
τ Michel Decay

- Charged Scalar interactions

$$+i\sqrt{2}\bar{\nu}_{\mu L}\left\{\frac{M_{\mu}}{v}\frac{1}{t_{2\beta}}+\frac{M_{\tau}}{v}\frac{1}{s_{2\beta}}\right\}\mu_R H^+ + i\sqrt{2}\bar{\nu}_{\tau L}\left\{\frac{M_{\tau}}{v}\frac{1}{t_{2\beta}}+\frac{M_{\mu}}{v}\frac{1}{s_{2\beta}}\right\}\tau_R H^+$$

- Lepton Universality

$$R_{\tau\rightarrow\mu/\tau\rightarrow e} = \frac{\Gamma(\tau \rightarrow \mu\bar{\nu}_{\mu}\nu_{\tau})}{\Gamma(\tau \rightarrow e\bar{\nu}_e\nu_{\tau})}$$



τ Michel Decay

- Charged Scalar interactions

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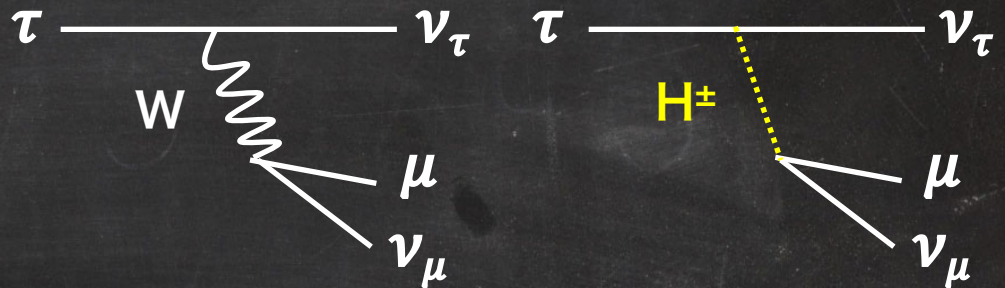
- Lepton Universality

$$R_{\tau\rightarrow\mu/\tau\rightarrow e} = \frac{\Gamma(\tau\rightarrow\mu\bar{\nu}_{\mu}\nu_{\tau})}{\Gamma(\tau\rightarrow e\bar{\nu}_e\nu_{\tau})}$$

$$\eta^{H^{\pm}} \equiv \frac{R^{H^{\pm}}_{\tau\rightarrow\mu/\tau\rightarrow e}}{R^{\text{SM}}_{\tau\rightarrow\mu/\tau\rightarrow e}}$$

$$\approx \left\{ 1 + \frac{1}{4} \left[-\frac{v^2}{M_{H^{\pm}}^2} \left(\frac{M_{\mu}}{v} \frac{1}{t_{2\beta}} + \frac{M_{\tau}}{v} \frac{1}{s_{2\beta}} \right) \left(\frac{M_{\tau}}{v} \frac{1}{t_{2\beta}} + \frac{M_{\mu}}{v} \frac{1}{s_{2\beta}} \right) \right]^2 + 2 \left[-\frac{v^2}{M_{H^{\pm}}^2} \left(\frac{M_{\mu}}{v} \frac{1}{t_{2\beta}} + \frac{M_{\tau}}{v} \frac{1}{s_{2\beta}} \right) \left(\frac{M_{\tau}}{v} \frac{1}{t_{2\beta}} + \frac{M_{\mu}}{v} \frac{1}{s_{2\beta}} \right) \right] \frac{M_{\mu} g(\rho^{-2})}{M_{\tau} f(\rho^{-2})} \right\}$$

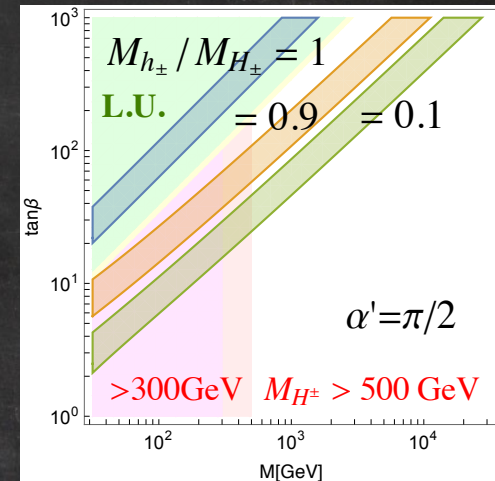
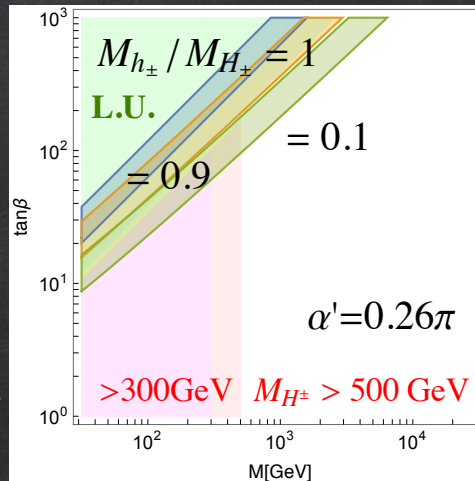
Chirality Suppressed Interference



τ Michel Decay

- Charged Scalar interactions

$$M^2 \equiv M_H^2 = M_A^2 = M_{H^\pm}^2 = M_{h^\pm}^2$$



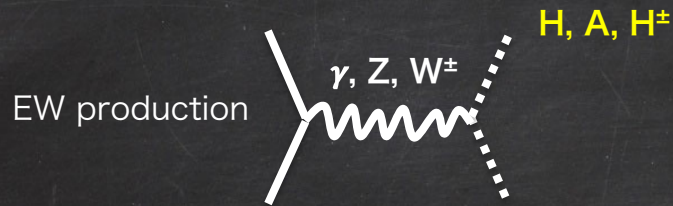
$$\eta^{H^\pm} \equiv \frac{R_{\tau \rightarrow \mu / \tau \rightarrow e}^{H^\pm}}{R_{\tau \rightarrow \mu / \tau \rightarrow e}^{\text{SM}}}$$

$$\approx \left\{ 1 + \frac{1}{4} \left[-\frac{v^2}{M_{H^\pm}^2} \left(\frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right) \left(\frac{M_\tau}{v} \frac{1}{t_{2\beta}} + \frac{M_\mu}{v} \frac{1}{s_{2\beta}} \right) \right]^2 + 2 \left[-\frac{v^2}{M_{H^\pm}^2} \left(\frac{M_\mu}{v} \frac{1}{t_{2\beta}} + \frac{M_\tau}{v} \frac{1}{s_{2\beta}} \right) \left(\frac{M_\tau}{v} \frac{1}{t_{2\beta}} + \frac{M_\mu}{v} \frac{1}{s_{2\beta}} \right) \right] \frac{M_\mu g(\rho^{-2})}{M_\tau f(\rho^{-2})} \right\}$$

Chirality Suppressed Interference

Collider Signature

- Flavor conserving μ - τ Specific Higgs bosons



$$\mathcal{B}(H/A \rightarrow \tau\tau) = \mathcal{B}(H/A \rightarrow \mu\mu) = 50\%$$

$$\mathcal{B}(H^\pm \rightarrow \tau\nu) = \mathcal{B}(H^\pm \rightarrow \mu\nu) = 50\%$$

- Flavor changing Scalar bosons



$$\mathcal{B}(h_\pm/H_\pm \rightarrow \tau\mu) = 100\%$$

$$\mathcal{B}(\phi_\pm^\pm \rightarrow \tau\nu) = \mathcal{B}(\phi_\pm^\pm \rightarrow \mu\nu) = 50\%$$

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$(S_+, \phi_+^0, \phi_-^{0*})$	$((\cancel{NG})_L, H_+, h_+)$	0	—	changing
(ϕ_+^+, ϕ_-^+)	(ϕ_+^+, ϕ_-^+)	+1	—	changing

Implication for Neutrinos

Witten Anomaly ?

- Odd Number of Weyl Fermions \rightarrow Witten Anomaly

	SM leptons			
	$L = L_e$	$L^\alpha = \begin{matrix} L_\mu \\ L_\tau \end{matrix}$	e_R	$e_R^\alpha = \begin{matrix} \mu_R \\ \tau_R \end{matrix}$
$SU(2)_{\mu\tau}$	1	2	1	2
$SU(2)_L$	2	2	1	1
$U(1)_Y$	-1/2	-1/2	-1	-1

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- Sign ambiguity $\int (d\psi d\bar{\psi})_{\text{Weyl}} \exp(\bar{\psi} i \not{D} \psi) = \pm (\det i \not{D})^{1/2}$

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- Sign ambiguity $\int (d\psi d\bar{\psi})_{\text{Weyl}} \exp(\bar{\psi} i \not{D} \psi) = \pm (\det i \not{D})^{1/2}$

Sign is not invariant unde the topologically non-trivial transf.

$$[\det i \not{D}(A_\mu)]^{1/2} = -[\det i \not{D}(A_\mu^U)]^{1/2}$$

$$A_\mu^U = U^{-1} A_\mu U - i U^{-1} \partial_\mu U$$

Witten Anomaly ?

- Odd Number of Weyl Fermions \rightarrow Witten Anomaly

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Sign is not invariant unde the topologically non-trivial transf.

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$$A_\mu^U = U^{-1} A_\mu U - i U^{-1} \partial_\mu U$$

No such theory after path integral of gauge fields

$$Z = \int dA_\mu (\det i \not{D})^{1/2} \exp(-\frac{1}{2g^2} \int d^4x \text{Tr} F_{\mu\nu}^2) \rightarrow 0$$

Witten Anomaly ?

- Odd Number of Weyl Fermions \rightarrow Witten Anomaly

	SM leptons			
	$L = L_e$	$L^\alpha = \begin{bmatrix} L_\mu \\ L_\tau \end{bmatrix}$	e_R	$e_R^\alpha = \begin{bmatrix} \mu_R \\ \tau_R \end{bmatrix}$
$SU(2)_{\mu\tau}$	1	2	1	2
$SU(2)_L$	2	2	1	1
$U(1)_Y$	-1/2	-1/2	-1	-1

- Sign ambiguity $\int (d\psi d\bar{\psi})_{\text{Weyl}} \exp(\bar{\psi} i \not{D} \psi) = \pm (\det i \not{D})^{1/2}$

Sign is not invariant unde the topologically non-trivial transf.

$$[\det i \not{D}(A_\mu)]^{1/2} = -[\det i \not{D}(A_\mu^U)]^{1/2}$$

$$A_\mu^U = U^{-1} A_\mu U - i U^{-1} \partial_\mu U$$

No such theory after path integral of gauge fields

$$Z = \int dA_\mu (\det i \not{D})^{1/2} \exp\left(-\frac{1}{2g^2} \int d^4x \text{Tr} F_{\mu\nu}^2\right) \rightarrow 0$$

\rightarrow A simple solution : introduce a flavor doublet N_R

Neutrino Mass

- ν Dirac Yukawa $+ \bar{L} \lambda_e \widetilde{\Phi}_0 N_R + \bar{L}_\alpha (\lambda_+ \widetilde{\Phi}_0 \delta^\alpha_\beta + 2\lambda_- \widetilde{\Phi}^\alpha_\beta) N_R^\beta$
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Φ_0	Φ^α_β	S^α	Σ^α_β
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ν mass and mixing are predicted in a specific parameter region

[The prediction is the same as in the $U(1)_{\mu\tau}$ model]

Asai-Hamaguchi-Nagata (17)

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- Seesaw Relation : $(M_\nu)^{-1} = -(M_D^{-1}) M_N M_D^{-1}$

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Strong correlation among v params

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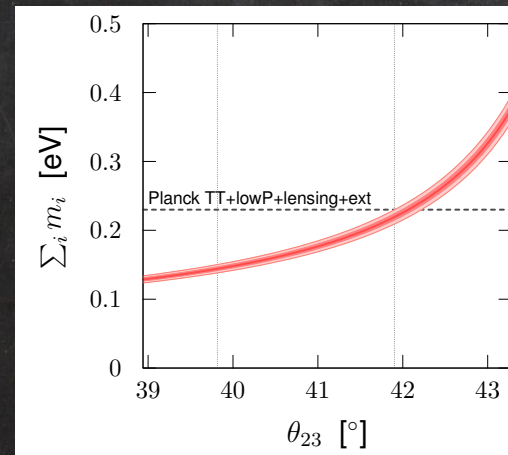
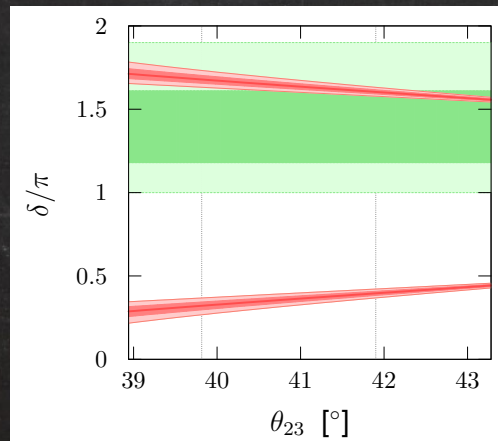
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Red : correlation
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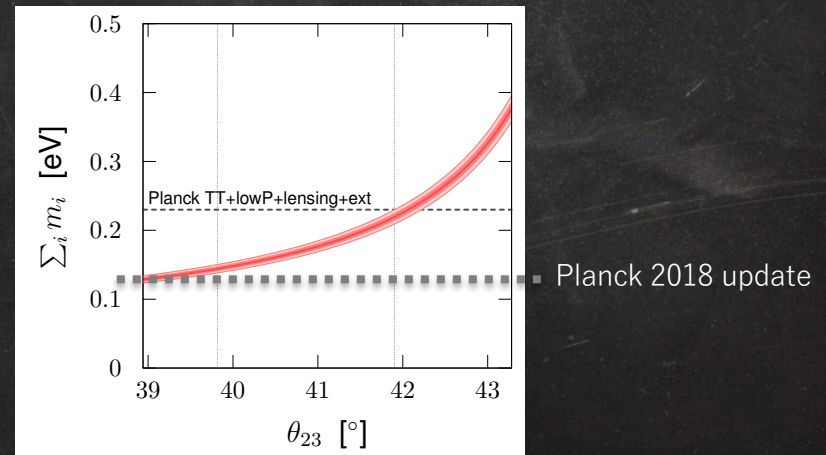
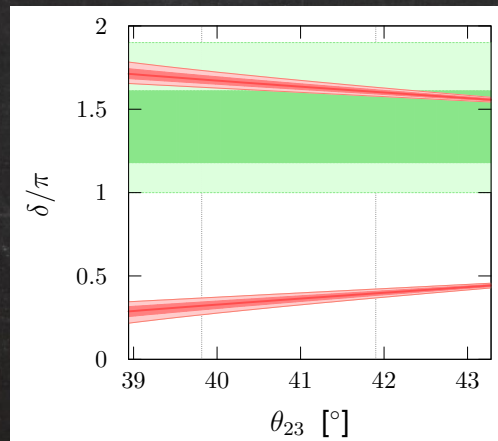
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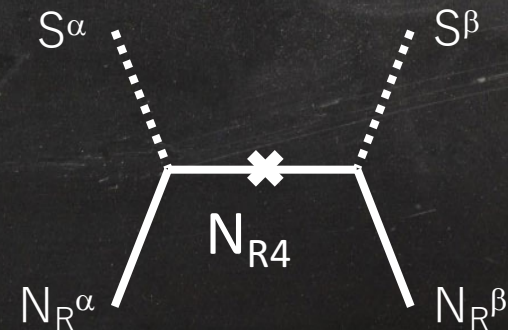
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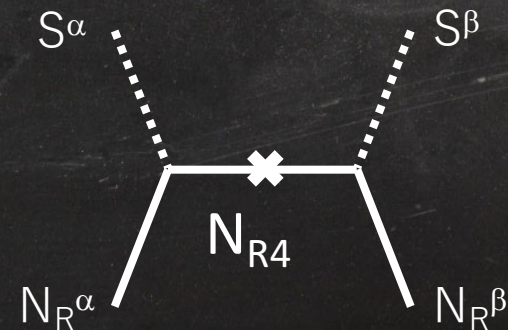
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Summary

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 - Muon $g-2$ anomaly
 - $U(1)_{\mu\tau}$ Model
- $SU(2)_{\mu\tau}$ Model
 - X_3 and X_{\pm}
 - Higgs Sector
 - Michel Decay
 - Collider Signature
 - Neutrino Mass
- Conclusion

